Going global

Combining electroweak precision, diboson, Higgs, and top data to search for new physics

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High Energy Physics Seminar, Imperial College London

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[J. Ellis, M. Madigan, KM, V. Sanz & T. You; JHEP 04 (2021) 279]

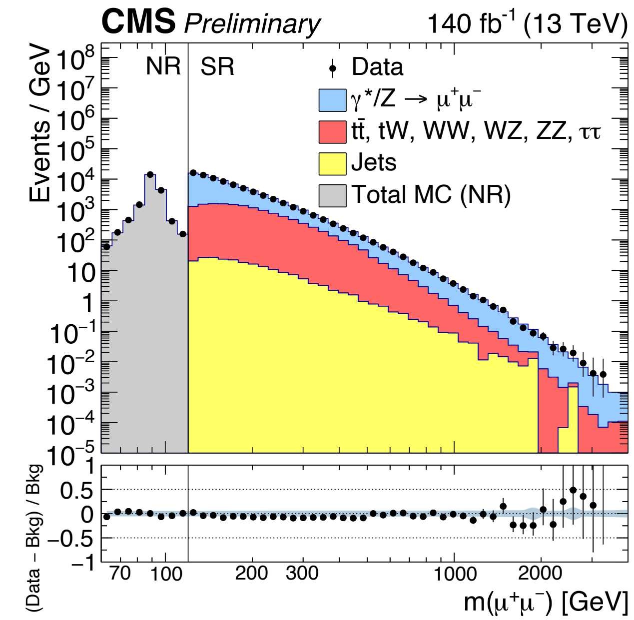
`fitmaker` https://gitlab.com/kenmimasu/fitrepo

[G. Durieux, C. Degrande, F. Maltoni, KM, C. Zhang, E. Vryonidou; PRD 103 (2021) 9, 096024]

`SMEFTatNLO` http://feynrules.irmp.ucl.ac.be/wiki/SMEFTatNLO

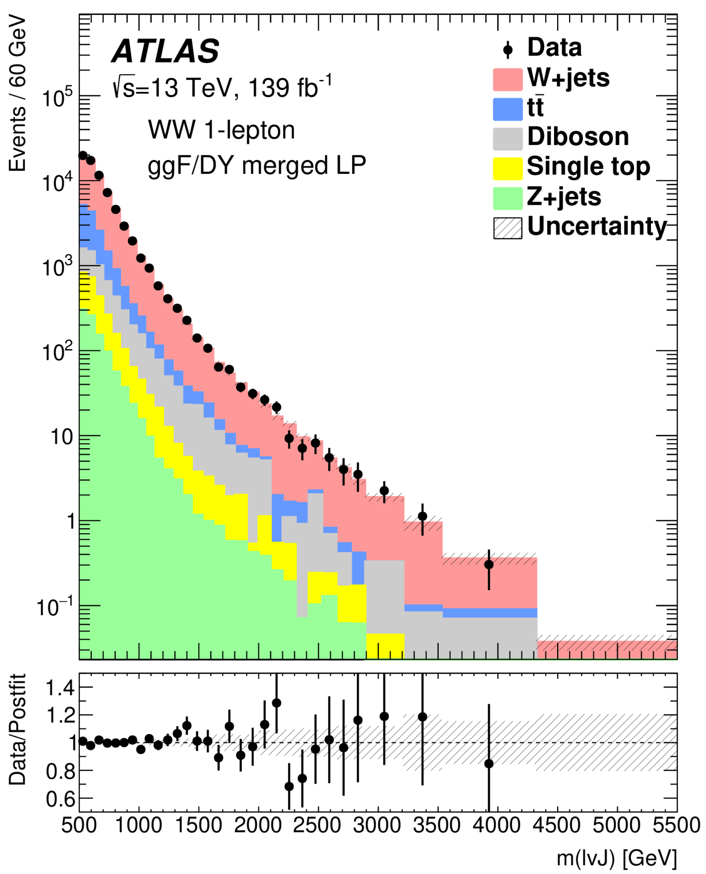
Where are we?

[CMS-PAS-EXO-19-019]

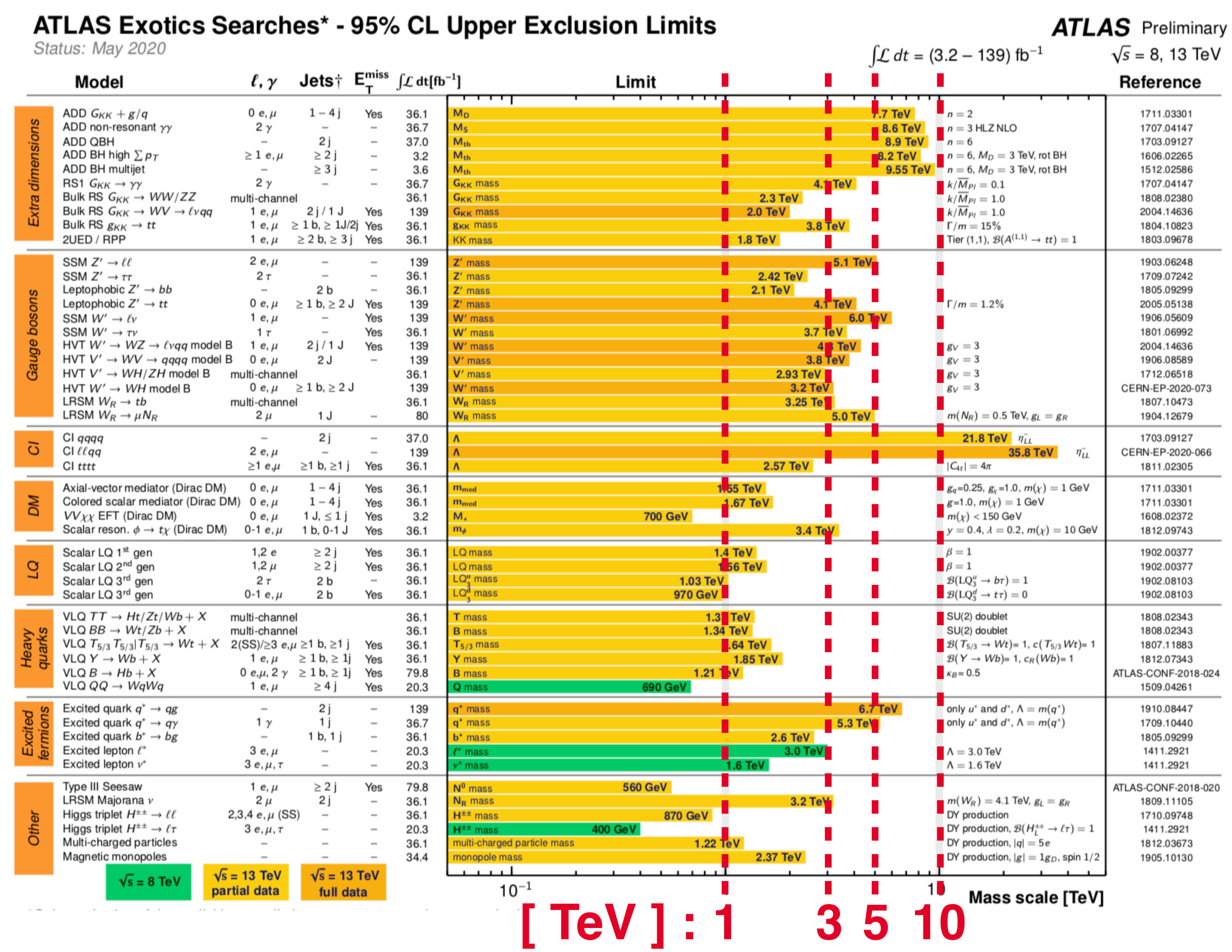


10 years since the start of LHC Run 1

- No clear sign of new physics at the TeV scale
- Direct searches are saturating the energy frontier



[CERN-EP-2020-049]



What have we learnt?

BSM states are too…

Weakly coupled *rate limited*	Room for improvement with increasing luminosity Still 20x more data to come
Exotic *we aren't looking in the right place*	Limited by our creativity Work for theorists & experimentalists to motivate & enable searches for new signatures
Heavy *kinematically out of reach*	Worst-case scenario from direct search point of view Complemented by indirect searches

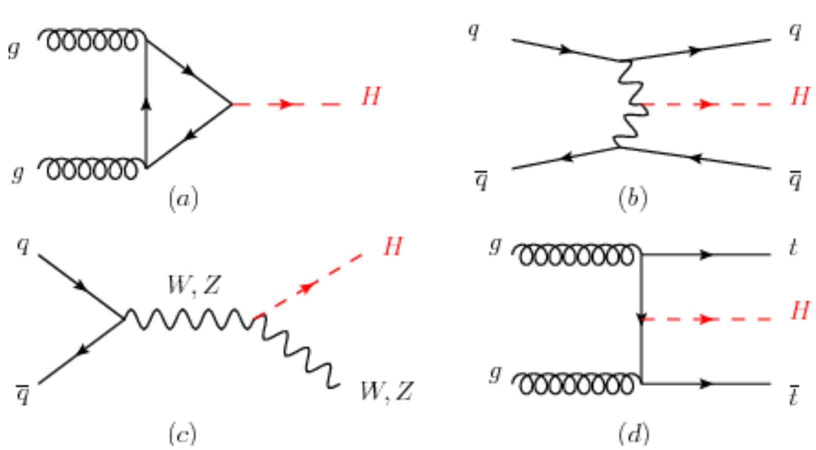
A tremendous amount about the SM!

- Higgs discovery & properties ⇒ precision LHC programme

The LHC explorer

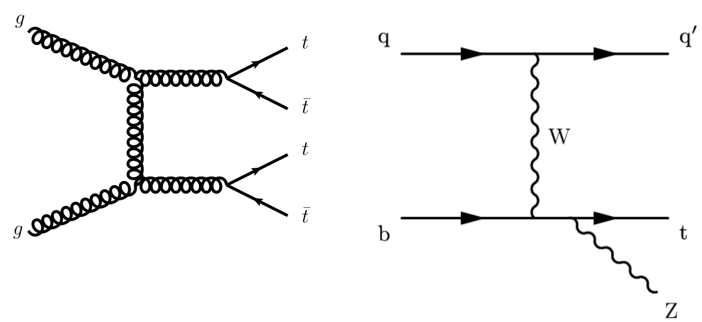
Many new processes observed at the LHC for the first time

Main Higgs production modes



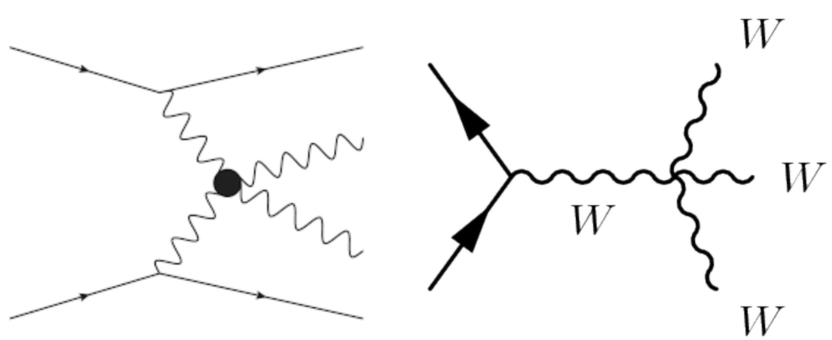
ggF, VH, VBF, ttH

Rare top production



tttt, ttbb   *ttV, tW, tZ*

Weak boson scattering



VBS, VVV

Each opens a new window, through which we can

- Improve our understanding of the SM
- Search for new physics via new interactions

The SM is broken

Theory & matter content rich with symmetry & structure

$SU(3)_C \times SU(2)_L \times U(1)_Y$

$$\varphi = \begin{pmatrix} G^+ \\ v + h + iG^0 \end{pmatrix} : \mathbf{2}_{\frac{1}{2}}$$

Electroweak symmetry breaking

- Offers a **parametrisation**: lacks dynamical origin for the weak scale

Symmetry ↔ Constraints/Relations

$$y_f \, \bar{F}_L f_R \, \varphi \qquad (D^\mu \varphi)^\dagger \, (D_\mu \varphi)$$

Mass ↔ Higgs coupling

$$\frac{1}{4} W^a_{\mu\nu} W^{\mu\nu}_a \qquad i\bar{F} \not{D} F$$

Self-interactions ↔ Gauge currents

New physics can **indirectly** perturb this delicate balance

The indirect way

"…the direct method may be used…but indirect methods will be needed in order to secure victory."

"…there are not more than two methods of attack – the direct and the indirect;…Who can exhaust the possibilities of their combination?"

— Sun Tzu, *The Art of War*

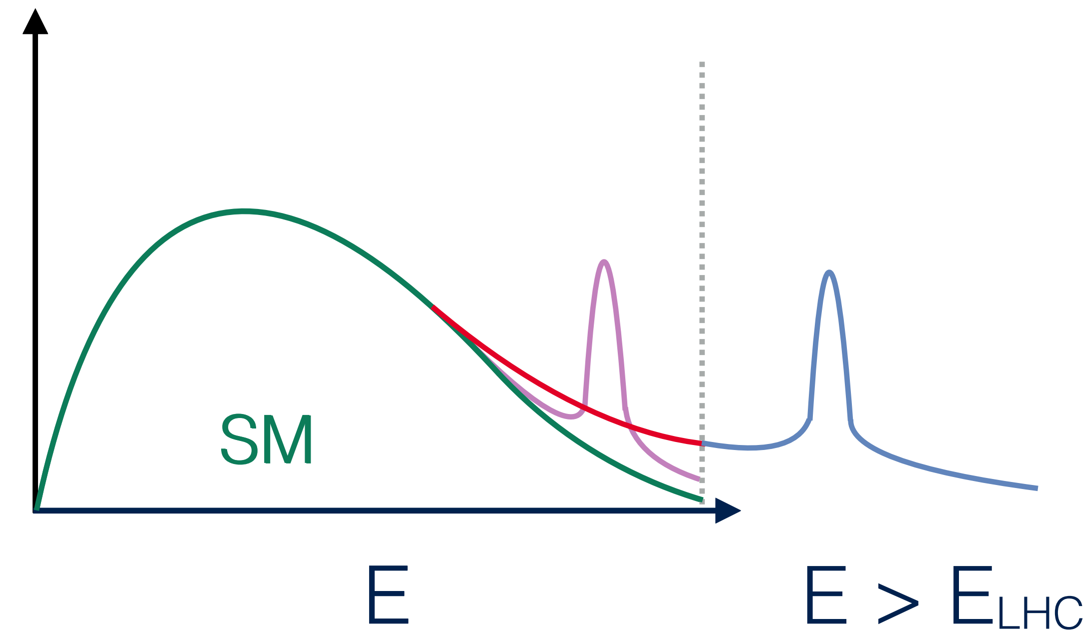
Energy & precision

Paradigm shift at the energy frontier for BSM searches

Direct (bumps)

Indirect (tails)

⇒ New physics is heavy



Heavy new physics
Precision measurements
High energy

→ **Standard Model Effective Field Theory (SMEFT)**

A QFT parameter space for BSM interactions between SM particles

SMEFT: SM v2.0

$$\mathcal{L}_{\text{eff}} = \sum_i \frac{c_i \mathcal{O}_i^D}{\Lambda^{D-4}}$$

SM is low energy effective description

- Supplemented by a tower of irrelevant operators
- Respecting low energy field content & symmetries

SU(3)$_C$ x SU(2)$_L$ x U(1)$_Y$

$$\varphi = \begin{pmatrix} G^+ \\ v + h + iG^0 \end{pmatrix} : \mathbf{2}_{\frac{1}{2}}$$

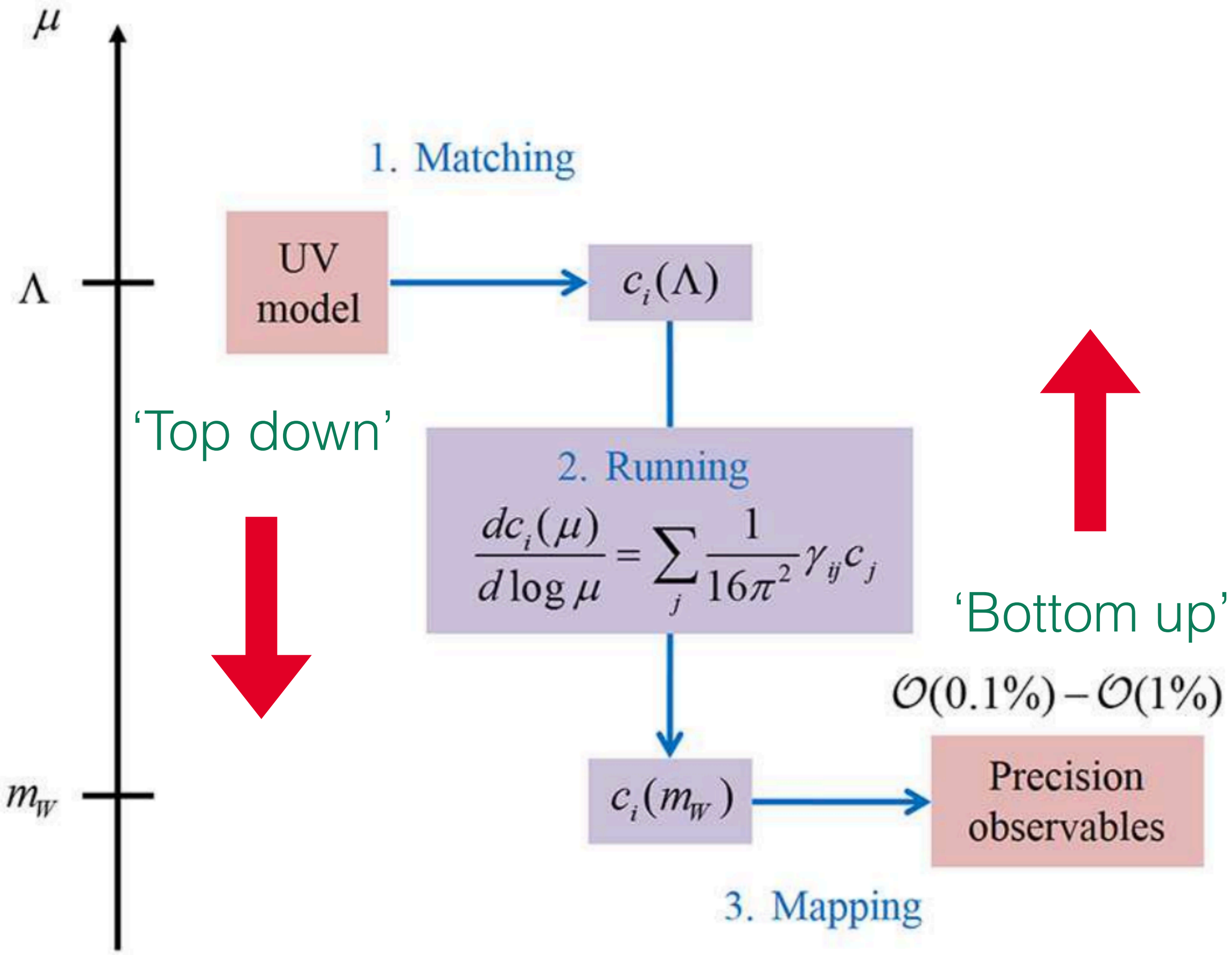
aTGC	$X^3 : \epsilon_{IJK} W^I_{\mu\nu} W^{J,\nu\rho} W^{K,\mu}_{\rho}$	$X^2 H^2 : (\varphi^\dagger \varphi)^2 G^a_{\mu\nu} G^{\mu\nu}_a$	ggh(h)
λ_h	$H^6 : (\varphi^\dagger \varphi)^3$	$H^4 D^2 : (\varphi^\dagger D^\mu \varphi)^* (\varphi^\dagger D^\mu \varphi)$	δM_Z
y_f	$\psi^2 H^3 : (\varphi^\dagger \varphi)^2 (\bar{q}_i u_j \tilde{\varphi})$	$\psi^2 X H : (\bar{q}_i \sigma^{\mu\nu} u_j \tilde{\varphi}) B_{\mu\nu}$	‘dipole’
ffV	$\psi^2 H^2 D : (\varphi^\dagger \overleftrightarrow{D}_\mu \varphi)(\bar{q}_i \gamma^\mu q_j)$	$\psi^4 : (\bar{q}_i \gamma^\mu q_j)(\bar{q}_k \gamma_\mu q_l)$	4F

More than ‘just’ a parametrisation of ignorance

- Unlike anomalous couplings
- Finite energy range ($\sim\Lambda$)
- Renormalisable QFT (order-by-order)
- Well defined matching procedure

SMEFT strategy

$$\mathcal{L}_{\rm EFT} = \mathcal{L}_{\rm SM} + \sum_i \frac{c_i^{(6)}}{\Lambda^2} O_i^{(6)} + \mathcal{O}(\Lambda^{-4})$$



Map coefficients to the data
once and for all

SMEFT is a way to test many BSM scenarios

- Economical
- Well developed

UV matching quasi-automated

- Tree-level dictionary

[de Blas et al.; JHEP 03 (2018) 109]

- Universal one loop effective action

[Henning, Lu & Murayama; JHEP 01 (2016) 023]

[Drozd et al.; JHEP 03 (2016) 180]

RGE are known

[Alonso, Jenkins, Manohar & Trott; JHEP 1310 (2013) 087, JHEP 1401 (2014) 035 JHEP 1404 (2014) 159*]*

Mature MC tools

SMEFTsim, SMEFTatNLO, dim6top,…

SMEFT is…

$$\mathcal{L}_{\rm eff} = \sum_i \frac{c_i \mathcal{O}_i^D}{\Lambda^{D-4}}$$

Model independent

- Underlying assumptions

Heavy new physics: $M > E_{exp}$

SM field content & gauge symmetries

Linear EWSB: Higgs = doublet

Systematically improvable

- Double expansion *higher dim.* $\frac{E^2}{\Lambda^2}$ & $\{g_S,\, g,\, g'\}$ *more loops*

Global

- Model independence: we don't know what operators NP will generate
- Patterns & correlations among observables are key
- Ultimate goal: complete SMEFT likelihood confronted with HEP data

EWPO, Higgs, multiboson, top, DY, flavor,…

Established part of LHC programme

SMEFT interpretation

Ingredients:

$$\Delta o_n = o_n^{\mathrm{EXP}} - o_n^{\mathrm{SM}} = \sum_i \frac{a_{n,i}^{(6)}(\mu)\, c_i^{(6)}(\mu)}{\Lambda^2} + \mathcal{O}\left(\frac{1}{\Lambda^3}\right)$$

Global nature

As many observables as possible

Identify patterns & correlations in fits

Exploit energy-growth

Sensitivity

Experiment:

Best measurements & understanding of uncertainties and correlations

Theory:

Best available predictions for observables (NLO, NNLO, N3LO,…)

Interpretation

Relies on accurate knowledge of the size & correlation among a_i

Determining $c_i^{(6)}$ requires most precise available SMEFT predictions

Status in a nutshell

Global new physics searches via high precision/energy

- Z & W-pole data: handle on the EW gauge sector
- LHC: thriving Higgs & top programmes
- Probing gauge interactions at high energy (VV, VBS, VVV, …)

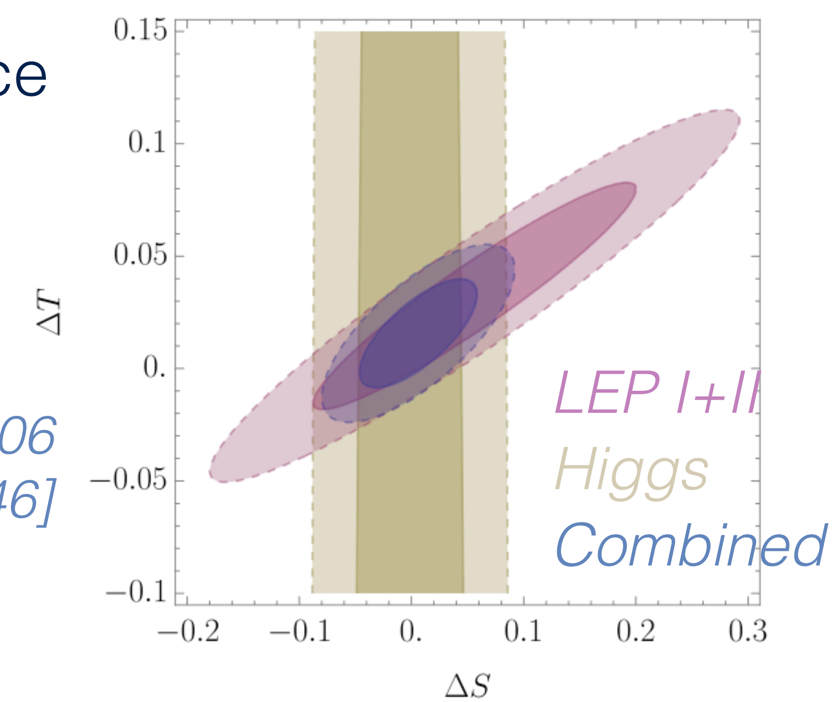
[Han & Skiba; PRD 71 (2005) 075009]
[Falkowski & Riva; JHEP 02 (2015) 039]

How much cross-talk? Where does being global matter?

We know that Higgs data greatly complements LEP

- Access unconstrained directions in parameter space
- Allows for a closed fit to flavor-universal SMEFT
- Crucial to combine EWPO, Diboson & Higgs data

[Corbett et al.; PRD 87 (2013) 015022]
[Pomarol & Riva; JHEP 01 (2014) 151]
[Ellis, Sanz & You; JHEP 03 (2015) 157]
[Biekötter Corbett & Plehn; SciPost Phys 6 (2019) 6, 064]…
[Ellis et al.; JHEP 06 (2018) 146]

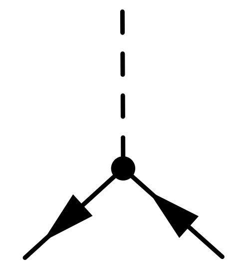


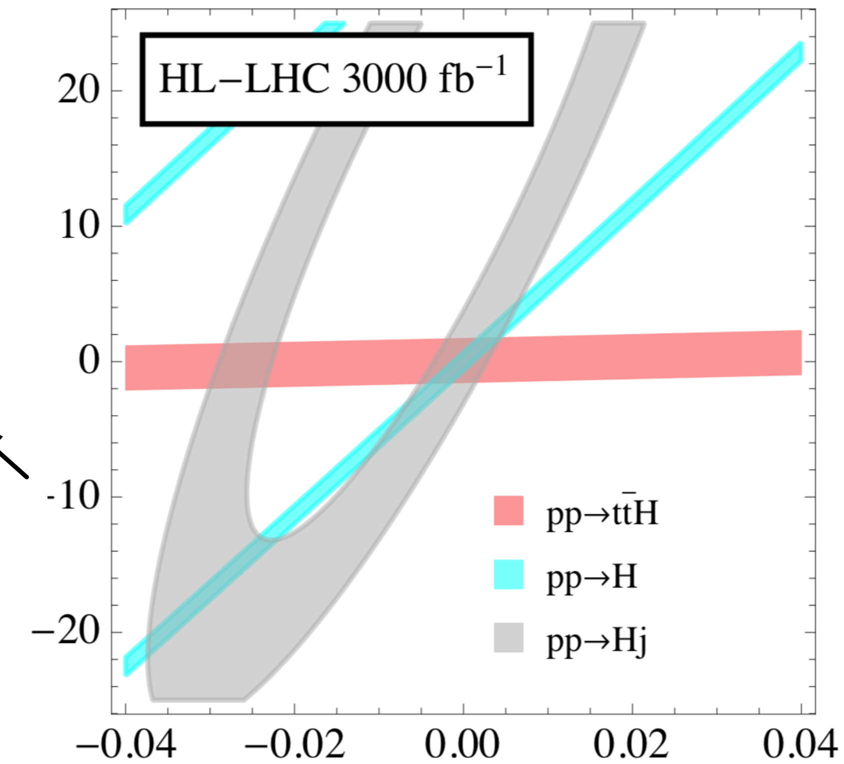
[Maltoni, Vryonidou & Zhang; JHEP 1610 (2016) 123]

Top & Higgs

Inextricably linked in the SM

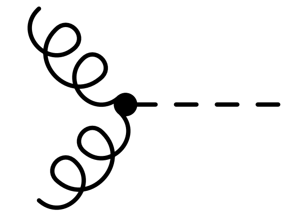
- Yukawa interaction controls ggF
- Strong BSM motivation to study tops



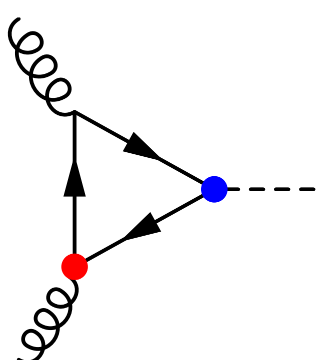


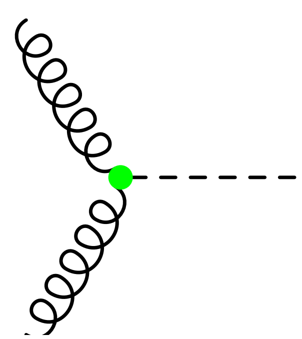
ggF is well measured now

- Does not exclude top partners, anomalous Yukawa!



C_{HG} *Point-like*
C_{tH} *Yukawa*
C_{tG} *Dipole*

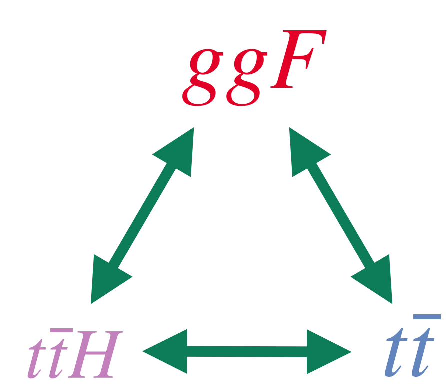




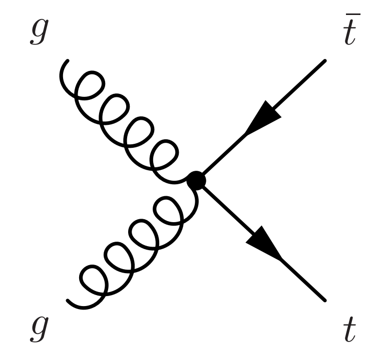
Blind direction in BSM scenarios
=
Effective coupling degeneracy

Need more data to break degeneracy

- $t\bar{t}H$ production for direct Yukawa measurement
- $t\bar{t}$ data to constrain dipole

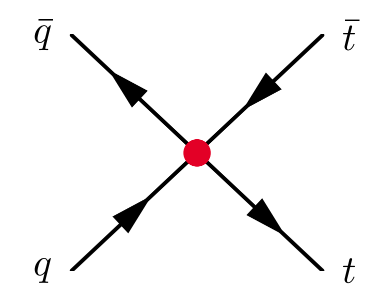


The role of top data



$t\bar{t}$ cross section measurements constrain C_{tG}

- Indirectly improve bounds on C_{HG} and C_{tH}



Several other new interactions can affect $t\bar{t}$

- Notably $q\bar{q}t\bar{t}$ operators, of which there are many (14)
- To what extent do these limit ultimate NP sensitivity in top/Higgs sector?

Can only be addressed in combined fit

- Beyond tree-level (at least for ggF)

[Degrande et al.; arXiv:2008.11743]
http://feynrules.irmp.ucl.ac.be/wiki/SMEFTatNLO

- Identify other cross-talk (non-trivial correlations)
- Crystallisation of knowledge gained after LHC Run 2
- Broaden range of applicability to UV models

The fit

Top, Higgs, Diboson and Electroweak Fit to the Standard Model Effective Field Theory

John Ellis,[a,b,c] Maeve Madigan,[d] Ken Mimasu,[a] Veronica Sanz[e,f] and Tevong You[b,d,g]

[JHEP 04 (2021) 279]

Global SMEFT interpretation of 4 categories of data

Based on [Ellis et al.; JHEP 06 (2018) 146]

14 • Electroweak Precision Observables (EWPO): Z-pole & W-mass

118 • LEP2 & LHC diboson production: differential WW, WZ, Zjj

72 • Higgs measurements: signal strengths & STXS

137 • Top data: single-top, ttbar & asymmetries, ttV, tZ, tW

Big thanks to authors of SMEFiT analysis [JHEP 04 (2019) 100] for sharing some of their top predictions

341 measurements across categories

- Chosen to be statistically independent & maximise reach
- Correlations included when publicly available (mostly are)

Linear EFT approximation:
$$\mu_X \equiv \frac{X}{X_{SM}} = 1 + \sum_i a_i^X \frac{C_i}{\Lambda^2} + \mathcal{O}\left(\frac{1}{\Lambda^4}\right)$$

Theory

[Grzadkowski et al.; JHEP 10 (2010) 085]

X^3		H^6 and H^4D^2		ψ^2H^3	
$\mathcal{O}_G$	$f^{ABC}G_\mu^{A\nu}G_\nu^{B\rho}G_\rho^{C\mu}$	$\mathcal{O}_H$	$(H^\dagger H)^3$	$\mathcal{O}_{eH}$	$(H^\dagger H)(\bar l_p e_r H)$
$\mathcal{O}_{\widetilde G}$	$f^{ABC}\widetilde G_\mu^{A\nu}G_\nu^{B\rho}G_\rho^{C\mu}$	$\mathcal{O}_{H\Box}$	$(H^\dagger H)\Box(H^\dagger H)$	$\mathcal{O}_{uH}$	$(H^\dagger H)(\bar q_p u_r \widetilde H)$
$\mathcal{O}_W$	$\varepsilon^{IJK}W_\mu^{I\nu}W_\nu^{J\rho}W_\rho^{K\mu}$	$\mathcal{O}_{HD}$	$(H^\dagger D^\mu H)^\star(H^\dagger D_\mu H)$	$\mathcal{O}_{dH}$	$(H^\dagger H)(\bar q_p d_r H)$
$\mathcal{O}_{\widetilde W}$	$\varepsilon^{IJK}\widetilde W_\mu^{I\nu}W_\nu^{J\rho}W_\rho^{K\mu}$				
X^2H^2		ψ^2XH		ψ^2H^2D	
$\mathcal{O}_{HG}$	$H^\dagger H\, G^A_{\mu\nu}G^{A\mu\nu}$	$\mathcal{O}_{eW}$	$(\bar l_p\sigma^{\mu\nu}e_r)\tau^I H W^I_{\mu\nu}$	$\mathcal{O}_{Hl}^{(1)}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar l_p\gamma^\mu l_r)$
$\mathcal{O}_{H\widetilde G}$	$H^\dagger H\, \widetilde G^A_{\mu\nu}G^{A\mu\nu}$	$\mathcal{O}_{eB}$	$(\bar l_p\sigma^{\mu\nu}e_r)H B_{\mu\nu}$	$\mathcal{O}_{Hl}^{(3)}$	$(H^\dagger i\overleftrightarrow{D}^I_\mu H)(\bar l_p\tau^I\gamma^\mu l_r)$
$\mathcal{O}_{HW}$	$H^\dagger H\, W^I_{\mu\nu}W^{I\mu\nu}$	$\mathcal{O}_{uG}$	$(\bar q_p\sigma^{\mu\nu}T^A u_r)\widetilde H\, G^A_{\mu\nu}$	$\mathcal{O}_{He}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar e_p\gamma^\mu e_r)$
$\mathcal{O}_{H\widetilde W}$	$H^\dagger H\, \widetilde W^I_{\mu\nu}W^{I\mu\nu}$	$\mathcal{O}_{uW}$	$(\bar q_p\sigma^{\mu\nu}u_r)\tau^I\widetilde H\, W^I_{\mu\nu}$	$\mathcal{O}_{Hq}^{(1)}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar q_p\gamma^\mu q_r)$
$\mathcal{O}_{HB}$	$H^\dagger H\, B_{\mu\nu}B^{\mu\nu}$	$\mathcal{O}_{uB}$	$(\bar q_p\sigma^{\mu\nu}u_r)\widetilde H\, B_{\mu\nu}$	$\mathcal{O}_{Hq}^{(3)}$	$(H^\dagger i\overleftrightarrow{D}^I_\mu H)(\bar q_p\tau^I\gamma^\mu q_r)$
$\mathcal{O}_{H\widetilde B}$	$H^\dagger H\, \widetilde B_{\mu\nu}B^{\mu\nu}$	$\mathcal{O}_{dG}$	$(\bar q_p\sigma^{\mu\nu}T^A d_r)H\, G^A_{\mu\nu}$	$\mathcal{O}_{Hu}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar u_p\gamma^\mu u_r)$
$\mathcal{O}_{HWB}$	$H^\dagger\tau^I H\, W^I_{\mu\nu}B^{\mu\nu}$	$\mathcal{O}_{dW}$	$(\bar q_p\sigma^{\mu\nu}d_r)\tau^I H\, W^I_{\mu\nu}$	$\mathcal{O}_{Hd}$	$(H^\dagger i\overleftrightarrow{D}_\mu H)(\bar d_p\gamma^\mu d_r)$
$\mathcal{O}_{H\widetilde WB}$	$H^\dagger\tau^I H\, \widetilde W^I_{\mu\nu}B^{\mu\nu}$	$\mathcal{O}_{dB}$	$(\bar q_p\sigma^{\mu\nu}d_r)H\, B_{\mu\nu}$	$\mathcal{O}_{Hud}$	$i(\widetilde H^\dagger D_\mu H)(\bar u_p\gamma^\mu d_r)$

$(\bar LL)(\bar LL)$		$(\bar RR)(\bar RR)$		$(\bar LL)(\bar RR)$	
$\mathcal{O}_{ll}$	$(\bar l_p\gamma_\mu l_r)(\bar l_s\gamma^\mu l_t)$	$\mathcal{O}_{ee}$	$(\bar e_p\gamma_\mu e_r)(\bar e_s\gamma^\mu e_t)$	$\mathcal{O}_{le}$	$(\bar l_p\gamma_\mu l_r)(\bar e_s\gamma^\mu e_t)$
$\mathcal{O}_{qq}^{(1)}$	$(\bar q_p\gamma_\mu q_r)(\bar q_s\gamma^\mu q_t)$	$\mathcal{O}_{uu}$	$(\bar u_p\gamma_\mu u_r)(\bar u_s\gamma^\mu u_t)$	$\mathcal{O}_{lu}$	$(\bar l_p\gamma_\mu l_r)(\bar u_s\gamma^\mu u_t)$
$\mathcal{O}_{qq}^{(3)}$	$(\bar q_p\gamma_\mu\tau^I q_r)(\bar q_s\gamma^\mu\tau^I q_t)$	$\mathcal{O}_{dd}$	$(\bar d_p\gamma_\mu d_r)(\bar d_s\gamma^\mu d_t)$	$\mathcal{O}_{ld}$	$(\bar l_p\gamma_\mu l_r)(\bar d_s\gamma^\mu d_t)$
$\mathcal{O}_{lq}^{(1)}$	$(\bar l_p\gamma_\mu l_r)(\bar q_s\gamma^\mu q_t)$	$\mathcal{O}_{eu}$	$(\bar e_p\gamma_\mu e_r)(\bar u_s\gamma^\mu u_t)$	$\mathcal{O}_{qe}$	$(\bar q_p\gamma_\mu q_r)(\bar e_s\gamma^\mu e_t)$
$\mathcal{O}_{lq}^{(3)}$	$(\bar l_p\gamma_\mu\tau^I l_r)(\bar q_s\gamma^\mu\tau^I q_t)$	$\mathcal{O}_{ed}$	$(\bar e_p\gamma_\mu e_r)(\bar d_s\gamma^\mu d_t)$	$\mathcal{O}_{qu}^{(1)}$	$(\bar q_p\gamma_\mu q_r)(\bar u_s\gamma^\mu u_t)$
		$\mathcal{O}_{ud}^{(1)}$	$(\bar u_p\gamma_\mu u_r)(\bar d_s\gamma^\mu d_t)$	$\mathcal{O}_{qu}^{(8)}$	$(\bar q_p\gamma_\mu T^A q_r)(\bar u_s\gamma^\mu T^A u_t)$
		$\mathcal{O}_{ud}^{(8)}$	$(\bar u_p\gamma_\mu T^A u_r)(\bar d_s\gamma^\mu T^A d_t)$	$\mathcal{O}_{qd}^{(1)}$	$(\bar q_p\gamma_\mu q_r)(\bar d_s\gamma^\mu d_t)$
				$\mathcal{O}_{qd}^{(8)}$	$(\bar q_p\gamma_\mu T^A q_r)(\bar d_s\gamma^\mu T^A d_t)$
$(\bar LR)(\bar RL)$ and $(\bar LR)(\bar LR)$		B-violating			
$\mathcal{O}_{ledq}$	$(\bar l_p^j e_r)(\bar d_s q_t^j)$	$\mathcal{O}_{duq}$	$\varepsilon^{\alpha\beta\gamma}\varepsilon_{jk}\left[(d_p^\alpha)^T C u_r^\beta\right]\left[(q_s^{\gamma j})^T C l_t^k\right]$		
$\mathcal{O}_{quqd}^{(1)}$	$(\bar q_p^j u_r)\varepsilon_{jk}(\bar q_s^k d_t)$	$\mathcal{O}_{qqu}$	$\varepsilon^{\alpha\beta\gamma}\varepsilon_{jk}\left[(q_p^{\alpha j})^T C q_r^{\beta k}\right]\left[(u_s^\gamma)^T C e_t\right]$		
$\mathcal{O}_{quqd}^{(8)}$	$(\bar q_p^j T^A u_r)\varepsilon_{jk}(\bar q_s^k T^A d_t)$	$\mathcal{O}_{qqq}$	$\varepsilon^{\alpha\beta\gamma}\varepsilon_{jn}\varepsilon_{km}\left[(q_p^{\alpha j})^T C q_r^{\beta k}\right]\left[(q_s^{\gamma m})^T C l_t^n\right]$		
$\mathcal{O}_{lequ}^{(1)}$	$(\bar l_p^j e_r)\varepsilon_{jk}(\bar q_s^k u_t)$	$\mathcal{O}_{duu}$	$\varepsilon^{\alpha\beta\gamma}\left[(d_p^\alpha)^T C u_r^\beta\right]\left[(u_s^\gamma)^T C e_t\right]$		
$\mathcal{O}_{lequ}^{(3)}$	$(\bar l_p^j\sigma_{\mu\nu}e_r)\varepsilon_{jk}(\bar q_s^k\sigma^{\mu\nu}u_t)$				

Warsaw basis with CP & B conservation

- Full 'bosonic' sector: Higgs, triple-gauge & gauge-Higgs
- Scenario 1: Flavor-**universal** degrees of freedom

$U(3)_L \times U(3)_e \times U(3)_Q \times U(3)_u \times U(3)_d$ + Yukawas: $\mathcal{O}_{tH}, \mathcal{O}_{bH}, \mathcal{O}_{\tau H}, \mathcal{O}_{\mu H}$

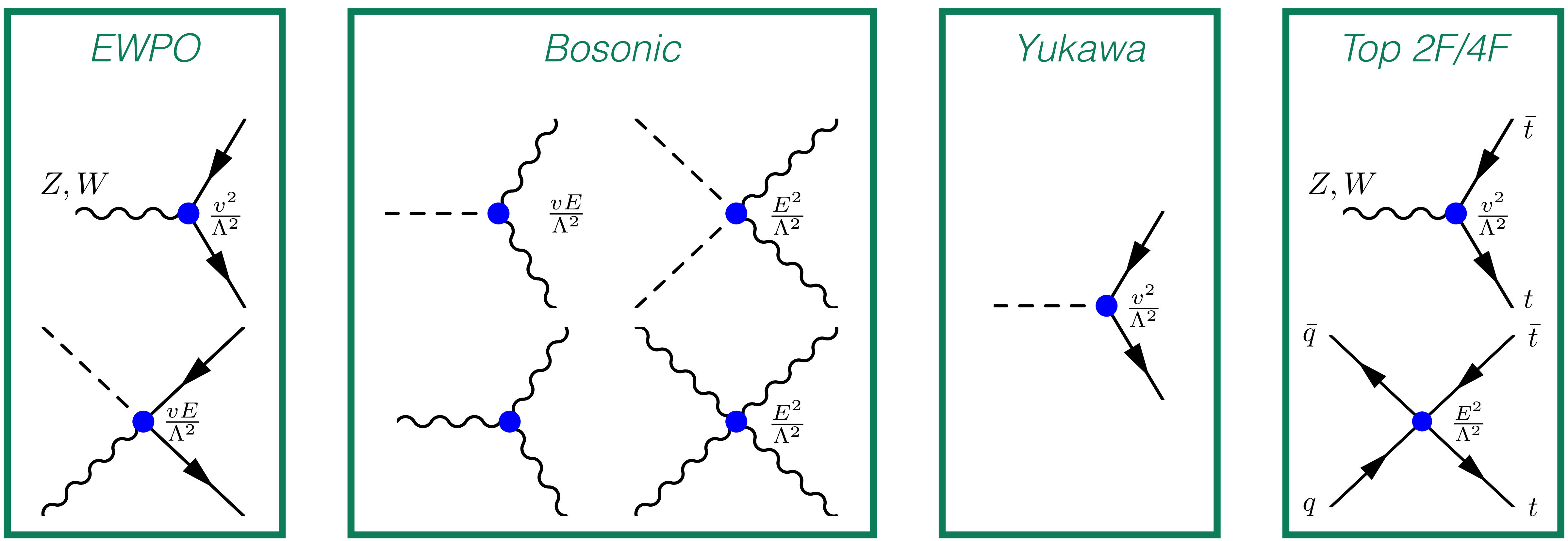
- Scenario 2: **top**-centric flavor symmetry

$U(3)_L \times U(3)_e \times U(2)_Q \times U(2)_u \times U(3)_d$

cf. Minimal flavor violation
[Buras et al.; PLB 500 (2001) 161]
[D'Ambrosio et al.; NPB 645 (2002) 155]
[Aguilar-Saavedra et al.; arXiv:1802.07237]

Degrees of freedom

EWPO:	$\mathcal{O}_{HWB}\,,\ \mathcal{O}_{HD}\,,\ \mathcal{O}_{ll}\,,\ \mathcal{O}_{Hl}^{(3)}\,,\ \mathcal{O}_{Hl}^{(1)}\,,\ \mathcal{O}_{He}\,,\ \mathcal{O}_{Hq}^{(3)}\,,\ \mathcal{O}_{Hq}^{(1)}\,,\ \mathcal{O}_{Hd}\,,\ \mathcal{O}_{Hu}\,,$	
Bosonic:	$\mathcal{O}_{H\Box}\,,\ \mathcal{O}_{HG}\,,\ \mathcal{O}_{HW}\,,\ \mathcal{O}_{HB}\,,\ \mathcal{O}_{W}\,,\ \mathcal{O}_{G}\,,$	
Yukawa:	$\mathcal{O}_{\tau H}\,,\ \mathcal{O}_{\mu H}\,,\ \mathcal{O}_{bH}\,,\ \mathcal{O}_{tH}\,,$	20
Top 2F:	$\mathcal{O}_{HQ}^{(3)}\,,\ \mathcal{O}_{HQ}^{(1)}\,,\ \mathcal{O}_{Ht}\,,\ \mathcal{O}_{tG}\,,\ \mathcal{O}_{tW}\,,\ \mathcal{O}_{tB}\,,$	
Top 4F:	$\mathcal{O}_{Qq}^{3,1}\,,\ \mathcal{O}_{Qq}^{3,8}\,,\ \mathcal{O}_{Qq}^{1,8}\,,\ \mathcal{O}_{Qu}^{8}\,,\ \mathcal{O}_{Qd}^{8}\,,\ \mathcal{O}_{tQ}^{8}\,,\ \mathcal{O}_{tu}^{8}\,,\ \mathcal{O}_{td}^{8}\,.$	+14



Degrees of freedom

EWPO:	$\mathcal{O}_{HWB}\,,\mathcal{O}_{HD}\,,\mathcal{O}_{ll}\,,\mathcal{O}_{Hl}^{(3)}\,,\mathcal{O}_{Hl}^{(1)}\,,\mathcal{O}_{He}\,,\mathcal{O}_{Hq}^{(3)}\,,\mathcal{O}_{Hq}^{(1)}\,,\mathcal{O}_{Hd}\,,\mathcal{O}_{Hu}\,,$	
Bosonic:	$\mathcal{O}_{H\Box}\,,\mathcal{O}_{HG}\,,\mathcal{O}_{HW}\,,\mathcal{O}_{HB}\,,\mathcal{O}_{W}\,,\mathcal{O}_{G}\,,$	
Yukawa:	$\mathcal{O}_{\tau H}\,,\mathcal{O}_{\mu H}\,,\mathcal{O}_{bH}\,,\mathcal{O}_{tH}\,,$	20
Top 2F:	$\mathcal{O}_{HQ}^{(3)}\,,\mathcal{O}_{HQ}^{(1)}\,,\mathcal{O}_{Ht}\,,\mathcal{O}_{tG}\,,\mathcal{O}_{tW}\,,\mathcal{O}_{tB}\,,$	
Top 4F:	$\mathcal{O}_{Qq}^{3,1}\,,\mathcal{O}_{Qq}^{3,8}\,,\mathcal{O}_{Qq}^{1,8}\,,\mathcal{O}_{Qu}^{8}\,,\mathcal{O}_{Qd}^{8}\,,\mathcal{O}_{tQ}^{8}\,,\mathcal{O}_{tu}^{8}\,,\mathcal{O}_{td}^{8}\,.$	+14

In total: 20(34) d.o.f. for the two flavor scenarios

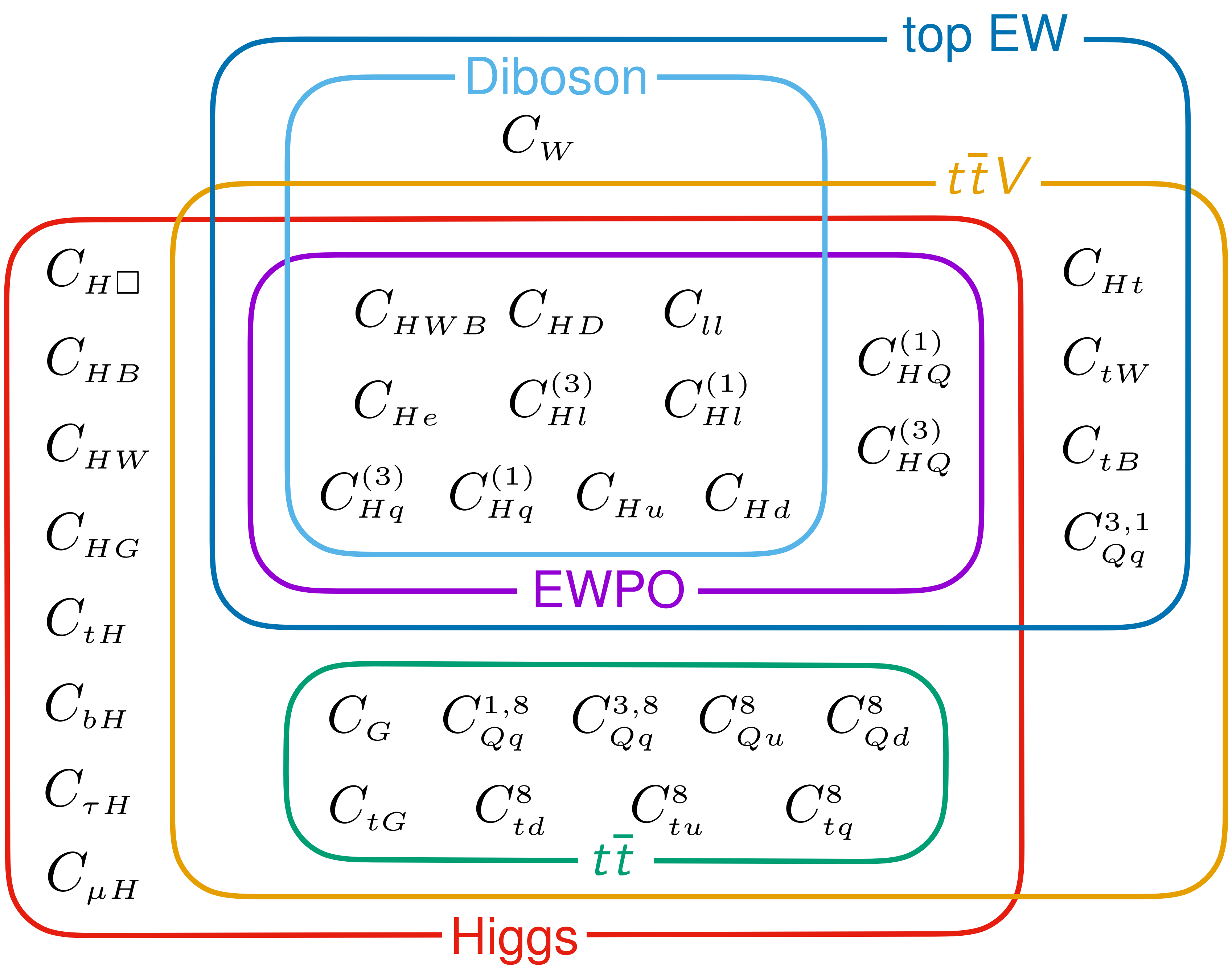
Dim6top conventions: [Aguilar-Saavedra et al.; arXiv:1802.07237]

Dictated by flavor symmetry & sensitivity of dataset

Linear EFT fit: precludes sensitivity to some ops

- Those that cannot interfere due to helicity/symmetries
- e.g. neutral colour-singlet top 4F operators: $(\bar{q}\gamma^\mu q)(\bar{t}\gamma^\mu t)$ (x 6)
- Four-heavy quark operators in 4top & ttbb (quadratic dominated)

Interplay



top EW
Diboson
C_W
$t\bar{t}V$
$C_{H\Box}$
C_{HB}
C_{HW}
C_{HG}
C_{tH}
C_{bH}
$C_{\tau H}$
$C_{\mu H}$
C_{HWB} C_{HD} C_{ll}
C_{He} $C^{(3)}_{Hl}$ $C^{(1)}_{Hl}$
$C^{(3)}_{Hq}$ $C^{(1)}_{Hq}$ C_{Hu} C_{Hd}
$C^{(1)}_{HQ}$
$C^{(3)}_{HQ}$
C_{Ht}
C_{tW}
C_{tB}
$C^{3,1}_{Qq}$
EWPO
C_G $C^{1,8}_{Qq}$ $C^{3,8}_{Qq}$ C^{8}_{Qu} C^{8}_{Qd}
C_{tG} C^{8}_{td} C^{8}_{tu} C^{8}_{tq}
$t\bar{t}$
Higgs

Technical details

$$\mu_X \equiv \frac{X}{X_{SM}} = 1 + \sum_i a_i^X \frac{C_i}{\Lambda^2} + \mathcal{O}\left(\frac{1}{\Lambda^4}\right)$$

Exp. data: **`HEPdata`**, **`WebPlotDigitizer`**, …

- Construct ‘signal strength’, w.r.t. SM prediction from exp. paper
- Otherwise computed with `MG5`, `fastnlo`, directly from theory papers
- Combine all sources of uncertainty in quadrature (stat., syst., th.)

Theory predictions: **`MG5 (SMEFTsim & SMEFTatNLO)`**

- LO, parton-level, linear dependence in (α, G_F, M_Z) scheme
- Tree-level + 1-loop gluon fusion Higgs production
- a_i: Effects from production, decays, total width
- No theory error from EFT, assume SM error dominant

The code

`fitmaker` *https://gitlab.com/kenmimasu/fitrepo*

public-friendly version w/ example notebooks in progress

Main analysis: linearised least-squares fit

$$\chi^2(C_i) = (\vec{y} - \vec{\mu}(C_i))^T\, \mathbf{V}^{-1}\, (\vec{y} - \vec{\mu}(C_i)) \qquad \mu_\alpha(C_i) = \mu_\alpha^{\rm SM} + \mathbf{H}_{\alpha i} C_i$$

Best fit
$$\hat{\vec{C}} = \left(\mathbf{H}^T \mathbf{V}^{-1} \mathbf{H}\right)^{-1} \mathbf{H}^T \mathbf{V}^{-1} (\vec{y} - \vec{\mu}^{\rm SM}) \equiv \mathbf{F}^{-1} \vec{\omega}$$

$$\mathbf{F} \equiv \mathbf{H}^T \mathbf{V}^{-1} \mathbf{H} \quad , \quad \vec{\omega} \equiv \mathbf{H}^T \mathbf{V}^{-1} (\vec{y} - \vec{\mu}^{\rm SM}) \,,$$

Fisher information

$\mathbf{F}^{-1} \equiv \mathbf{U}$ Covariance matrix of least-squares estimators

$(\chi^2_{SM}, \hat{\vec{C}}, \mathbf{U})$ fully characterise likelihood

- Individual, profiled/marginalised bounds & correlations
- Principal component analysis (eigensystem of F)

Implemented as part of the `fitmaker` framework

Also nested sampling routine for general likelihoods

The code

fitmaker *https://gitlab.com/kenmimasu/fitrepo*

public-friendly version w/ example notebooks in progress

Database of input measurements encoded in `.json` format

- Values, errors, metadata,…

Python-class based definition of theoretical models

- Predictions for observables can be hard-coded
- …or read-in from a `.json` file

$$\mu^{ggF}_{H\to 4\ell} = 0.98^{+0.12}_{-0.11}$$


```
{
  "observable_name": "mu_ggF_H_ZZ_13",
  "measurement_name": "mu_ggF_H_ZZ_CMS_Run2",
  "CDS": "http://cds.cern.ch/record/2706103",
  "reportnumber": "CMS-PAS-HIG-19-005",
  "DOI": "",
  "date": "2020/01/10",
  "experiment": "CERN LHC Run 2, CMS",
  "description": "Higgs boson signal strength for
  "value": 0.98,
  "uncertainty": {
    "tot": [ 0.12, 0.11 ]
  },
  "uncertainty_sigma": 1,
  "th_flat": true
}
```


$$\mu^{ggF} = 1 + 35.8C_{HG} - 0.122C_{tH}$$
$$-0.959C_{tG} - 0.121C_{H\Box} + \ldots$$


```
{
  "observable": "ggF0j",
  "params": [ "CHG", "CuH", "CuG", "CHbox" ],
  "constant": 1.0,
  "linear": [ 35.8, -0.122, 0.959, -0.121 ],
  "quadratic": [
  [ 321.0, -1.095, 8.45, -1.085 ],
 [ -1.095, 0.00371, -0.02925, 0.003695 ],
 [ 8.45, -0.02925, 0.23, -0.0291 ],
 [ -1.085, 0.003695, -0.0291, 0.00367 ]
  ],
  "lambda_gen": 1000.0
}
```

SMEFT@NLO

Loops & SMEFT: active field in recent years

- Non-universal K-factors in EFT space ⇔ new information at NLO
- Loop-induced sensitivity (e.g. $gg \to H$)
- Control theoretical uncertainties
- Experimental interest in higher precision for SMEFT analyses/interpretations

Challenge: many processes x many operators

- LO ⇒ NLO = more cross-talk/operators/complexity
- Automated tools for fixed-order/NLO+PS are essential to the LHC programme

Solution: SMEFT@NLO

- UFO model for `MadGraph5_aMC@NLO`
- Process-independent implementation: SMEFT in top-specific flavor limit

http://feynrules.irmp.ucl.ac.be/wiki/SMEFTatNLO

Standard Model Effective Theory at One-Loop in QCD

Céline Degrande, Gauthier Durieux, Fabio Maltoni, Ken Mimasu, Eleni Vryonidou & Cen Zhang, arXiv:2008.11743

The implementation is based on the Warsaw basis of dimension-six SMEFT operators, after canonical normalization. Electroweak input parameters are taken to be G_F, M_Z, M_W. The CKM matrix is approximated as a unit matrix, and a $U(2)_q \times U(2)_u \times U(3)_d \times (U(1)_l \times U(1)_e)^3$ flavor symmetry is enforced. It forbids all fermion masses and Yukawa couplings except that only of the top quark. The model therefore implements the five-flavor scheme for PDFs.

A new coupling order, `NP=2`, is assigned to SMEFT interactions. The cutoff scale `Lambda` takes a default value of 1 TeV^{-2} and can be modified along with the Wilson coefficients in the `param_card`. Operators definitions, normalisations and coefficient names in the UFO model are specified in definitions.pdf. The notations and normalizations of top-quark operator coefficients comply with the LHC TOP WG standards of 1802.07237. Note however that the flavor symmetry enforced here is slightly more restrictive than the baseline assumption there (see the dim6top page for more information). This model has been validated at tree level against the `dim6top` implementation (see 1906.12310 and the comparison details).

Current implementation

UFO model: SMEFTatNLO_v1.0.tar.gz

The current implementation imposes CP conservation. In the quark sector, it focuses primarily on top-quark interactions. The light-quark current operator, qqHDH, uuHDH, ddHDH, with coefficients `cpq3i`, `cpqMi`, `cpu`, `cpd` are however included. The triple-gluon operator, with coefficient `cG`, is currently not available (see the loop-capable GGG implementation). Vertices including more than four scalars or four leptons are not included. Scalar and tensor `QQl1` operators, with coefficients `ctlS3`, `ctlT3`, and `cblS3`, break our flavor symmetry assumption and are not available for one-loop computations. Top-quark flavor-changing interactions, not compatible with the imposed flavor symmetry, are not included (see the loop-capable TopFCNC implementation).

Unlike prescribed by the LHC TOP WG, the top quark chromomagnetic-dipole operator coefficient `ctG` is normalized with a factor of the strong coupling, g_S. This normalization factor temporarily ensures compatibility with the 2.X.X series of MadGraph5_aMC@NLO but may be dropped in the future. As with every other appearance of this coupling in MadGraph5_aMC@NLO, its value is renormalisation-group evolved to the QCD renormalization scale (set in the run_card).

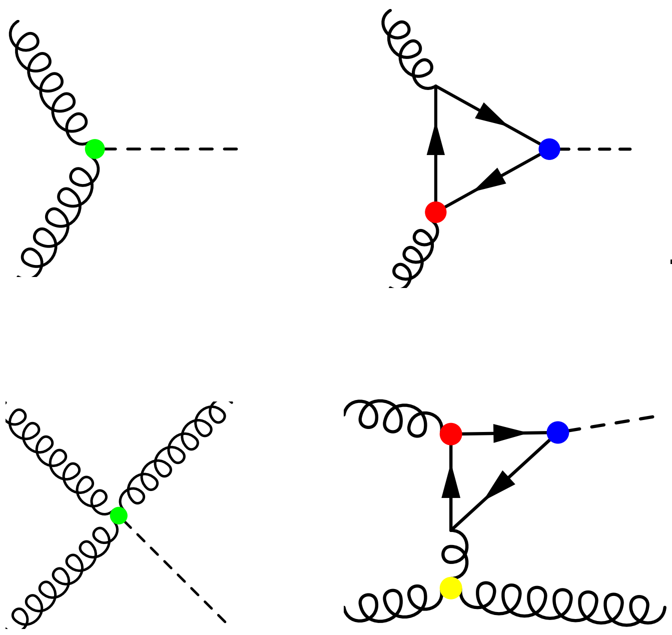
```
MG5_aMC>import model SMEFTatNLO
MG5_aMC>generate p p > t t~ NP=2 [QCD]
MG5_aMC>output
MG5_aMC>launch
```

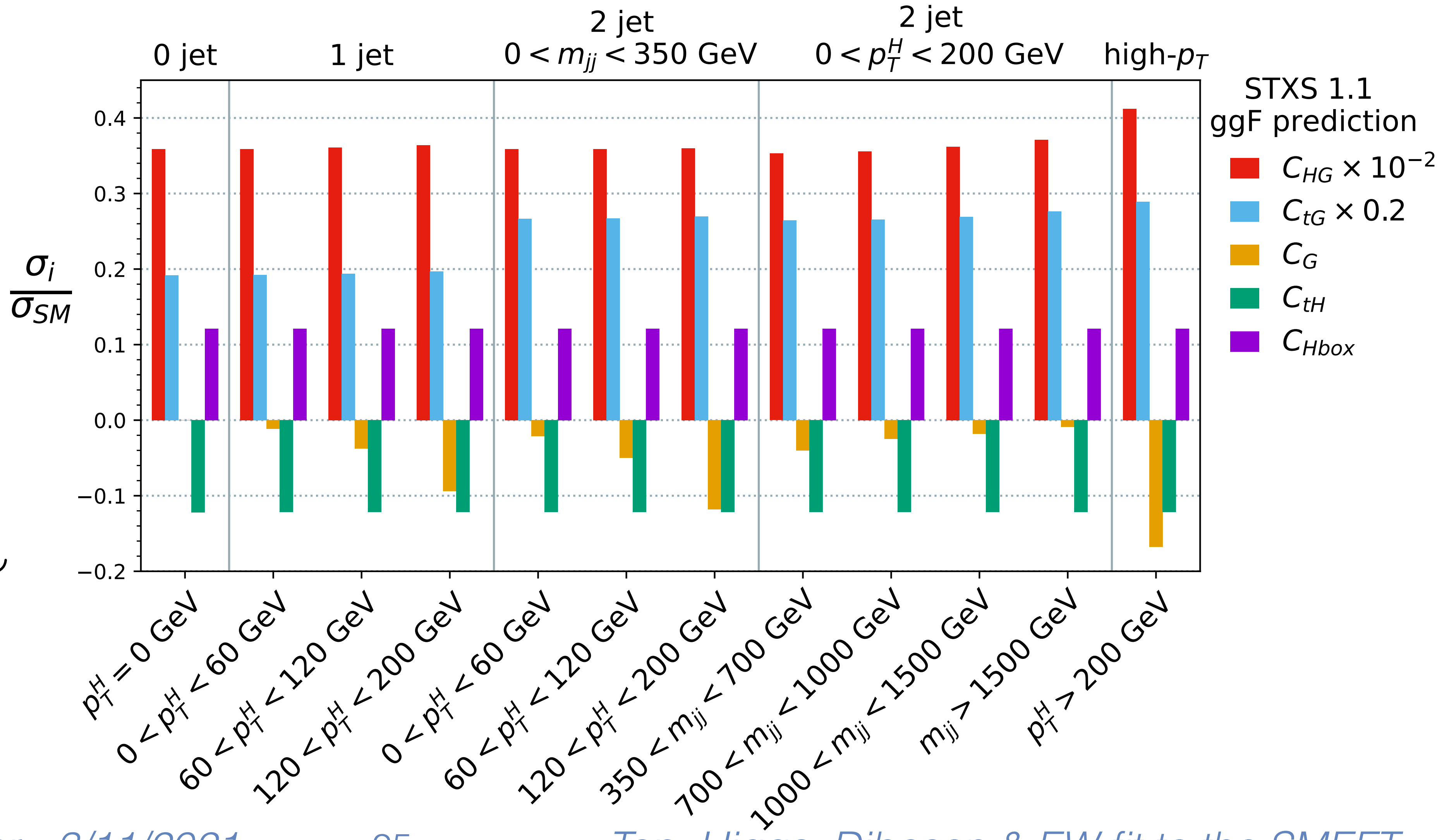
See also: [ATLAS-CONF-2020-053]

SMEFT@NLO in STXS

Gluon fusion Simplified Template Cross Sections bins

- LO in the SM is one-loop
- Tree-EFT x loop-SM + loop-EFT x loop-SM interference terms
- Heavy top limit is OK for 0-jet, breaks down at high-p_T





Results roadmap

1. Flavor universal: EWPO + diboson + Higgs

2. Top only: EWPO + top

Interlude: Top-Higgs interplay

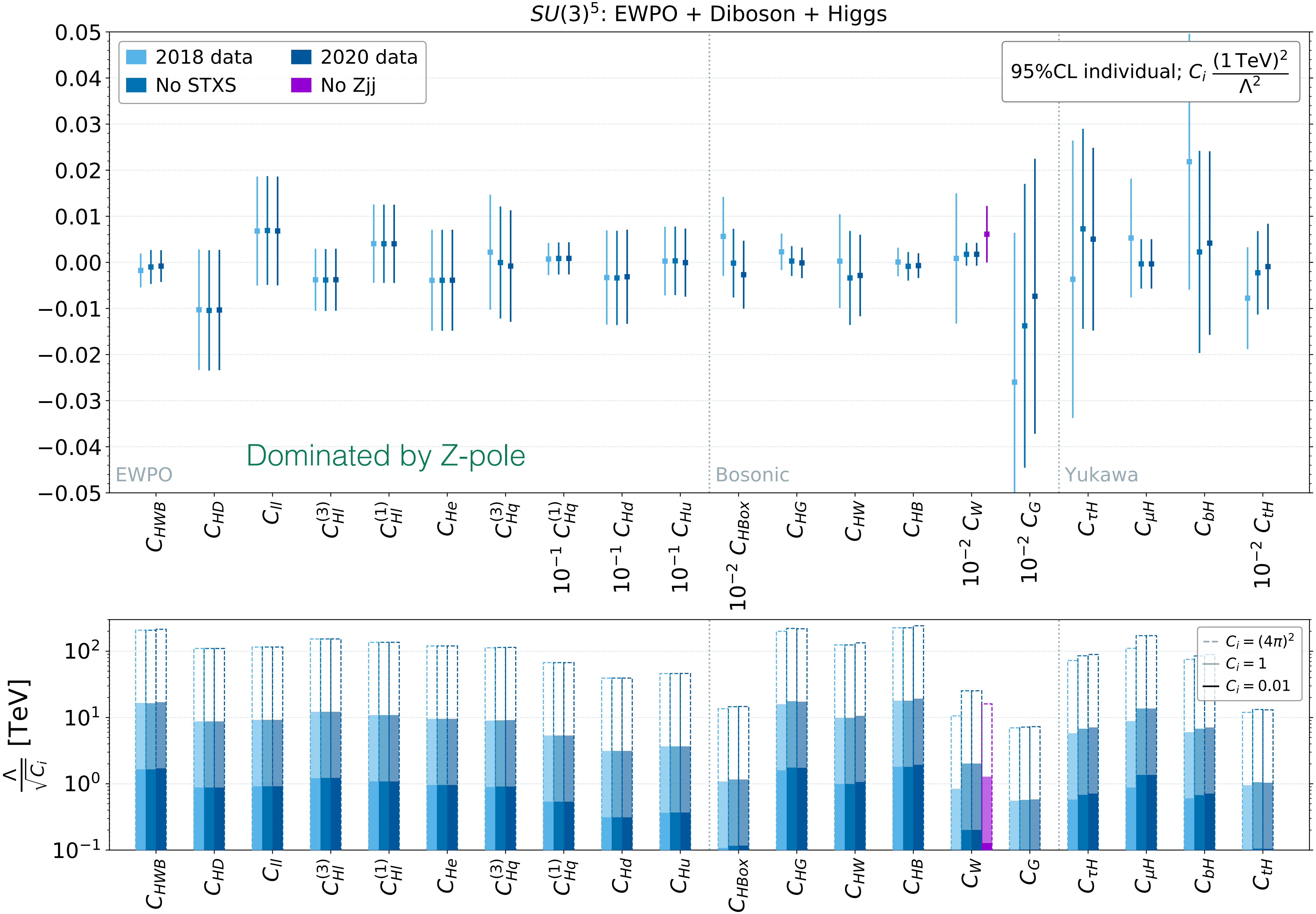
3. Top-specific : EWPO + diboson + Higgs + top

$U(3)^5$ ↓ $U(2)^2 \times U(3)^3$

EWPO:	$\mathcal{O}_{HWB}\,,\ \mathcal{O}_{HD}\,,\ \mathcal{O}_{ll}\,,\ \mathcal{O}_{Hl}^{(3)}\,,\ \mathcal{O}_{Hl}^{(1)}\,,\ \mathcal{O}_{He}\,,\ \mathcal{O}_{Hq}^{(3)}\,,\ \mathcal{O}_{Hq}^{(1)}\,,\ \mathcal{O}_{Hd}\,,\ \mathcal{O}_{Hu}\,,$
Bosonic:	$\mathcal{O}_{H\Box}\,,\ \mathcal{O}_{HG}\,,\ \mathcal{O}_{HW}\,,\ \mathcal{O}_{HB}\,,\ \mathcal{O}_{W}\,,\ \mathcal{O}_{G}\,,$
Yukawa:	$\mathcal{O}_{\tau H}\,,\ \mathcal{O}_{\mu H}\,,\ \mathcal{O}_{bH}\,,\ \mathcal{O}_{tH}\,,$
Top 2F:	$\mathcal{O}_{HQ}^{(3)}\,,\ \mathcal{O}_{HQ}^{(1)}\,,\ \mathcal{O}_{Ht}\,,\ \mathcal{O}_{tG}\,,\ \mathcal{O}_{tW}\,,\ \mathcal{O}_{tB}\,,$ $+\mathcal{O}_G$
Top 4F:	$\mathcal{O}_{Qq}^{3,1}\,,\ \mathcal{O}_{Qq}^{3,8}\,,\ \mathcal{O}_{Qq}^{1,8}\,,\ \mathcal{O}_{Qu}^{8}\,,\ \mathcal{O}_{Qd}^{8}\,,\ \mathcal{O}_{tQ}^{8}\,,\ \mathcal{O}_{tu}^{8}\,,\ \mathcal{O}_{td}^{8}\,.$

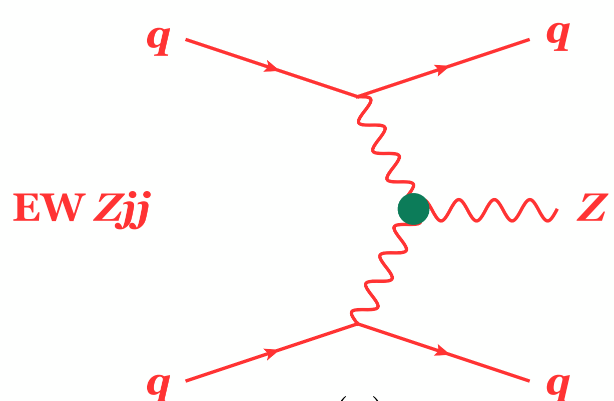
Individual limits: U(3)^5

2018 data: *[Ellis et al.; JHEP 06 (2018) 146]*



Zjj for triple gauge coupling

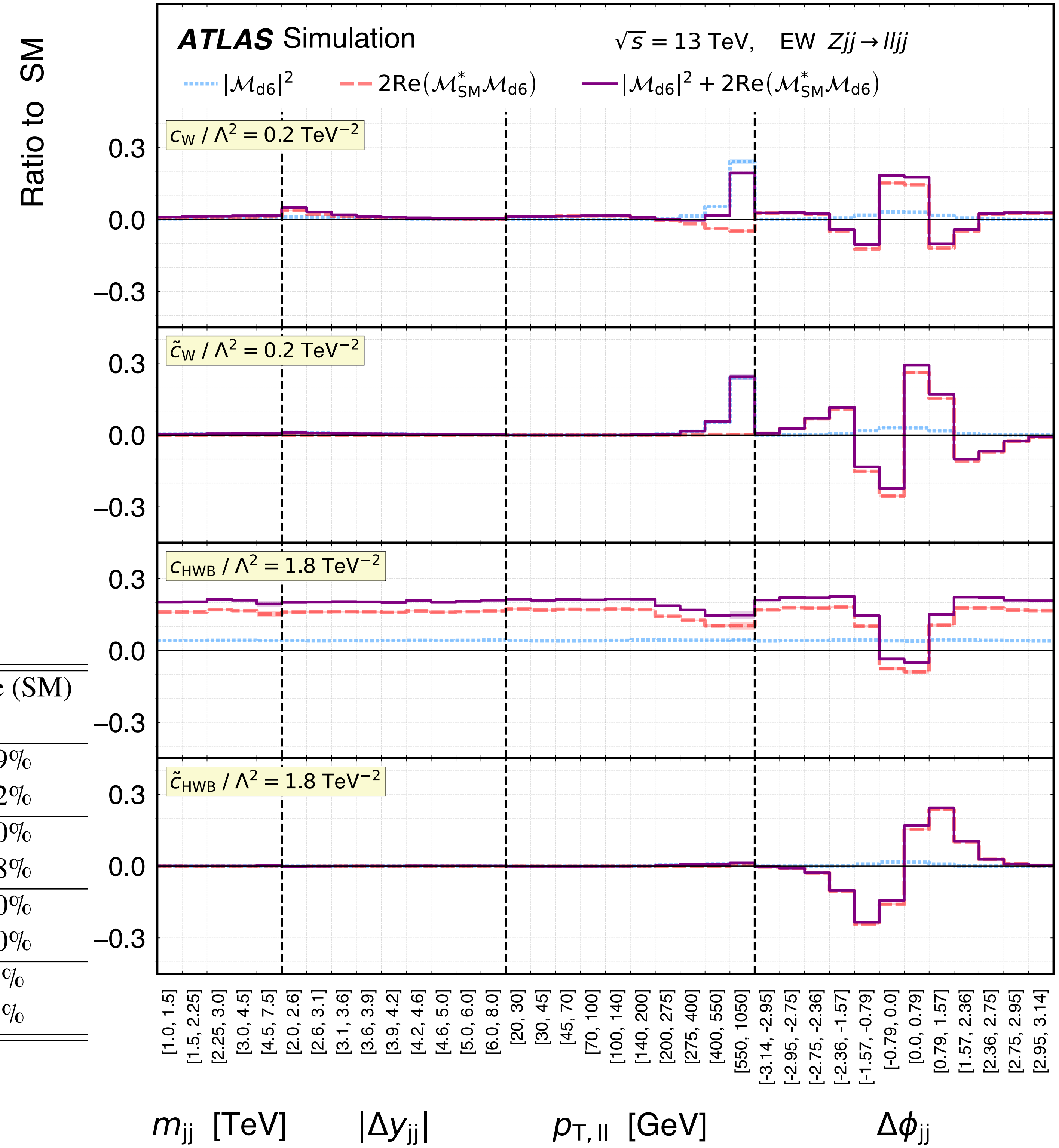
[ATLAS; CERN-EP-2020-045]



(a)

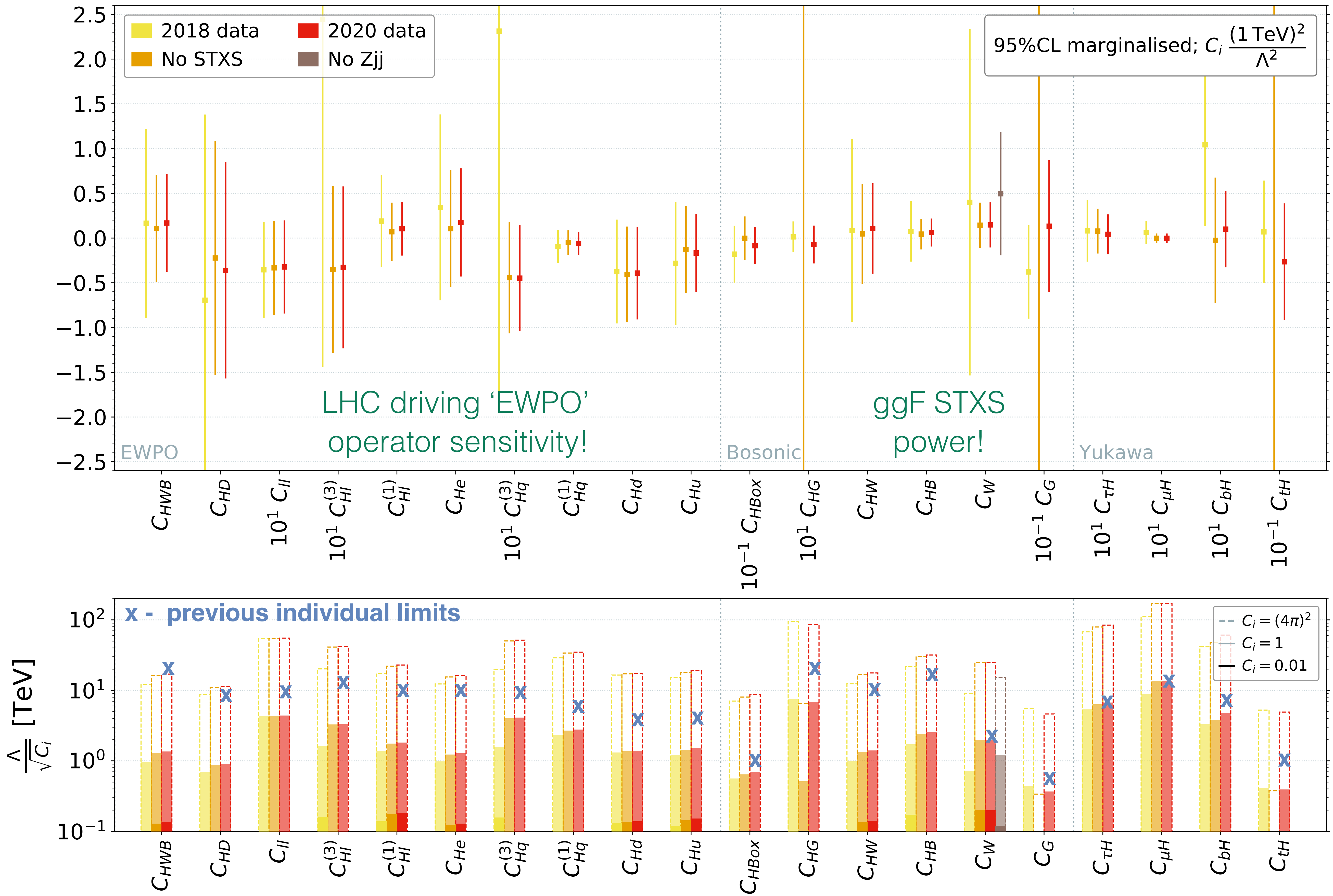
$\Delta\phi_{jj}$ distribution sensitive to linear C_W contributions

Wilson coefficient	Includes $\lvert\mathcal{M}_{d6}\rvert^2$	95% confidence interval [TeV^{-2}] Expected	Observed	p-value (SM)
c_W/Λ^2	no	[−0.30, 0.30]	[−0.19, 0.41]	45.9%
	yes	[−0.31, 0.29]	[−0.19, 0.41]	43.2%
$\tilde{c}_W/\Lambda^2$	no	[−0.12, 0.12]	[−0.11, 0.14]	82.0%
	yes	[−0.12, 0.12]	[−0.11, 0.14]	81.8%
c_{HWB}/Λ^2	no	[−2.45, 2.45]	[−3.78, 1.13]	29.0%
	yes	[−3.11, 2.10]	[−6.31, 1.01]	25.0%
$\tilde{c}_{HWB}/\Lambda^2$	no	[−1.06, 1.06]	[0.23, 2.34]	1.7%
	yes	[−1.06, 1.06]	[0.23, 2.35]	1.6%

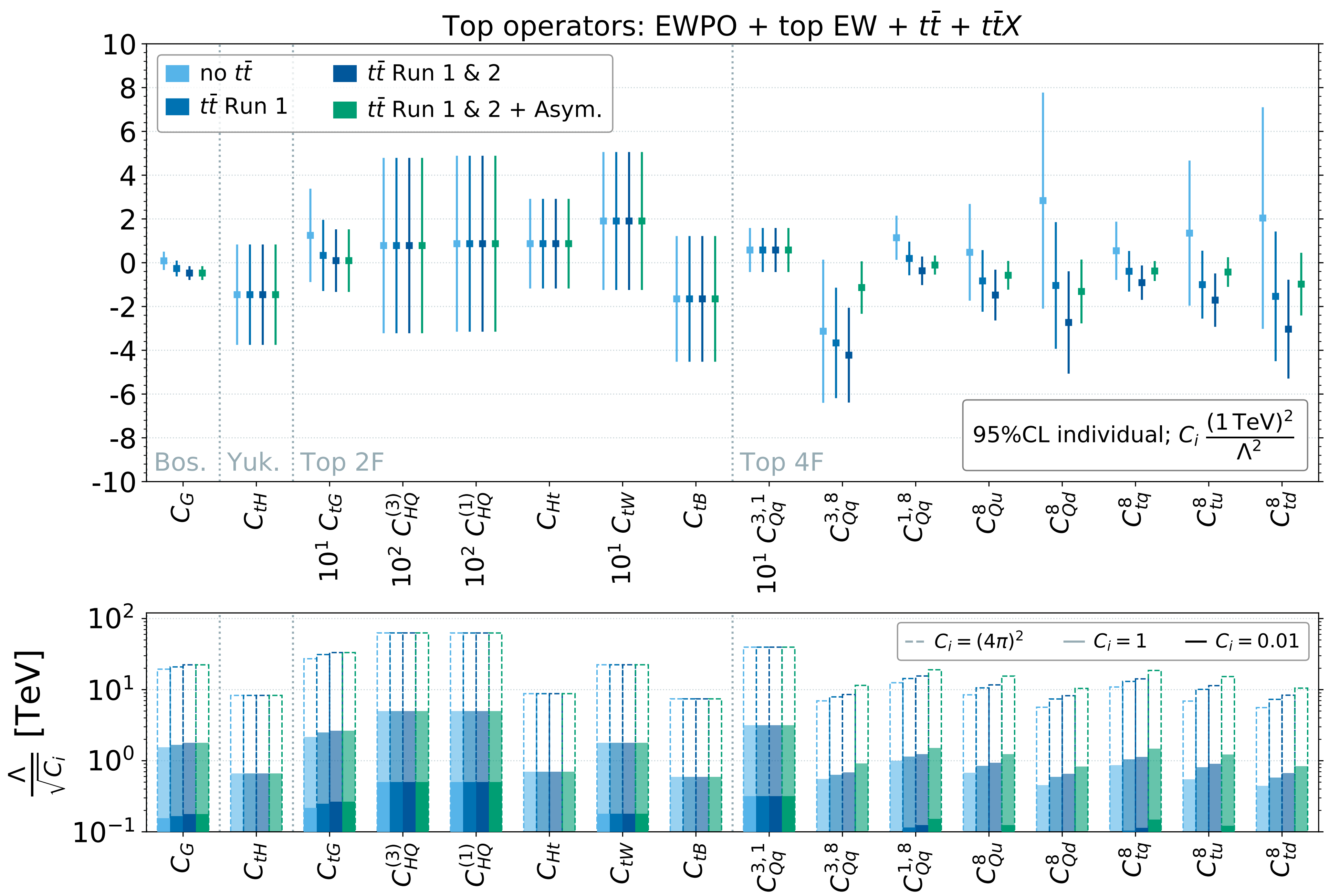


Marginalised limits: $U(3)^5$

2018 ⇒ 2020: only LHC data changed



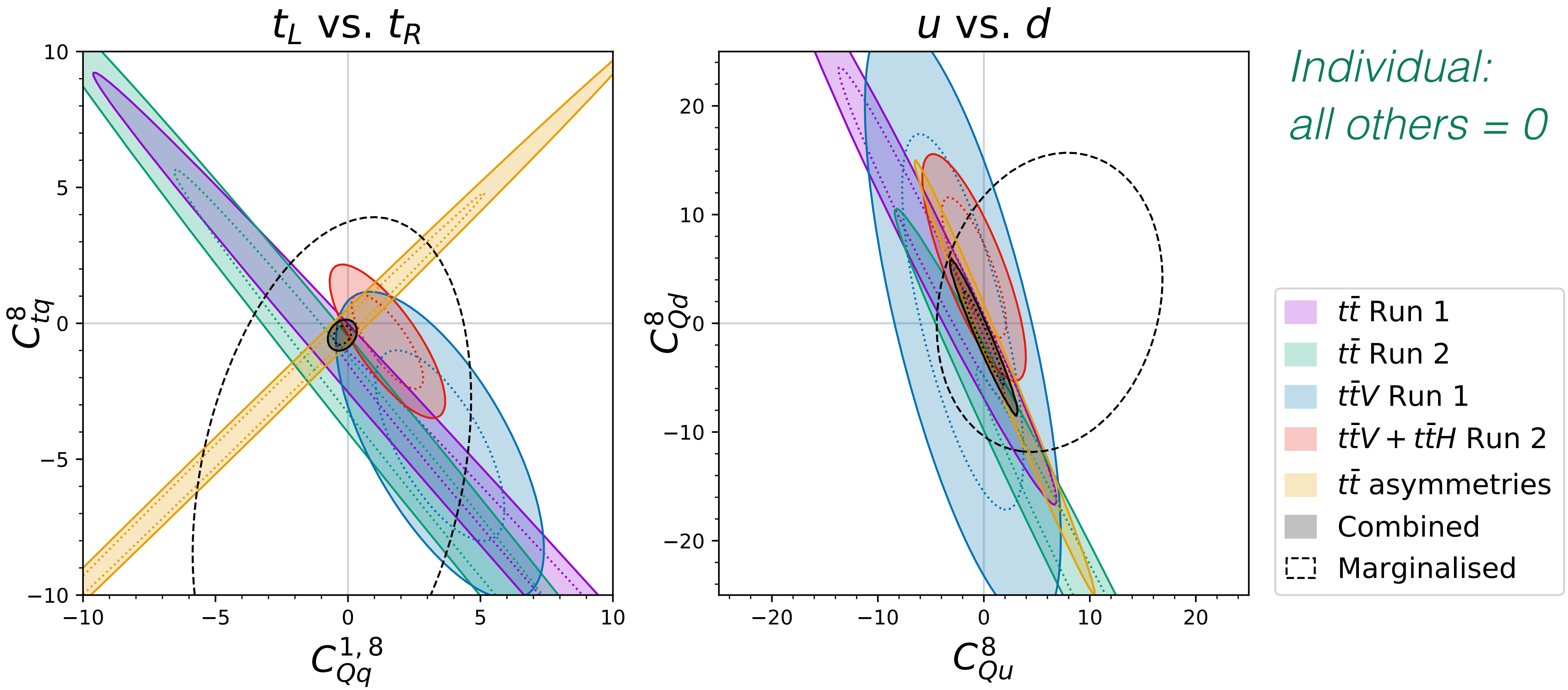
Top-only: top + EWPO individual



- Some tension in $t\bar{t}$ data
- Asymmetries help to improve agreement

See also dedicated top fits: SMEFiT [Hartland et al.; JHEP 04 (2019) 100]
SFitter [Brivio et al.; JHEP 02 (2020) 131]

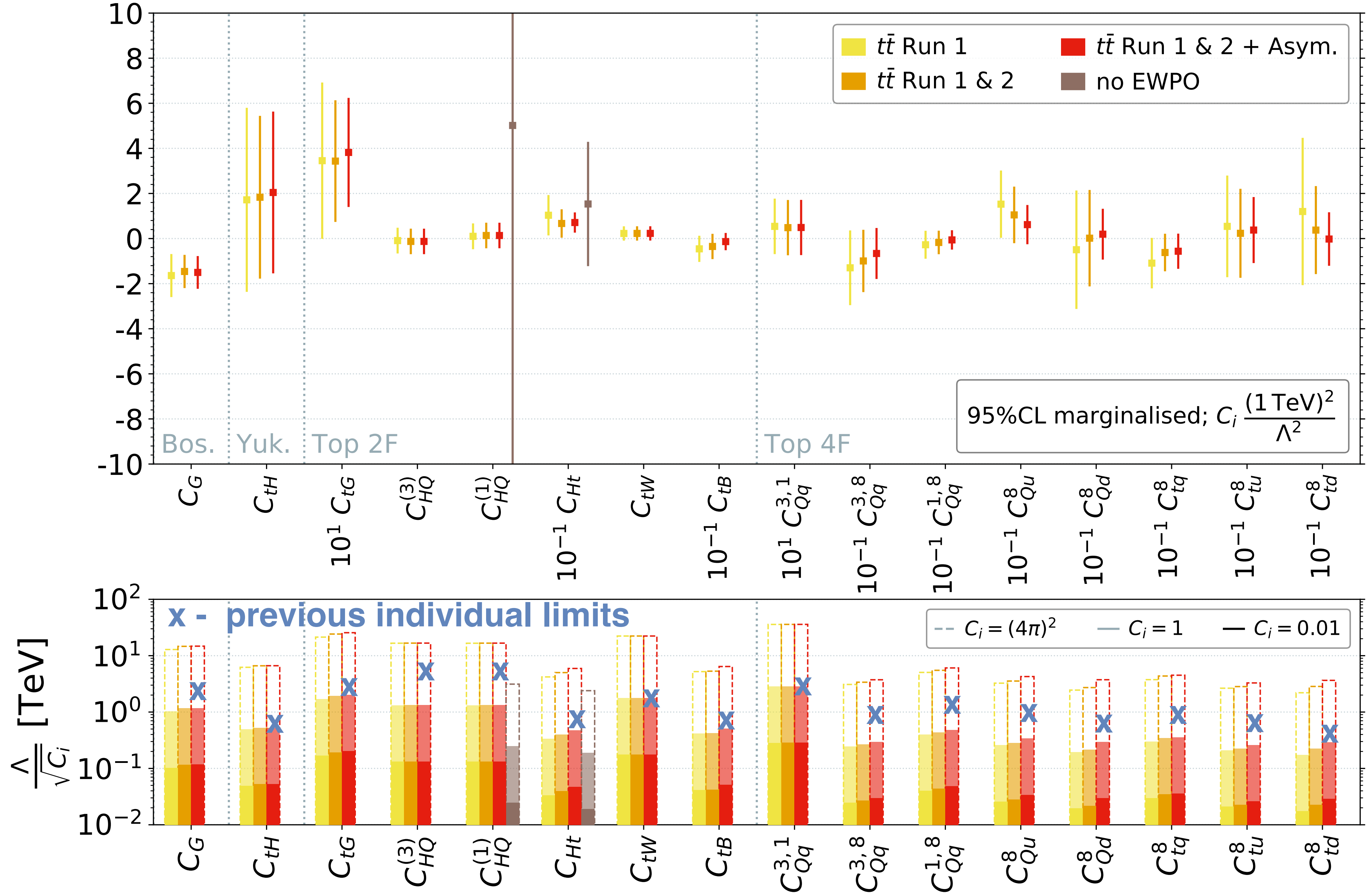
Top-only: breakdown



Individual: all others = 0

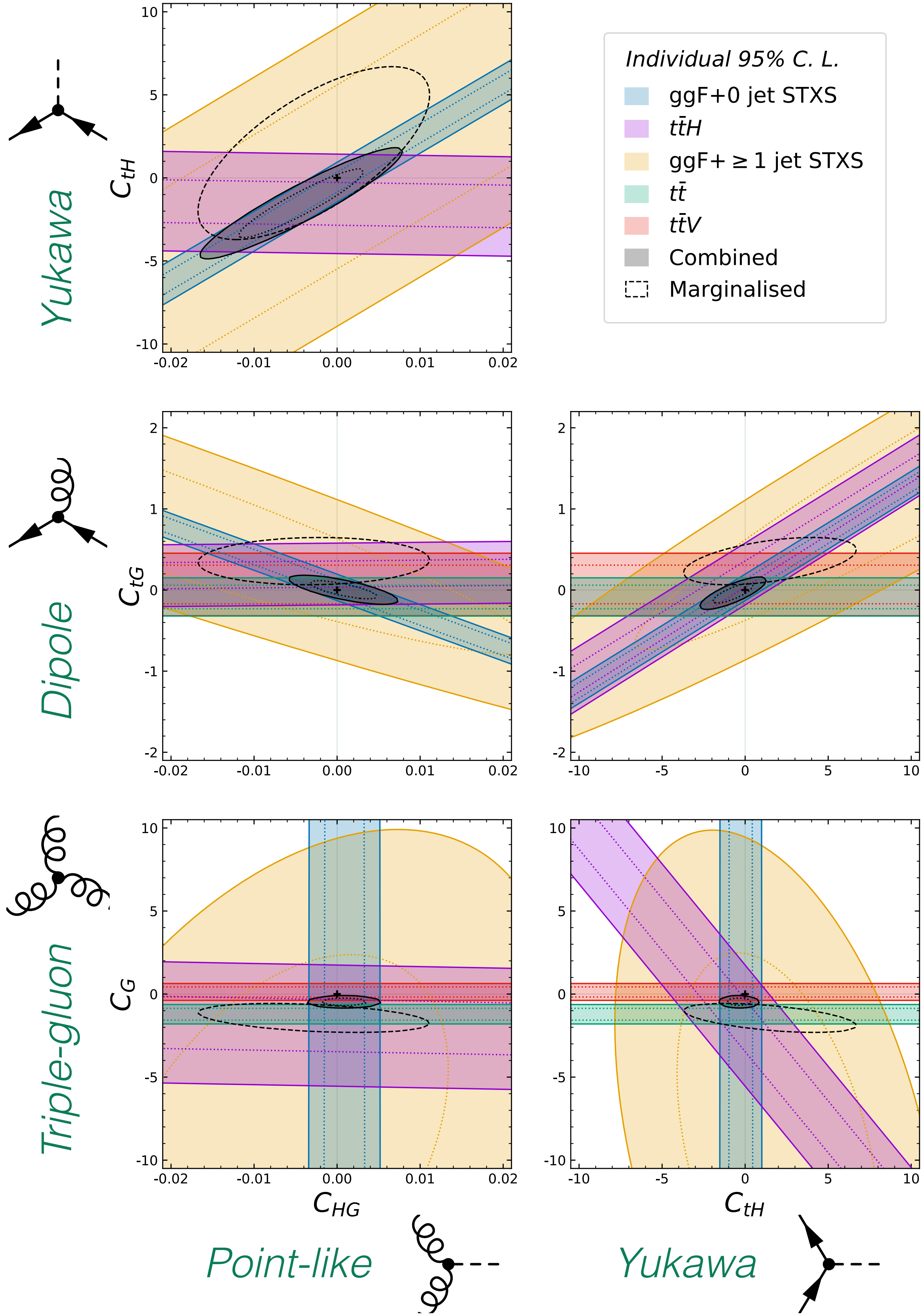
- $t\bar{t}$ asymmetries constrain orthogonal direction to cross section
- Large marginalisation effects: many similar operators
- $t\bar{t}V$ & $t\bar{t}H$ help to close the space
- Marginalised linear sensitivity: $C_{4F}\left[\frac{1\,\mathrm{TeV}^2}{\Lambda^2}\right] \sim (5-15)$ *significant* $\frac{1}{\Lambda^4}$ *effects*

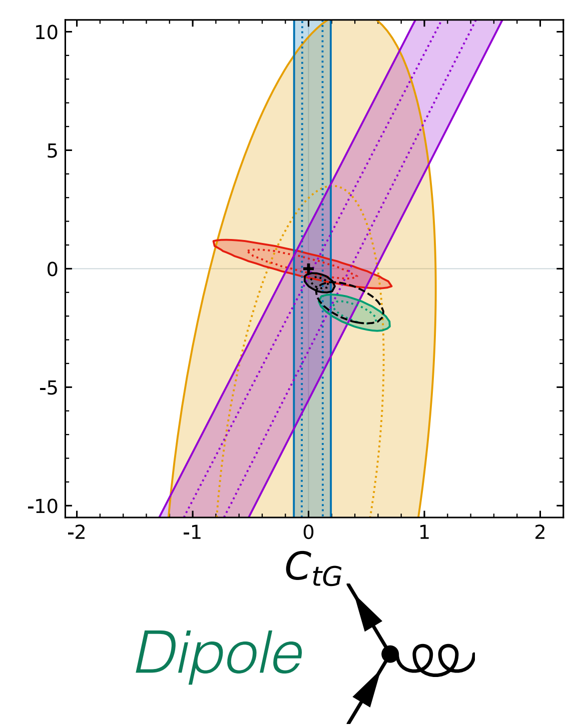
Top-only: top + EWPO marginalised



- C_{tH}: $t\bar{t}H$ bound alone is quite weak
- C_{tG}: Strong constraint but tension with SM
- Neutral top couplings poorly constrained
- EWPO closes $Zb\bar{b}$ coupling direction
- Impact of asymmetries in 4F
- Somewhat low scales (validity?)

Top-Higgs interplay

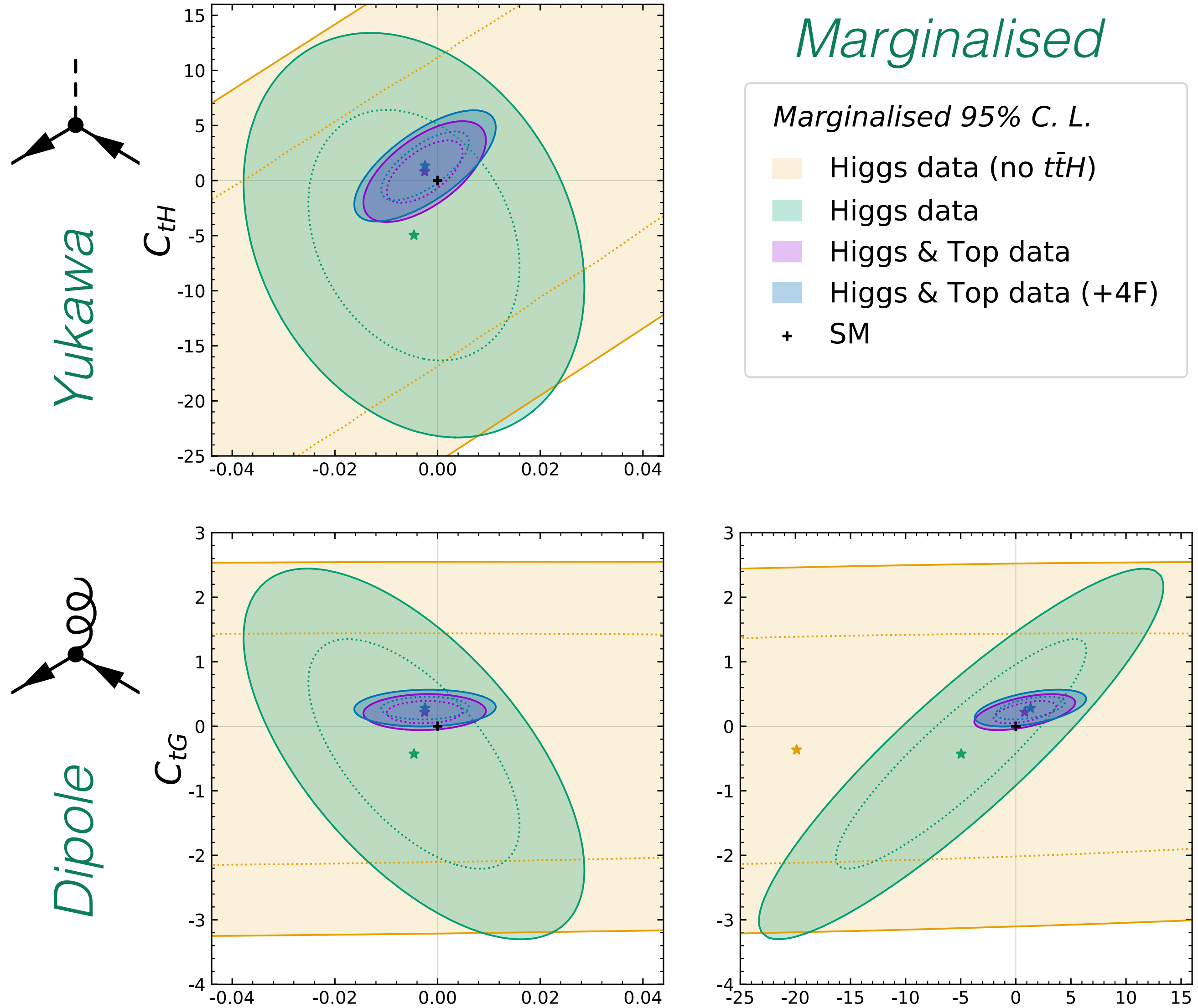


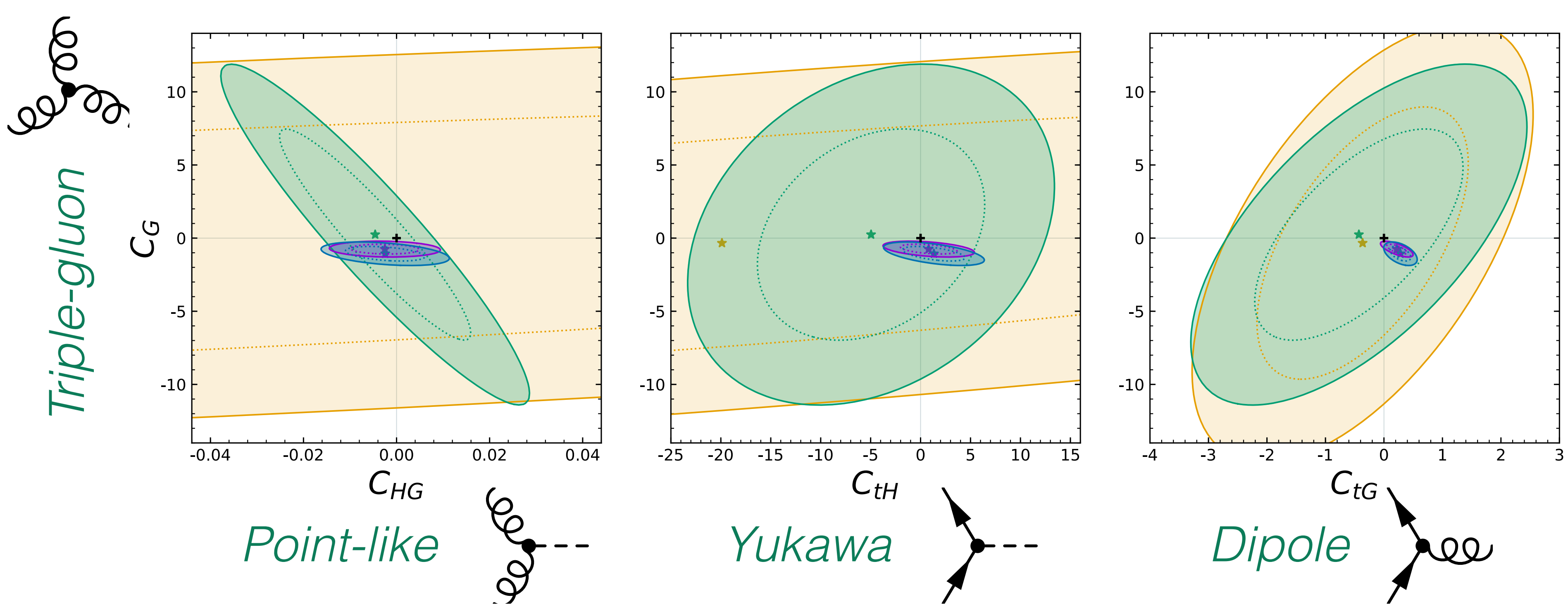


2D individual constraints

- All others set to 0
- $ggF/t\bar{t}H$ complementarity for (C_{HG}, C_{tH})
- H+jets STXS & $t\bar{t}V$ not yet competitive
- Strong impact of $t\bar{t}$ evident for (C_{tG}, C_G)
- Tension with SM $\sim 2\sigma$
- Significant correlations remain
- Large marginalisation effects

What is the concrete impact of 4F?





4F impact

Fit to 'Higgs-only' subspace

$C_{H\Box}, C_{HG}, C_{HW}, C_{HB}, C_{tH}, C_{bH}, C_{\tau H}, C_{\mu H}$
$+ \ C_{tG} \ \& \ C_G$

- Allow a closed fit to Higgs data only
- Emphasises impact of $t\bar{t}H$ & $t\bar{t}$

Now add in $t\bar{t}$ 4F operators

$+ \ C^{3,8}_{Qq}, C^{1,8}_{Qq}, C^{8}_{Qu}, C^{8}_{Qd}, C^{8}_{tq}, C^{8}_{tu}, C^{8}_{td}$

- Relatively mild impact
- Preferred $t\bar{t}$ phase space is different

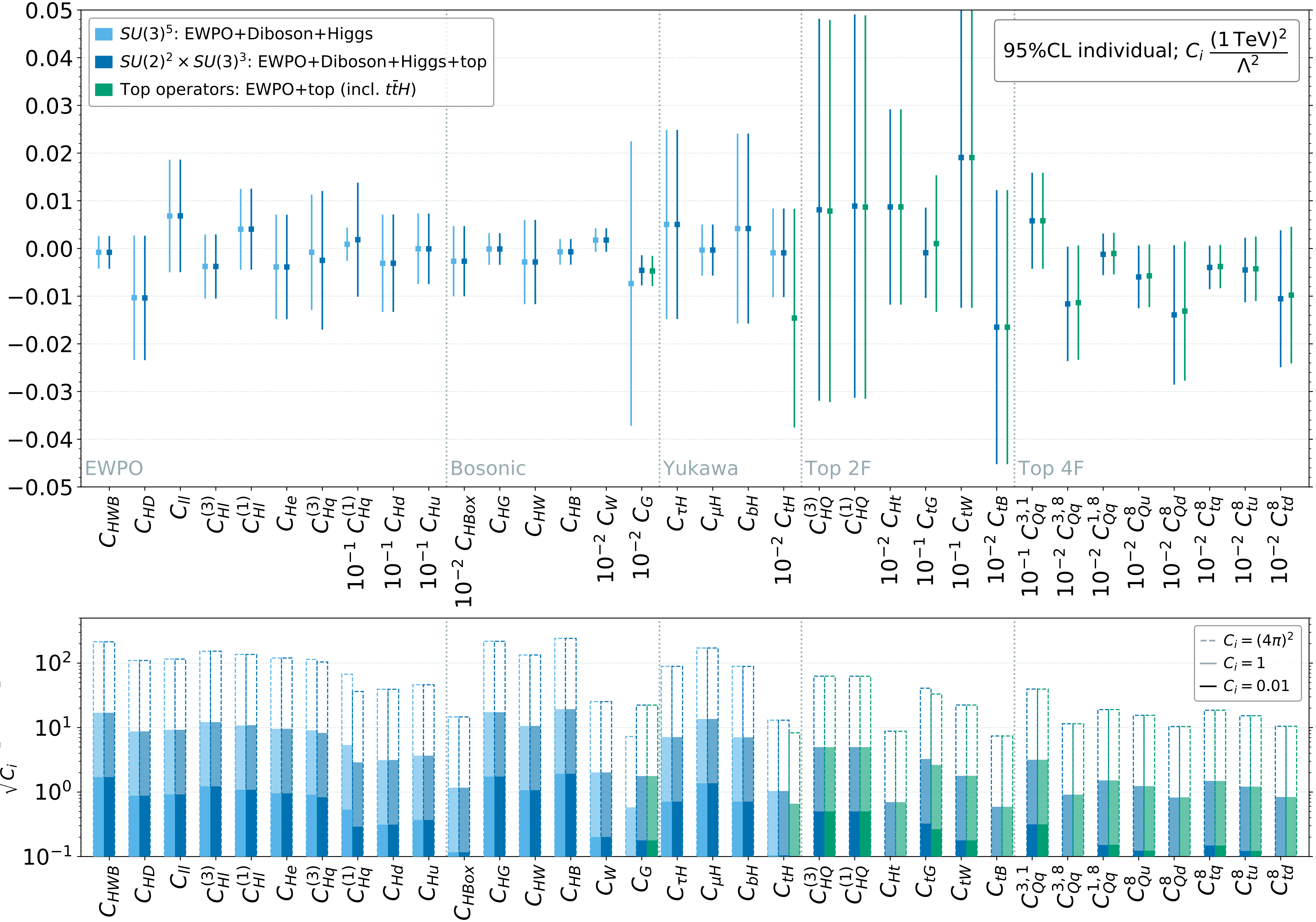
C_{tG} : low $m_{t\bar{t}}$

$4F$: high $m_{t\bar{t}}$

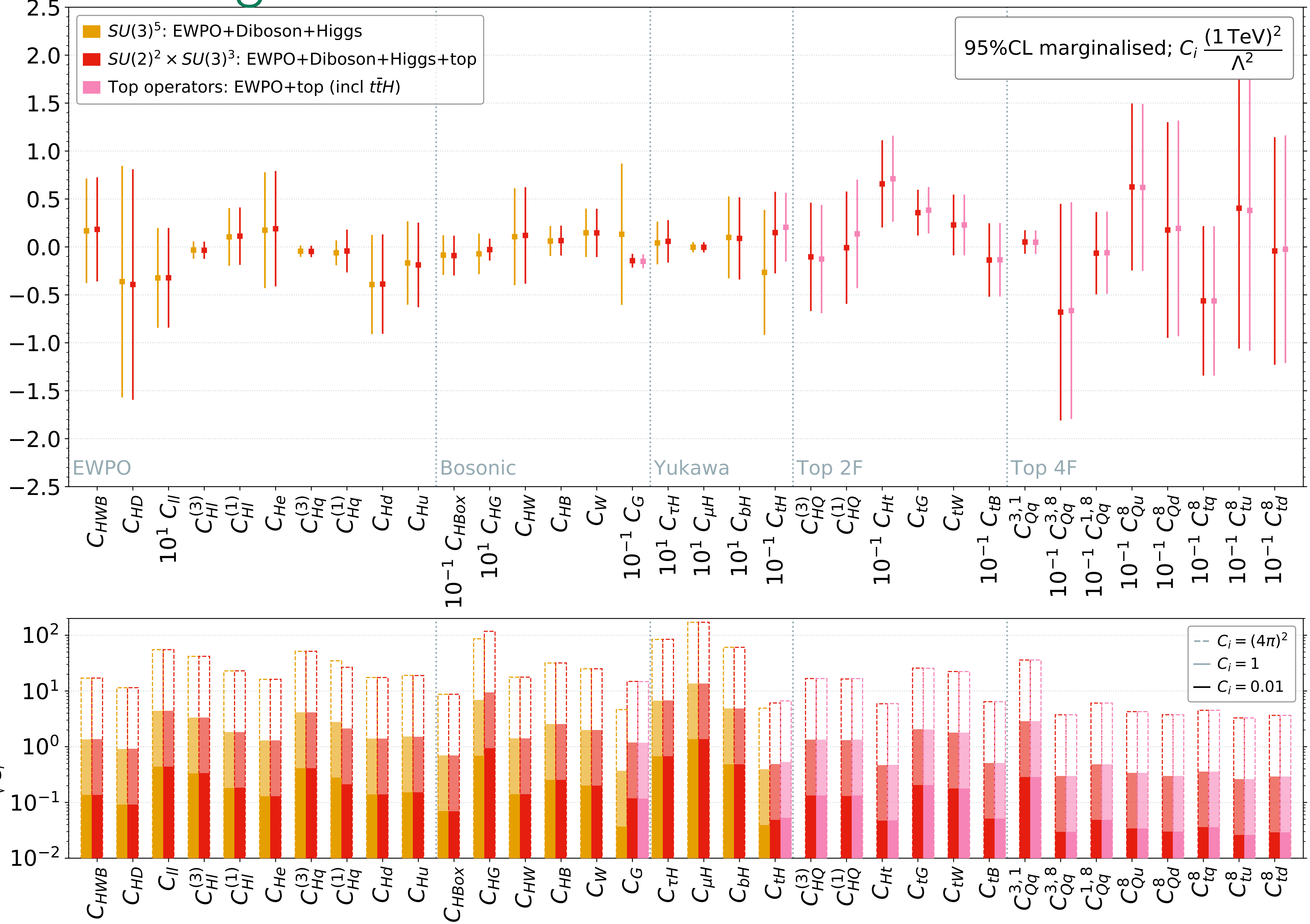
- Able to constrain them independently

[ATLAS-CONF-2021-031]

Full fit: individual



Full fit: marginalised



Correlations

Block diagonal: correlations *within* 'sector'

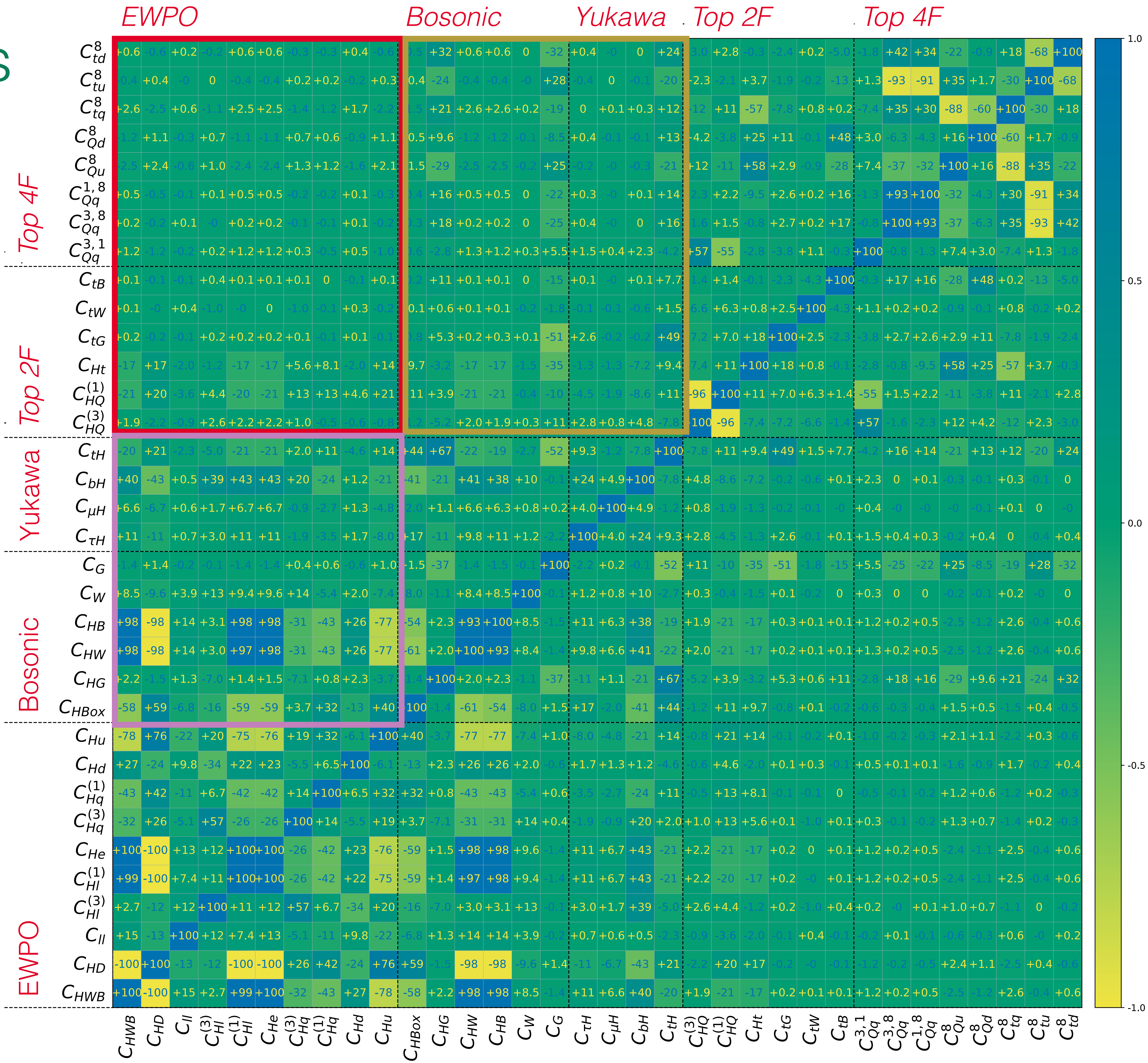
Block off-diagonal: correlations *among* 'sectors'

EWPO & top ~uncorrelated

EWPO-Higgs $C_{HB}, C_{HW}, C_{H\Box}$ & Yukawa with EWPO

Higgs precision rivalling LEP

Top-Higgs C_{HG}, C_G, C_{tH} with 4F



PCA

Relative constraining power (%)

EWPO	LEP WW	Run 1, μ_F	Run 2, μ_F	STXS	LHC WW/…	LHC Zjj	$t\bar{t}$	single to…	$t\bar{t}V$
6	–	15	29	50	–	–	–	–	–
35	–	13	28	23	–	–	–	–	–
59	–	7	16	18	–	–	–	–	–
–	–	5	94	1	–	–	–	–	–
99	–	–	–	–	–	–	–	–	–
95	–	–	–	–	3	–	–	–	–
99	–	–	–	–	–	–	–	–	–
–	–	21	45	33	–	–	–	–	–
74	–	2	4	14	6	–	–	–	–
18	–	8	19	55	–	–	–	–	–
73	22	–	–	–	3	–	–	–	–
7	2	5	53	4	29	–	–	–	–
–	–	–	3	–	–	–	–	96	–
7	–	2	39	2	46	–	–	3	–
–	–	–	10	4	–	–	57	–	28
94	2	–	–	–	2	–	–	–	–
–	2	–	–	–	11	86	–	–	–
65	14	2	4	10	4	1	–	–	–
–	–	–	–	–	–	–	99	–	–
–	–	–	–	–	–	–	84	15	–
–	–	–	13	6	–	–	48	–	32
90	5	–	–	1	2	–	–	–	–
–	–	–	5	2	–	–	79	–	13
–	–	–	–	–	–	–	3	78	18
2	3	27	40	12	6	–	4	–	4
–	–	1	2	–	–	–	15	4	77
4	4	14	33	23	7	1	12	–	2
–	–	1	3	1	–	–	4	–	89
–	–	3	10	4	–	–	28	2	53
1	–	8	27	16	2	–	32	–	12
–	–	–	–	–	–	–	58	–	41
–	–	–	–	–	–	–	90	–	10
–	–	–	–	–	–	–	61	–	38
–	–	–	–	–	–	–	95	–	5

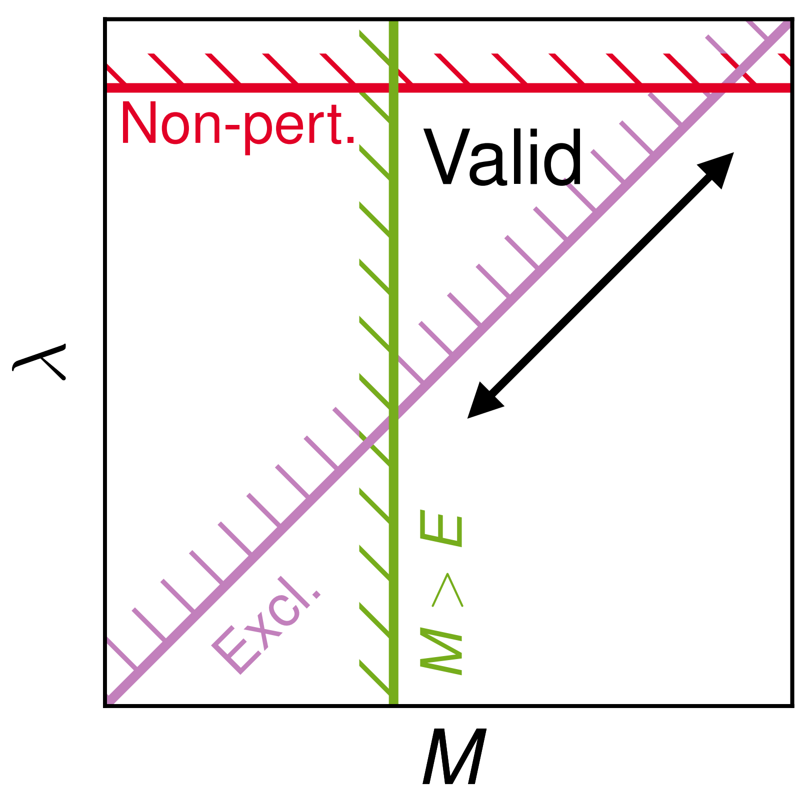
BSM implications

SMEFT-UV connection is model dependent by construction

- Implications on heavy new physics & validity of EFT is ***a posteriori***
- Depends on sensitivity & energy scale probed by data
- Bottom-up philosophy: new physics scale unknown

arbitrary dimensionful parameter $$\frac{c_S}{\Lambda^2} = \frac{\lambda^2}{M^2}$$ *coupling/mass scale of new physics*

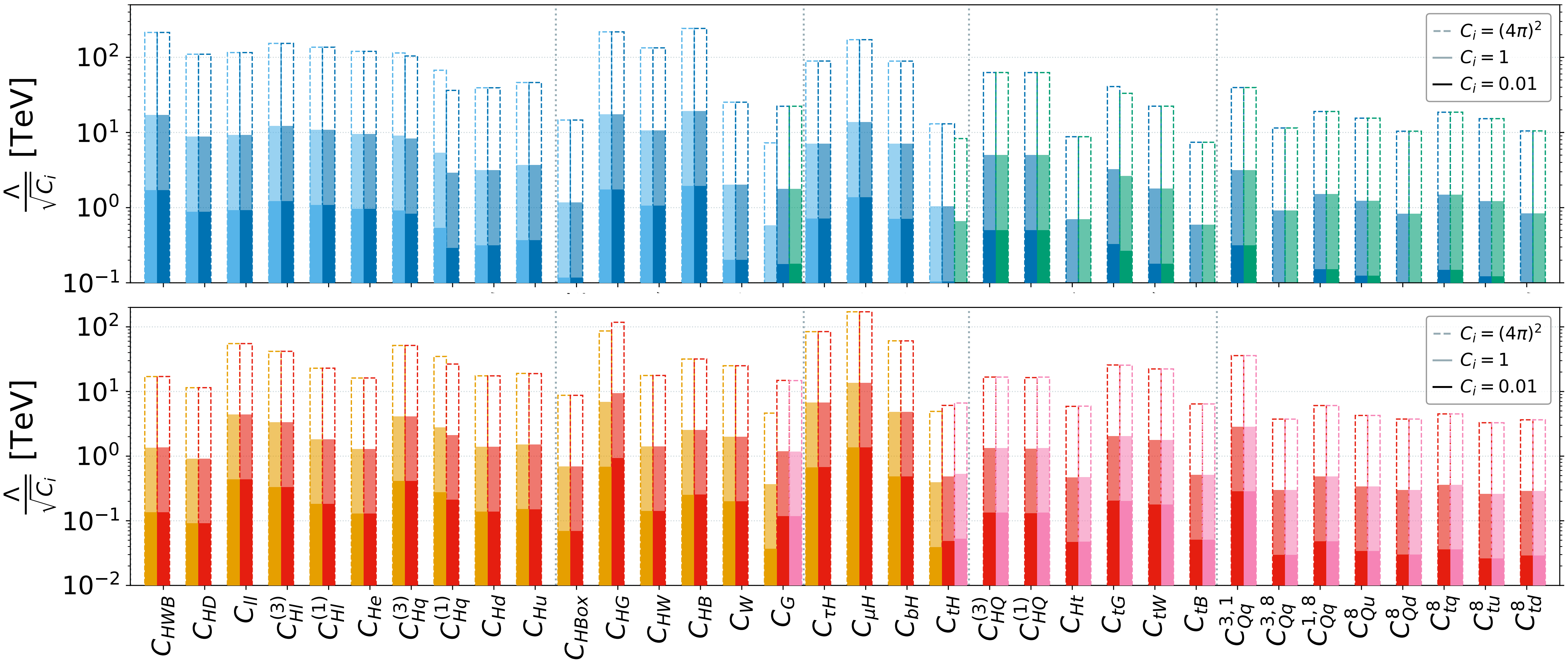
constraint: $c/\Lambda^2 < X$



Difficult to address in a general way

- Today we are probing TeV scale new physics
- Hierarchies in sensitivity EWPO > Higgs > top (EW)
- Moderate-to-strong coupling scenarios most safe
- Generic NP in loops looks challenging for the LHC

BSM implications



Individual/marginalised = optimistic/pessimistic

- Real models should lie somewhere in between
- Less underlying parameters - more correlations
- Need to ‘re-run’ the fits to infer on underlying model parameters

Single field extensions

	Name	Spin	SU(3)	SU(2)	U(1)	Param.	Name	Spin	SU(3)	SU(2)	U(1)	Param.	
	S	0	1	1	0	$(M_S,\ \kappa_S)$	Δ_1	$\frac{1}{2}$	1	2	$-\frac{1}{2}$	$(M_{\Delta_1},\lambda_{\Delta_1})$	
	S_1	0	1	1	1	(M_{S_1},y_{S_1})	Δ_3	$\frac{1}{2}$	1	2	$-\frac{1}{2}$	$(M_{\Delta_3},\lambda_{\Delta_3})$	
Scalars	φ	0	1	2	$\frac{1}{2}$	$(M_\varphi,Z_6\cos\beta)$	Σ	$\frac{1}{2}$	1	3	0	$(M_\Sigma,\lambda_\Sigma)$	VLL
	Ξ	0	1	3	0	(M_Ξ,κ_Ξ)	Σ_1	$\frac{1}{2}$	1	3	-1	$(M_{\Sigma_1},\lambda_{\Sigma_1})$	
	Ξ_1	0	1	3	1	$(M_{\Xi_1},\kappa_{\Xi_1})$	U	$\frac{1}{2}$	3	1	$\frac{2}{3}$	(M_U,λ_U)	
Z'	B	1	1	1	0	$(M_B,\hat{g}_H^B)$	D	$\frac{1}{2}$	3	1	$-\frac{1}{3}$	(M_D,λ_D)	
W'	B_1	1	1	1	1	(M_{B_1},g_{B_1})	Q_1	$\frac{1}{2}$	3	2	$\frac{1}{6}$	(M_{Q_1},λ_{Q_1})	
	W	1	1	3	0	$(M_W,\hat{g}_H^W)$	Q_5	$\frac{1}{2}$	3	2	$-\frac{5}{6}$	(M_{Q_5},λ_{Q_5})	VLQ
	W_1	1	1	3	1	$(M_{W_1},\hat{g}_{W_1}^\varphi)$	Q_7	$\frac{1}{2}$	3	2	$\frac{7}{6}$	(M_{Q_7},λ_{Q_7})	
VLL	N	$\frac{1}{2}$	1	1	0	(M_N,λ_N)	T_1	$\frac{1}{2}$	3	3	$-\frac{1}{3}$	(M_{T_1},λ_{T_1})	
	E	$\frac{1}{2}$	1	1	-1	(M_E,λ_E)	T_2	$\frac{1}{2}$	3	3	$\frac{2}{3}$	(M_{T_2},λ_{T_2})	
VLQ	T	$\frac{1}{2}$	3	1	$\frac{2}{3}$	(M_T,s_L^t)	TB	$\frac{1}{2}$	3	2	$\frac{1}{6}$	$(M_{TB},s_L^{t,b})$	

Considered single field extensions of the SM

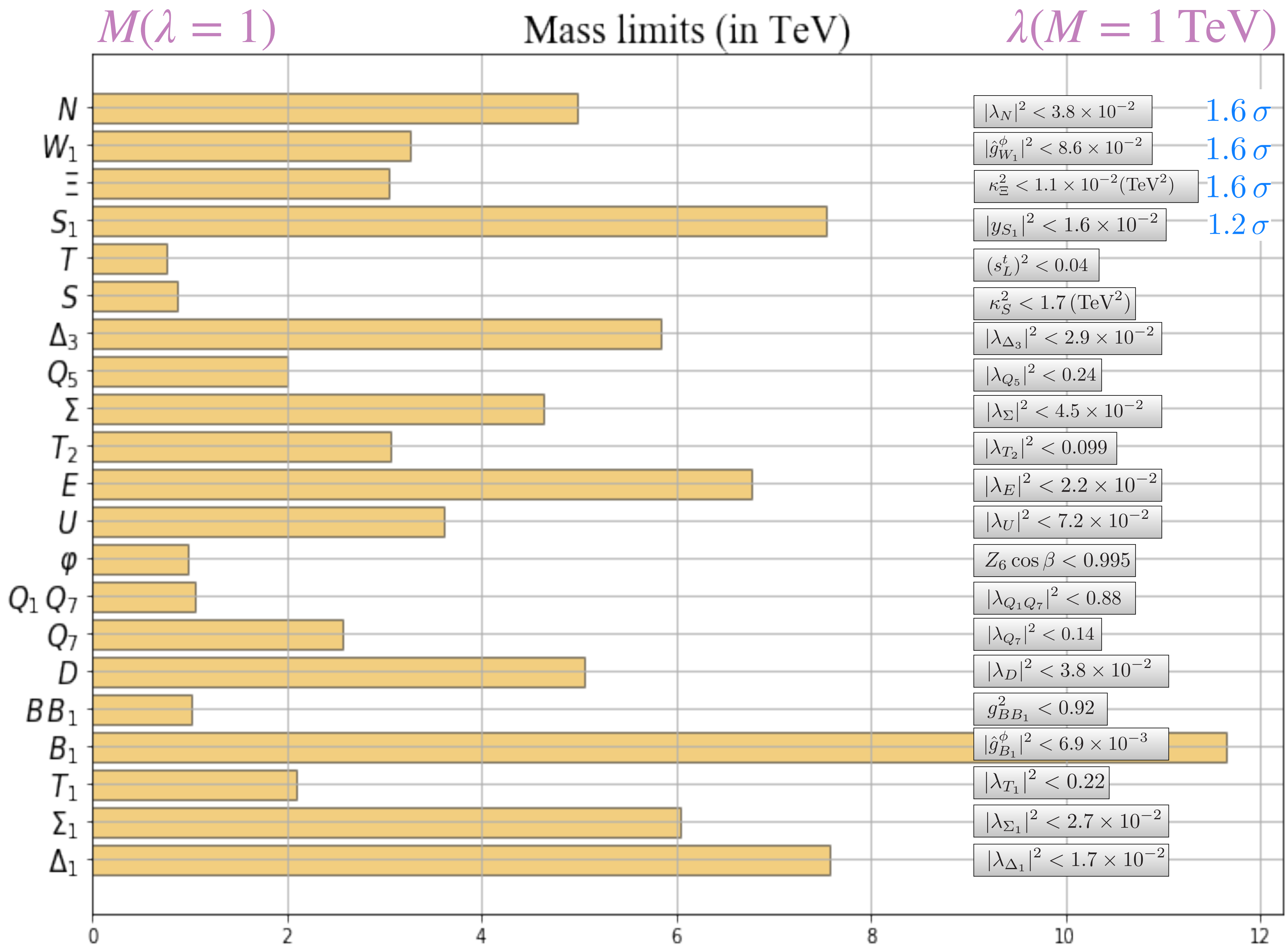
- Complete tree-level matching dictionary is known *[de Blas et al.; JHEP 03 (2018) 109]*
- Interpret in terms of simplified 1 & 2 parameter versions of the models

One parameter models

Model	C_{HD}	C_{ll}	C_{Hl}^{3}	C_{Hl}^{1}	C_{He}	$C_{H\Box}$	$C_{\tau H}$	C_{tH}	C_{bH}
S						$-\frac{1}{2}$			
S_1		1							
Σ			$\frac{1}{16}$	$\frac{3}{16}$			$\frac{y_\tau}{4}$		
Σ_1			$-\frac{1}{16}$	$-\frac{3}{16}$			$\frac{y_\tau}{8}$		
N			$-\frac{1}{4}$	$\frac{1}{4}$					
E			$-\frac{1}{4}$	$-\frac{1}{4}$			$\frac{y_\tau}{2}$		
Δ_1					$\frac{1}{2}$		$\frac{y_\tau}{2}$		
Δ_3					$-\frac{1}{2}$		$\frac{y_\tau}{2}$		
B_1	1					$-\frac{1}{2}$	$-\frac{y_\tau}{2}$	$-\frac{y_t}{2}$	$-\frac{y_b}{2}$
Ξ	-2					$\frac{1}{2}$	y_τ	y_t	y_b
W_1	$-\frac{1}{4}$					$-\frac{1}{8}$	$-\frac{y_\tau}{8}$	$-\frac{y_t}{8}$	$-\frac{y_b}{8}$
φ							$-y_\tau$	$-y_t$	$-y_b$
$\{B, B_1\}$						$-\frac{3}{2}$	$-y_\tau$	$-y_t$	$-y_b$
$\{Q_1, Q_7\}$								y_t	
Model	C_{Hq}^{3}	C_{Hq}^{1}	$(C_{Hq}^{3})_{33}$	$(C_{Hq}^{1})_{33}$	C_{Hu}	C_{Hd}	C_{tH}	C_{bH}	
U	$-\frac{1}{4}$	$\frac{1}{4}$	$-\frac{1}{4}$	$\frac{1}{4}$			$\frac{y_t}{2}$		
D	$-\frac{1}{4}$	$-\frac{1}{4}$	$-\frac{1}{4}$	$-\frac{1}{4}$				$\frac{y_b}{2}$	
Q_5						$-\frac{1}{2}$		$\frac{y_b}{2}$	
Q_7					$\frac{1}{2}$		$\frac{y_t}{2}$		
T_1	$-\frac{1}{16}$	$-\frac{3}{16}$	$-\frac{1}{16}$	$-\frac{3}{16}$			$\frac{y_t}{4}$	$\frac{y_b}{8}$	
T_2	$-\frac{1}{16}$	$\frac{3}{16}$	$-\frac{1}{16}$	$\frac{3}{16}$			$\frac{y_t}{8}$	$\frac{y_b}{4}$	
T			$-\frac{1}{2}\frac{M_T^2}{v^2}$	$\frac{1}{2}\frac{M_T^2}{v^2}$			$y_t \frac{M_T^2}{v^2}$		

$$\times \frac{\lambda^2}{M^2}$$

One parameter models

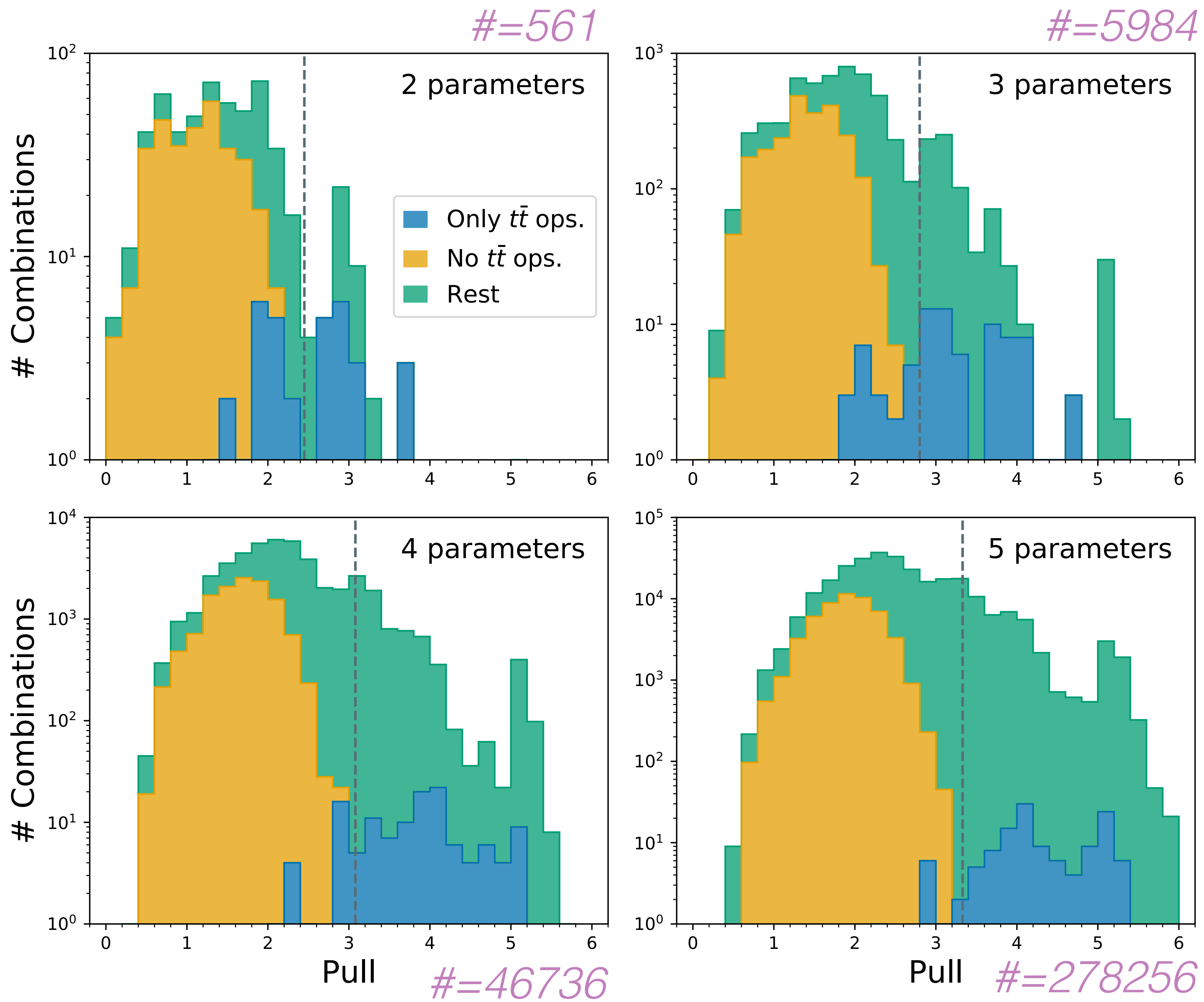


Pull-ology

Brute force: fit to all combinations of n-coefficients

$$P \equiv \sqrt{\chi^2_{SM} - \chi^2_{\text{fit}}}$$

- Agnostic search for improved fit w.r.t SM
- NP hints could show up in this way
- Advantage of fast, linear fit method
- Highlights tension in $t\bar{t}$ data
- No conclusive NP hints so far…



Conclusions

Presented first EWPO, Higgs, Diboson & Top fit in SMEFT

- Include leading contributions from top operators in ggF
- Top & Higgs sector are starting to talk to each other
- $t\bar{t}$ 4 fermion operators don't appear to spoil naive picture of interplay

Analytic, linear analysis has many benefits

- Simple likelihood described by best fit+correlations, PCA exact
- Easy to interpret/combine with other likelihoods
- Fast: repeat for subsets, BSM interpretations

& Drawbacks

- Potentially important quadratic effects, especially in top data
- Gaussian priors only, not really appropriate for th. errors

Outlook

Much more to be done

- Explore the likelihood further - how? opportunity for exp/th exchange
- Compare results to a quadratic fit to test validity
- SMEFT theory errors?
- Explore impact of new observables: VBS, VVV, rare top modes

Impact of loops

- Top operators in loops: Higgs decays + EWPO
- NLO corrections

BSM implications

- Go beyond 1 particle benchmarks towards realistic models
- Are there compelling top/Higgs scenarios that admit a valid EFT interpretation with LHC data?

Backup

Data: EWPO & Diboson

EW precision observables	$n_{\mathbf{obs}}$
Precision electroweak measurements on the Z resonance. Γ_Z, $\sigma^0_{\mathrm{had.}}$, R^0_ℓ, A^ℓ_{FB}, $A_\ell(\mathrm{SLD})$, $A_\ell(\mathrm{Pt})$, R^0_b, R^0_c A^b_{FB}, A^c_{FB}, A_b & A_c	12
Combination of CDF and D0 W-Boson Mass Measurements	1
LHC run 1 W boson mass measurement by ATLAS	1

Diboson LEP & LHC	$n_{\mathbf{obs}}$
W^+W^- angular distribution measurements at LEP II.	8
W^+W^- total cross section measurements at L3 in the $\ell\nu\ell\nu$, $\ell\nu qq$ & $qqqq$ final states for 8 energies	24
W^+W^- total cross section measurements at OPAL in the $\ell\nu\ell\nu$, $\ell\nu qq$ & $qqqq$ final states for 7 energies	21
W^+W^- total cross section measurements at ALEPH in the $\ell\nu\ell\nu$, $\ell\nu qq$ & $qqqq$ final states for 8 energies	21
ATLAS W^+W^- differential cross section in the $e\nu\mu\nu$ channel, $\frac{d\sigma}{dp^T_{\ell_1}}$, $p_T > 120$ GeV overflow bin	1
ATLAS W^+W^- fiducial differential cross section in the $e\nu\mu\nu$ channel, $\frac{d\sigma}{dp^T_{\ell_1}}$	14
ATLAS $W^\pm Z$ fiducial differential cross section in the $\ell^+\ell^-\ell^\pm\nu$ channel, $\frac{d\sigma}{dp^T_Z}$	7
CMS $W^\pm Z$ normalised fiducial differential cross section in the $\ell^+\ell^-\ell^\pm\nu$ channel, $\frac{1}{\sigma}\frac{d\sigma}{dp^T_Z}$	11
ATLAS Zjj fiducial differential cross section in the $\ell^+\ell^-$ channel, $\frac{d\sigma}{d\Delta\varphi_{jj}}$	12

Data: Higgs

LHC Run 1 Higgs	$n_{\mathbf{obs}}$
ATLAS and CMS LHC Run 1 combination of Higgs signal strengths. Production: ggF, VBF, ZH, WH & ttH Decay: $\gamma\gamma$, ZZ, W^+W^-, $\tau^+\tau^-$ & $b\bar{b}$	21
ATLAS inclusive $Z\gamma$ signal strength measurement	1

LHC Run 2 Higgs (new)	$n_{\mathbf{obs}}$
ATLAS combination of signal strengths and stage 1.0 STXS in $H \to 4\ell$ including ratios of branching fractions to $\gamma\gamma$, WW^*, $\tau^+\tau^-$ & $b\bar{b}$ Signal strengths|coarse STXS bins| fine STXS bins	16|19|25
CMS LHC combination of Higgs signal strengths. Production: ggF, VBF, ZH, WH & ttH Decay: $\gamma\gamma$, ZZ, W^+W^-, $\tau^+\tau^-$, $b\bar{b}$ & $\mu^+\mu^-$	23
CMS stage 1.0 STXS measurements for $H \to \gamma\gamma$. 13 parameter fit | 7 parameter fit	13|7
CMS stage 1.0 STXS measurements for $H \to \tau^+\tau^-$	9
CMS stage 1.1 STXS measurements for $H \to 4\ell$	19
CMS differential cross section measurements of inclusive Higgs production in the $WW^* \to \ell\nu\ell\nu$ final state. $\frac{d\sigma}{dn_{\rm jet}}$ | $\frac{d\sigma}{dp_H^T}$	5|6
ATLAS $H \to Z\gamma$ signal strength.	1
ATLAS $H \to \mu^+\mu^-$ signal strength.	1

Data: Tevatron, LHC Run 1 & 2 top

Tevatron & Run 1 top	n_{obs}	**Ref.**
Tevatron combination of differential $t\bar{t}$ forward-backward asymmetry, $A_{FB}(m_{t\bar{t}})$.	4	[7]
ATLAS $t\bar{t}$ differential distributions in the dilepton channel. $\frac{d\sigma}{dm_{t\bar{t}}}$	6	[31]
ATLAS $t\bar{t}$ differential distributions in the ℓ+jets channel. $\frac{d\sigma}{dm_{t\bar{t}}} \vert \frac{d\sigma}{d\vert y_{t\bar{t}}\vert} \vert \frac{d\sigma}{dp_T^t} \vert \frac{d\sigma}{d\vert y_t\vert}$.	$7\vert 5\vert 8\vert 5$	[24]
CMS $t\bar{t}$ differential distributions in the ℓ+jets channel. $\frac{d\sigma}{dm_{t\bar{t}}} \vert \frac{d\sigma}{dy_{t\bar{t}}} \vert \frac{d\sigma}{dp_T^t} \vert \frac{d\sigma}{dy_t}$.	$7\vert 10\vert 8 \vert 10$	[25, 34]
CMS measurement of differential $t\bar{t}$ charge asymmetry, $A_C(m_{t\bar{t}})$ in the dilepton channel.	3	[33]
ATLAS inclusive measurement $t\bar{t}$ charge asymmetry, $A_C(m_{t\bar{t}})$ in the dilepton channel.	1	[32]
ATLAS & CMS combination of differential $t\bar{t}$ charge asymmetry, $A_C(m_{t\bar{t}})$, in the ℓ+jets channel.	6	[38]
CMS $t\bar{t}$ double differential distributions in the dilepton channel. $\frac{d\sigma}{dm_{t\bar{t}}dy_t} \vert \frac{d\sigma}{dm_{t\bar{t}}dy_{t\bar{t}}} \vert \frac{d\sigma}{dm_{t\bar{t}}dp_T^{t\bar{t}}} \vert \frac{d\sigma}{dy_t dp_T^t}$.	$16\vert 16$ $\vert 16\vert 16$	[18, 35]
ATLAS & CMS Run 1 combination of W-boson helicity fractions in top decay. f_0, f_L & f_R	3	[40]
ATLAS measurement of W-boson helicity fractions in top decay. f_0, f_L & f_R	3	[30]
CMS measurement of W-boson helicity fractions in top decay. f_0, f_L & f_R	3	[29]
ATLAS $t\bar{t}W$ & $t\bar{t}Z$ cross section measurements. $\sigma_{t\bar{t}W}\vert\sigma_{t\bar{t}Z}$	2	[23]
CMS $t\bar{t}W$ & $t\bar{t}Z$ cross section measurements. $\sigma_{t\bar{t}W}\vert\sigma_{t\bar{t}Z}$	2	[26]
ATLAS $t\bar{t}\gamma$ cross section measurement in the ℓ+ jets channel.	1	[36]
CMS $t\bar{t}\gamma$ cross section measurement in the ℓ+ jets channel.	1	[37]
ATLAS t-channel single-top differential distributions. $\frac{d\sigma}{dp_T^t} \vert \frac{d\sigma}{dp_T^{\bar{t}}} \vert \frac{d\sigma}{d\vert y_t\vert} \vert \frac{d\sigma}{d\vert y_{\bar{t}}\vert}$	$4\vert 4\vert 4\vert 5$	[39]
CMS s-channel single-top cross section measurement.	1	[28]
CMS t-channel single-top differential distributions. $\frac{d\sigma}{dp_T^{t+\bar{t}}} \vert \frac{d\sigma}{d\vert y_{t+\bar{t}}\vert}$	$6 \vert 6$	[19]
CMS measurement of the t-channel single-top and anti-top cross sections. $\sigma_t \vert \sigma_{\bar{t}} \vert \sigma_{t+\bar{t}} \vert R_t$.	$1\vert 1\vert 1\vert 1$	[20]
ATLAS s-channel single-top cross section measurement.	1	[27]
CMS tW cross section measurement.	1	[21]
ATLAS tW cross section measurement in the single lepton channel.	1	[41]
ATLAS tW cross section measurement in the dilepton channel.	1	[22]

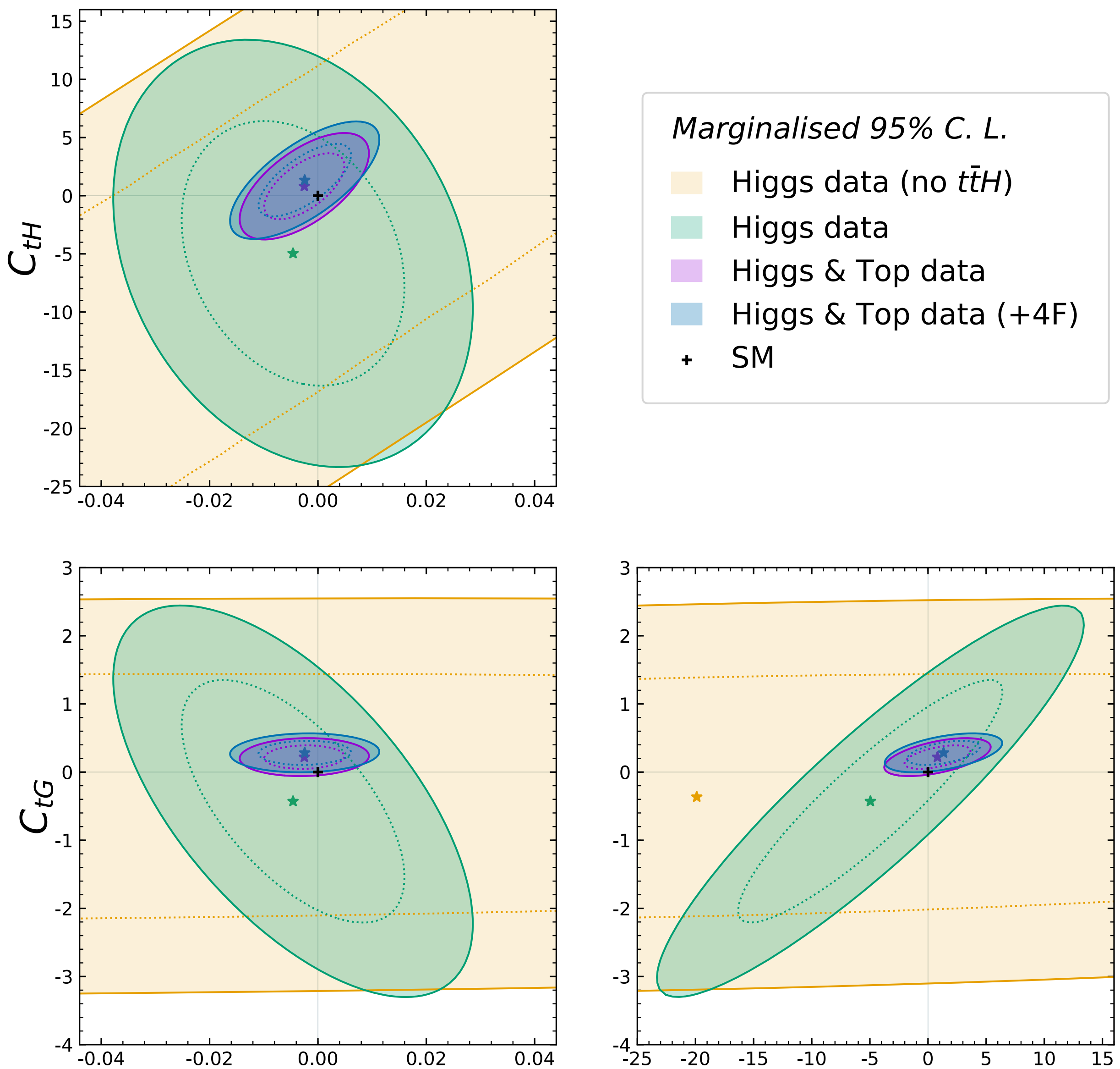
Run 2 top	n_{obs}	**Ref.**
CMS $t\bar{t}$ differential distributions in the dilepton channel. $\frac{d\sigma}{dm_{t\bar{t}}}$	6	[46, 50]
CMS $t\bar{t}$ differential distributions in the ℓ+jets channel. $\frac{d\sigma}{dm_{t\bar{t}}}$	10	[53]
ATLAS measurement of differential $t\bar{t}$ charge asymmetry, $A_C(m_{t\bar{t}})$.	5	[55]
ATLAS $t\bar{t}W$ & $t\bar{t}Z$ cross section measurements. $\sigma_{t\bar{t}W}\vert\sigma_{t\bar{t}Z}$	2	[58]
CMS $t\bar{t}W$ & $t\bar{t}Z$ cross section measurements. $\sigma_{t\bar{t}W}\vert\sigma_{t\bar{t}Z}$	$1\vert 1$	[48]
CMS $t\bar{t}Z$ differential distributions. $\frac{d\sigma}{dp_T^Z} \vert \frac{d\sigma}{d\cos\theta^*}$	$4\vert 4$	[60]
ATLAS $t\bar{t}\gamma$ differential distribution. $\frac{d\sigma}{dp_T^\gamma}$	11	[62]
CMS measurement of differential cross sections and charge ratios for t-channel single-top quark production. $\frac{d\sigma}{dp_T^{t+\bar{t}}} \vert R_t\,(p_{t+\bar{t}}^T)$	$5\vert 5$	[56]
CMS measurement of t-channel single-top and anti-top cross sections. σ_t, $\sigma_{\bar{t}}$, $\sigma_{t+\bar{t}}$ & R_t.	4	[42]
CMS measurement of the t-channel single-top and anti-top cross sections. $\sigma_t \vert \sigma_{\bar{t}} \vert \sigma_{t+\bar{t}} \vert R_t$.	$1\vert 1\vert 1\vert 1$	[45]
CMS t-channel single-top differential distributions. $\frac{d\sigma}{dp_T^{t+\bar{t}}} \vert \frac{d\sigma}{d\vert y_{t+\bar{t}}\vert}$	$4\vert 4$	[44]
ATLAS tW cross section measurement.	1	[43]
CMS tZ cross section measurement.	1	[47]
CMS tW cross section measurement.	1	[52]
ATLAS tZ cross section measurement.	1	[49]
CMS tZ $(Z \to \ell^+\ell^-)$ cross section measurement	1	[54]
ATLAS four-top search in the multi-lepton and same-sign dilepton channels.	1	[63]
ATLAS four-top search in the single-lepton and opposite-sign dilepton channels.	1	[51]
CMS four-top search in the multi-lepton and same-sign dilepton channels.	1	[61]
CMS four-top search in the single-lepton and opposite-sign dilepton channels.	1	[59]
CMS $t\bar{t}b\bar{b}$ cross section measurement in the all-jet channel.	1	[57]
CMS $t\bar{t}b\bar{b}$ cross section measurement in the dilepton channel.	1	[64]

Fisher information breakdown

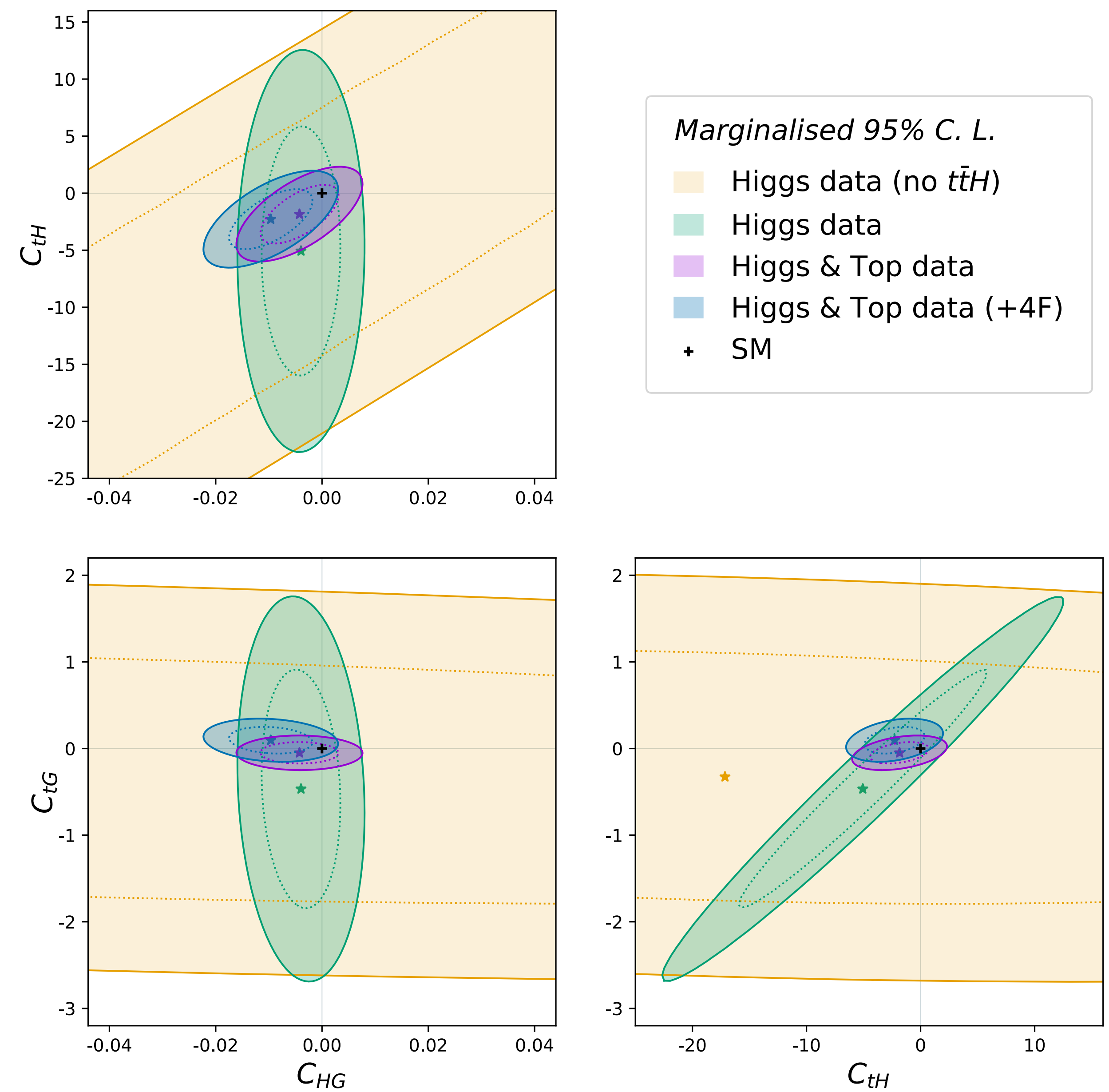
C_i	EWPO	LEP WW	Run 1 SS	Run 2 SS	STXS	LHC WW	WZ	Zjj	$t\bar{t}$	$W_{\rm hel.}$	tX	$t\bar{t}V$
C_{HWB}	51	–	7	14	28	–	–	–	–	–	–	–
C_{HD}	100	–	–	–	–	–	–	–	–	–	–	–
C_{ll}	99	–	–	–	–	–	–	–	–	–	–	–
$C_{Hl}^{(3)}$	99	–	–	–	–	–	–	–	–	–	–	–
$C_{Hl}^{(1)}$	100	–	–	–	–	–	–	–	–	–	–	–
C_{He}	100	–	–	–	–	–	–	–	–	–	–	–
$C_{Hq}^{(3)}$	89	1	–	–	2	–	6	–	–	–	–	–
$C_{Hq}^{(1)}$	99	–	–	–	–	–	–	–	–	–	–	–
C_{Hd}	99	–	–	–	–	–	–	–	–	–	–	–
C_{Hu}	98	–	–	–	1	–	–	–	–	–	–	–
$C_{H\Box}$	–	–	22	46	32	–	–	–	–	–	–	–
C_{HG}	–	–	22	42	36	–	–	–	–	–	–	–
C_{HW}	–	–	14	29	56	–	–	–	–	–	–	–
C_{HB}	–	–	14	29	57	–	–	–	–	–	–	–
C_W	–	3	–	–	–	–	13	84	–	–	–	–
C_G	–	–	–	–	–	–	–	–	43	–	–	56
$C_{\tau H}$	–	–	22	45	34	–	–	–	–	–	–	–
$C_{\mu H}$	–	–	5	95	–	–	–	–	–	–	–	–
C_{bH}	–	–	19	35	47	–	–	–	–	–	–	–
C_{tH}	–	–	21	45	34	–	–	–	–	–	–	–
$C_{HQ}^{(3)}$	99	–	–	–	–	–	–	–	–	–	–	–
$C_{HQ}^{(1)}$	100	–	–	–	–	–	–	–	–	–	–	–
C_{Ht}	–	–	–	–	–	–	–	–	–	–	–	100
C_{tG}	–	–	13	29	24	–	–	–	24	–	–	9
C_{tW}	–	–	–	–	–	–	–	–	–	84	15	–
C_{tB}	–	–	–	–	–	–	–	–	–	–	–	100
$C_{Qq}^{3,1}$	–	–	–	–	–	–	–	–	–	–	100	–
$C_{Qq}^{3,8}$	–	–	–	–	–	–	–	–	87	–	–	13
$C_{Qq}^{1,8}$	–	–	–	–	–	–	–	–	82	–	–	17
C_{Qu}^{8}	–	–	–	–	–	–	–	–	91	–	–	7
C_{Qd}^{8}	–	–	–	2	–	–	–	–	92	–	–	6
C_{tq}^{8}	–	–	–	1	–	–	–	–	89	–	–	10
C_{tu}^{8}	–	–	–	–	–	–	–	–	96	–	–	3
C_{td}^{8}	–	–	–	2	–	–	–	–	92	–	–	5

Removing C_G

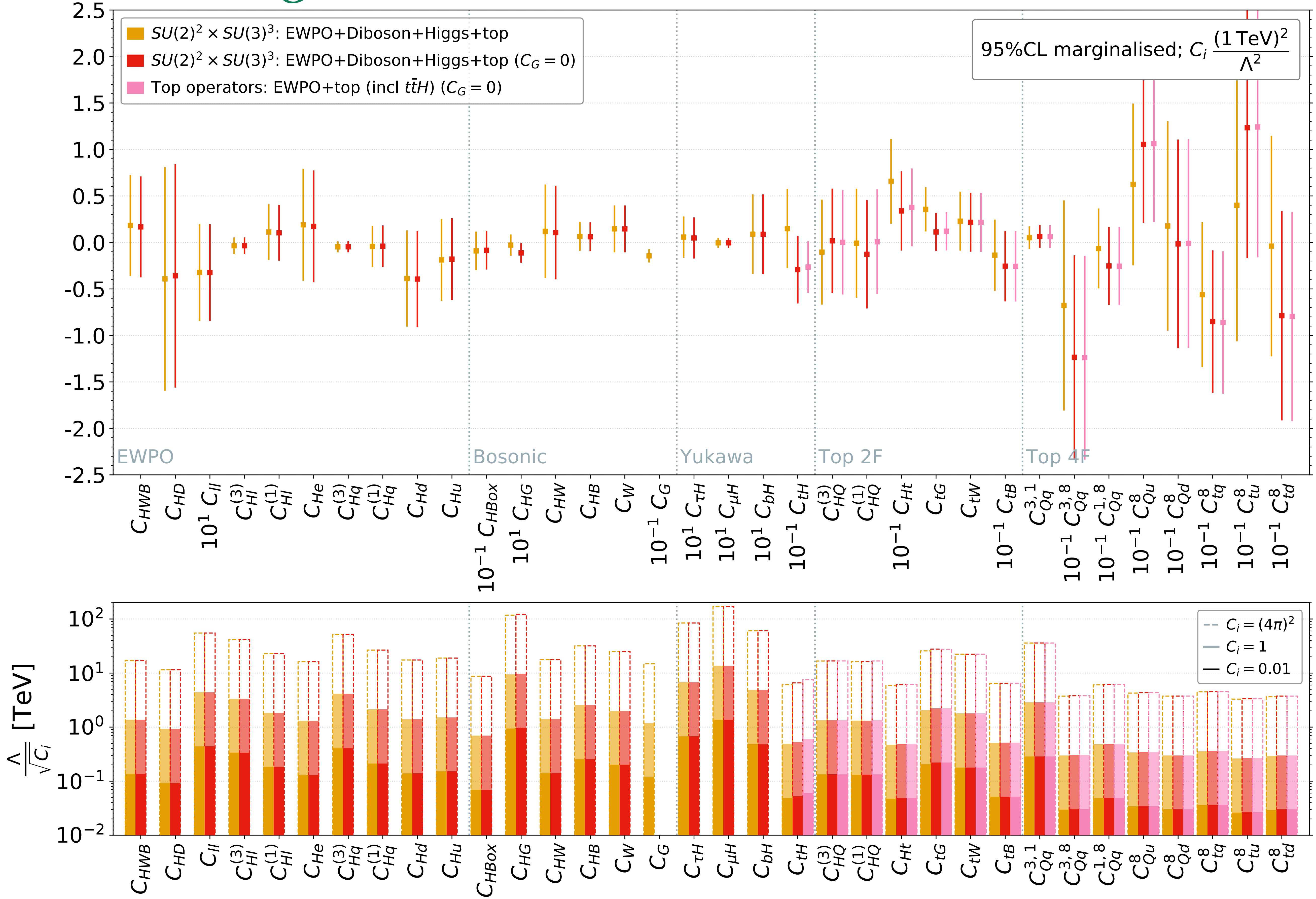
With



Without



Removing C_G



$SU(2)^2 \times SU(3)^3$: EWPO+Diboson+Higgs+top
$SU(2)^2 \times SU(3)^3$: EWPO+Diboson+Higgs+top ($C_G = 0$)
Top operators: EWPO+top (incl $t\bar{t}H$) ($C_G = 0$)
95%CL marginalised; $C_i \frac{(1\,\mathrm{TeV})^2}{\Lambda^2}$
EWPO
Bosonic
Yukawa
Top 2F
Top 4F
$C_i = (4\pi)^2$
$C_i = 1$
$C_i = 0.01$
$\frac{\Lambda}{\sqrt{C_i}}$ [TeV]