UNDERSTANDING **NEUTRINO PROPERTIES USING ASTROPHYSICAL FLUXES**

Imperial College Seminar



Jack Franklin



IPPP, Durham University

16/10/2024

JUNO AS A PROBE OF THE PSEUDO-DIRAC NATURE USING SOLAR NEUTRINOS

JF, Yuber F. Perez Gonzalez, Jessica Turner PhysRevD.108.035010

We know neutrinos have mass, but *how* do they get it?

We know neutrinos have mass, but *how* do they get it?

• We could try adding a coupling to the Higgs...

$${\cal L}_{
u} \supset -Y_{lpha i} ar{L}_{lpha} ilde{H} N^i_R$$

We know neutrinos have mass, but *how* do they get it?

• We could try adding a coupling to the Higgs...

$${\cal L}_{
u} \supset -Y_{lpha i} ar{L}_{lpha} ilde{H} N^i_R$$

• But we can also add a *Majorana* mass term

$$\mathcal{L}_{
u} \supset -Y_{lpha i} ar{L}_{lpha} ilde{H} N^i_R + rac{1}{2}ig(N^i_Rig)^c M^{ij}_R N^j_R + h.\,c.$$

• After SM spontaneous symmetry breaking via the Higgs vev we are left with:

$${\cal L}_
u \supset -rac{Y_{lpha i}v}{\sqrt{2}}ar{
u}^lpha_R N^i_R + rac{1}{2}ig(N^i_Rig)^c M^{ij}_R N^j_R + h.\,c.$$

• After SM spontaneous symmetry breaking via the Higgs vev we are left with:

$${\cal L}_
u \supset -rac{Y_{lpha i}v}{\sqrt{2}}ar{
u}^lpha_R N^i_R + rac{1}{2}ig(N^i_Rig)^c M^{ij}_R N^j_R + h.\,c.$$

• Which looks a bit nicer if we write it in matrix form:

$${\cal L}_{
u} \supset -rac{1}{2}ar{\Psi}^c M \Psi$$

with
$$\Psi = egin{pmatrix}
u_L \\ N_R^c \end{pmatrix}$$
 and $M = egin{pmatrix} 0_3 & rac{Y^T v}{\sqrt{2}} \\ rac{Y v}{\sqrt{2}} & M_R \end{pmatrix}$

• If lepton number is a global symmetry of the Universe, $M_R\equiv 0$

- If lepton number is a global symmetry of the Universe, $M_R\equiv 0$
- From an EFT perspective, M_R parameterises LNV. This could arise from:

- If lepton number is a global symmetry of the Universe, $M_R\equiv 0$
- From an EFT perspective, M_R parameterises LNV. This could arise from:
 - Quantum gravity (which may not allow global symmetries)

- If lepton number is a global symmetry of the Universe, $M_R\equiv 0$
- From an EFT perspective, M_R parameterises LNV. This could arise from:
 - Quantum gravity (which may not allow global symmetries)
 - UV theories (e.g. Majoron theory)

• If LNV is induced by QG we expect M_R to be small, and possibly even smaller than the Dirac mass!

- If LNV is induced by QG we expect M_R to be small, and possibly even smaller than the Dirac mass!
- This would result in neutrinos being *Pseudo-Dirac*

- If LNV is induced by QG we expect M_R to be small, and possibly even smaller than the Dirac mass!
- This would result in neutrinos being *Pseudo-Dirac*

$$egin{aligned} m_{k,\pm}^2 &= m_k^2 \pm rac{1}{2} \delta m_k^2 \
u_lpha &= U_{lpha k}^{3f} rac{
u_k^+ + i
u_k^-}{\sqrt{2}} \end{aligned}$$

- If LNV is induced by QG we expect M_R to be small, and possibly even smaller than the Dirac mass!
- This would result in neutrinos being *Pseudo-Dirac*

$$egin{aligned} m_{k,\pm}^2 &= m_k^2 \pm rac{1}{2} \delta m_k^2 \
u_lpha &= U_{lpha k}^{3f} rac{
u_k^+ + i
u_k^-}{\sqrt{2}} \end{aligned}$$





Source: doi:10.1103/PhysRevD.105.095019

• The new mass splittings will result in neutrino oscillations, occurring over lengths:

$$L_{
m osc} = rac{4\pi E_{
u}}{\delta m_k^2}$$

• The new mass splittings will result in neutrino oscillations, occurring over lengths:

$$egin{split} L_{
m osc} &= rac{4\pi E_
u}{\delta m_k^2} \ &pprox 25 imes 10^5\,{
m km}\left(rac{E_
u}{100\,{
m keV}}
ight)\left(rac{10^{-10}\,{
m eV}^2}{\delta m_k^2}
ight) \end{split}$$

• The new mass splittings will result in neutrino oscillations, occurring over lengths:

$$egin{split} L_{
m osc} &= rac{4\pi E_
u}{\delta m_k^2} \ &pprox 25 imes 10^5\,{
m km}\left(rac{E_
u}{100\,{
m keV}}
ight)\left(rac{10^{-10}\,{
m eV}^2}{\delta m_k^2}
ight) \end{split}$$

• Solar neutrinos:

- Solar neutrinos:
 - Low energies: $E_
 u \sim 100\,{
 m keV}$

- Solar neutrinos:
 - Low energies: $E_
 u \sim 100 \, {
 m keV}$
 - Long baseline: $L pprox 150 imes 10^6 \, {
 m km}$

- Solar neutrinos:
 - Low energies: $E_
 u \sim 100 \, {
 m keV}$
 - Long baseline: $Lpprox 150 imes 10^6~{
 m km}$



TESTING FOR THE PSEUDO-DIRAC SCENARIO





Credit: Xinhua / Alamy Stock Photo

 The Jiangmen Underground Neutrino Observatory (JUNO) Experiment:



Credit: Xinhua / Alamy Stock Photo

- The Jiangmen Underground Neutrino Observatory (JUNO) Experiment:
 - Liquid Scintillator Detector



Credit: Xinhua / Alamy Stock Photo

- The Jiangmen Underground Neutrino Observatory (JUNO) Experiment:
 - Liquid Scintillator Detector
 - 20kt fiducial volume



Credit: Xinhua / Alamy Stock Photo

- The Jiangmen Underground Neutrino Observatory (JUNO) Experiment:
 - Liquid Scintillator Detector
 - 20kt fiducial volume
 - Energy resolution of $3\%/\sqrt{E/{
 m MeV}}$



Credit: Xinhua / Alamy Stock Photo



Modelled Backgrounds at JUNO. Credit: JUNO Collaboration
Compare to events expected from solar flux models:



Events at JUNO

• Chi-squared analysis:

$$egin{split} \chi^2 = \sum_i rac{(\sum_a lpha_a N_{ ext{BSM}}^{i,a} + \sum_b (lpha_b - 1) N^{i,b} - N_{ ext{SM}}^i)^2}{N_{ ext{SM}}^i + \sum_b N_b^i} \ + \sum_a \left(rac{lpha_a - 1}{\sigma_a}
ight)^2 + \sum_b \left(rac{lpha_b - 1}{\sigma_b}
ight)^2 \end{split}$$

• Chi-squared analysis:

$$egin{split} \chi^2 = \sum_i rac{(\sum_a lpha_a N_{ ext{BSM}}^{i,a} + \sum_b (lpha_b - 1) N^{i,b} - N_{ ext{SM}}^i)^2}{N_{ ext{SM}}^i + \sum_b N_b^i} \ + \sum_a \left(rac{lpha_a - 1}{\sigma_a}
ight)^2 + \sum_b \left(rac{lpha_b - 1}{\sigma_b}
ight)^2 \end{split}$$

• a are the source fluxes (pp or $^7\mathrm{Be}$)

• Chi-squared analysis:

$$egin{split} \chi^2 = \sum_i rac{(\sum_a lpha_a N_{ ext{BSM}}^{i,a} + \sum_b (lpha_b - 1) N^{i,b} - N_{ ext{SM}}^i)^2}{N_{ ext{SM}}^i + \sum_b N_b^i} \ + \sum_a \left(rac{lpha_a - 1}{\sigma_a}
ight)^2 + \sum_b \left(rac{lpha_b - 1}{\sigma_b}
ight)^2 \end{split}$$

- a are the source fluxes (pp or $^7\mathrm{Be}$)
- *b* are the backgrounds

SCENARIO 1:



SCENARIO 1:

Single mass splitting, δm^2_{14}

• JUNO will put strong bounds on parameter space



SCENARIO 1:

- JUNO will put strong bounds on parameter space
- Only 6 years of data taking (~2031)



SCENARIO 1:

- JUNO will put strong bounds on parameter space
- Only 6 years of data taking (~2031)
- May be competitive with DARWIN (XLZD)



SCENARIO 1:

- JUNO will put strong bounds on parameter space
- Only 6 years of data taking (~2031)
- May be competitive with DARWIN (XLZD)
- Depends on ability to reduce/ model ^{14}C background



SCENARIO 2: MAXIMAL MIXING

• m_1 splitting only: $\delta m_{14}^2 \lesssim 1.5 imes 10^{-12} \, {
m eV}^2 o \delta m_{14}^2 \lesssim 3 imes 10^{-13} \, {
m eV}^2$

- m_1 splitting only: $\delta m_{14}^2 \lesssim 1.5 imes 10^{-12} \, \mathrm{eV}^2 o \delta m_{14}^2 \lesssim 3 imes 10^{-13} \, \mathrm{eV}^2$
- m_2 splitting only: $\delta m^2_{25} \lesssim 2 imes 10^{-11} \, {
 m eV}^2 o \delta m^2_{25} \lesssim 6 imes 10^{-13} \, {
 m eV}^2$

- m_1 splitting only: $\delta m_{14}^2 \lesssim 1.5 imes 10^{-12} \, {
 m eV}^2 o \delta m_{14}^2 \lesssim 3 imes 10^{-13} \, {
 m eV}^2$
- m_2 splitting only: $\delta m^2_{25} \lesssim 2 imes 10^{-11} \, {
 m eV}^2 o \delta m^2_{25} \lesssim 6 imes 10^{-13} \, {
 m eV}^2$
- All masses split equally: $\delta m^2_{14} \lesssim 1.5 imes 10^{-12} \, {
 m eV}^2 o \delta m^2_{14} \lesssim 3 imes 10^{-13} \, {
 m eV}^2$

- m_1 splitting only: $\delta m_{14}^2 \lesssim 1.5 imes 10^{-12} \, {
 m eV}^2 o \delta m_{14}^2 \lesssim 3 imes 10^{-13} \, {
 m eV}^2$
- m_2 splitting only: $\delta m^2_{25} \lesssim 2 imes 10^{-11} \, {
 m eV}^2 o \delta m^2_{25} \lesssim 6 imes 10^{-13} \, {
 m eV}^2$
- All masses split equally: $\delta m^2_{14} \lesssim 1.5 imes 10^{-12} \, {
 m eV}^2 o \delta m^2_{14} \lesssim 3 imes 10^{-13} \, {
 m eV}^2$
- Potential improvement of over an order of magnitude!

CONSTRAINS ON SECRET INTERACTIONS FROM POINT SOURCES AT ICECUBE

JF, Ivan Martinez-Soler, Yuber F. Perez Gonzalez, Jessica Turner

ICECUBE



IceCube Experiment. Credit: IceCube Collab.

Atmospheric neutrinos

Atmospheric neutrinos



Atmospheric neutrinos

Astrophysical neutrinos



Atmospheric neutrinos

Astrophysical neutrinos





Atmospheric neutrinos

Astrophysical neutrinos

• Diffuse



Atmospheric neutrinos

Astrophysical neutrinos

- Diffuse
- Point-like





Credit: IceCube Collab.

• Localised source in the sky



Credit: IceCube Collab.

- Localised source in the sky
- Associated with objects from optical catalogues



Credit: IceCube Collab.

- Localised source in the sky
- Associated with objects from optical catalogues
- Currently three significant sources:



Credit: IceCube Collab.

- Localised source in the sky
- Associated with objects from optical catalogues
- Currently three significant sources:
 - NGC 1068: AGN 14 Mpc from Milky Way



Credit: IceCube Collab.

- Localised source in the sky
- Associated with objects from optical catalogues
- Currently three significant sources:
 - NGC 1068: AGN 14 Mpc from Milky Way
 - TXS 0506+056: Blazar at redshift z = 0.45

Credit: IceCube Collab.

- Localised source in the sky
- Associated with objects from optical catalogues
- Currently three significant sources:
 - NGC 1068: AGN 14 Mpc from Milky Way
 - TXS 0506+056: Blazar at redshift z = 0.45
 - PKS 1424+240: AGN at z = 0.6

Credit: IceCube Collab.

Identify with unbinned log-likelihood ratio test:

$$\mathcal{L}(n_s,ec{ heta}|\{ec{x}_i\},N) = \prod_{i=1}^N \Big[\Big(rac{n_s}{N}\Big) f_S(ec{x}_i|ec{d}_{ ext{src}},ec{ heta}) + \Big(1-rac{n_s}{N}\Big) f_B(ec{x}_i) \Big]$$

Identify with unbinned log-likelihood ratio test:

$$\mathcal{L}(n_s,ec{ heta}|\{ec{x}_i\},N) = \prod_{i=1}^N \Big[\Big(rac{n_s}{N}\Big) f_S(ec{x}_i|ec{d}_{ ext{src}},ec{ heta}) + \Big(1-rac{n_s}{N}\Big) f_B(ec{x}_i) \Big]$$

• *N* : Total number of events

Identify with unbinned log-likelihood ratio test:

$$\mathcal{L}(n_s,ec{ heta}|\{ec{x}_i\},N) = \prod_{i=1}^N \Big[\Big(rac{n_s}{N}\Big) f_S(ec{x}_i|ec{d}_{ ext{src}},ec{ heta}) + \Big(1-rac{n_s}{N}\Big) f_B(ec{x}_i) \Big]$$

- N: Total number of events
- n_s : Estimator of number of signal events

Identify with unbinned log-likelihood ratio test:

$$\mathcal{L}(n_s,ec{ heta}|\{ec{x}_i\},N) = \prod_{i=1}^N \Big[\Big(rac{n_s}{N}\Big) f_S(ec{x}_i|ec{d}_{ ext{src}},ec{ heta}) + \Big(1-rac{n_s}{N}\Big) f_B(ec{x}_i) \Big]$$

- *N* : Total number of events
- n_s : Estimator of number of signal events

•
$$ec{x}_i = \{ \hat{E}_{\mu,i}, \hat{\delta}_i, \hat{lpha}_i, \hat{\sigma}_i \}$$
Identify with unbinned log-likelihood ratio test:

$$\mathcal{L}(n_s,ec{ heta}|\{ec{x}_i\},N) = \prod_{i=1}^N \Big[\Big(rac{n_s}{N}\Big) f_S(ec{x}_i|ec{d}_{ ext{src}},ec{ heta}) + \Big(1-rac{n_s}{N}\Big) f_B(ec{x}_i) \Big]$$

- N : Total number of events
- n_s : Estimator of number of signal events
- $ec{x}_i = \{ \hat{E}_{\mu,i}, \hat{\delta}_i, \hat{lpha}_i, \hat{\sigma}_i \}$
- $ec{d}_{
 m src} = \{\delta_{
 m src}, lpha_{
 m src}\}$: Location of source in sky

Identify with unbinned log-likelihood ratio test:

$$\mathcal{L}(n_s,ec{ heta}|\{ec{x}_i\},N) = \prod_{i=1}^N \Big[\Big(rac{n_s}{N}\Big) f_S(ec{x}_i|ec{d}_{ ext{src}},ec{ heta}) + \Big(1-rac{n_s}{N}\Big) f_B(ec{x}_i) \Big]$$

- N : Total number of events
- n_s : Estimator of number of signal events
- $ec{x}_i = \{ \hat{E}_{\mu,i}, \hat{\delta}_i, \hat{lpha}_i, \hat{\sigma}_i \}$
- $ec{d}_{
 m src} = \{\delta_{
 m src}, lpha_{
 m src}\}$: Location of source in sky
- $\vec{\theta}$: Source flux parameters

• Model fluxes as power law:

 $|\phi(E) \propto E^{-\gamma}$

• Model fluxes as power law:

 $|\phi(E) \propto E^{-\gamma}$

• Restricted to North Sky

• Model fluxes as power law:

 $|\phi(E) \propto E^{-\gamma}$

- Restricted to North Sky
- Test-statistic:

$$TS(ec{d}_{ ext{src}}) = -2\log\left(rac{\mathcal{L}(n_s=0|\{ec{x}_i\})}{\sup_{n_s,\gamma}\mathcal{L}(n_s,\gamma,ec{d}_{ ext{src}}|\{ec{x}_i\})}
ight)$$

Values from SkyLLH Analysis

Source	\hat{n}_s	γ	TS	$-\log_{10}(p)$
NGC 1068	56	3.15	19.56	4.25(5.04 σ)
PKS 1424+024	49	3.86	13.29	2.89(4.03 <i>o</i>)
TXS 0506+056	15	2.17	13.18	2.86(4.01 <i>o</i>)

• Neutrinos from point sources are travelling through the CuB

• Neutrinos from point sources are travelling through the CuB



• What if they were to interact?

• What if they were to interact?



• Example: Majoron model (source of neutrino mass)

- Example: Majoron model (source of neutrino mass)
- Also proposed as a way of easing the Hubble Tension (see e.g. PhysRevLett.123.191102

- Example: Majoron model (source of neutrino mass)
- Also proposed as a way of easing the Hubble Tension (see e.g. PhysRevLett.123.191102
- Neutrino interactions mediated by new scalar

- Example: Majoron model (source of neutrino mass)
- Also proposed as a way of easing the Hubble Tension (see e.g. PhysRevLett.123.191102
- Neutrino interactions mediated by new scalar



Relevant Feynman Diagrams. Source: PhysRevD.104.123014



$$egin{aligned} rac{d\phi(t,E_
u)}{dt} &= rac{\partial}{\partial E_
u} [H(t)E_
u\phi(t,E_
u)] \ &- \phi(t,E_
u)\sum_j n_j(t)\sigma_{ij}(E_
u) \end{aligned}$$

$$egin{aligned} rac{d\phi(t,E_
u)}{dt} &= rac{\partial}{\partial E_
u} [H(t)E_
u\phi(t,E_
u)] \ &- \phi(t,E_
u) \sum_j n_j(t)\sigma_{ij}(E_
u) \ &+ \sum_{jkl} n_j(t) \int_{E_
u}^\infty dE_
u'\phi(t,E_
u') rac{d\sigma_{jk
ightarrow il}}{dE_
u}(E_
u',E_
u) \end{aligned}$$

$$egin{aligned} rac{d\phi(t,E_
u)}{dt} &= rac{\partial}{\partial E_
u} [H(t)E_
u\phi(t,E_
u)] \ &- \phi(t,E_
u) \sum_j n_j(t)\sigma_{ij}(E_
u) \ &+ \sum_{jkl} n_j(t) \int_{E_
u}^\infty dE_
u'\phi(t,E_
u') rac{d\sigma_{jk
ightarrow il}}{dE_
u}(E_
u',E_
u) \end{aligned}$$

$$egin{aligned} rac{d\phi(t,E_
u)}{dt} &= rac{\partial}{\partial E_
u} [H(t)E_
u\phi(t,E_
u)] \ &- \phi(t,E_
u)\sum_j n_j(t)\sigma_{ij}(E_
u) \ &+ \sum_{jkl} n_j(t)\int_{E_
u}^\infty dE_
u'\phi(t,E_
u')rac{d\sigma_{jk
ightarrow il}}{dE_
u}(E_
u',E_
u) \end{aligned}$$



Effect of SI on flux from TXS

• Analysis performed using SkyLLH

- Analysis performed using SkyLLH
- Public data covering 10 years of IceCube

- Analysis performed using SkyLLH
- Public data covering 10 years of IceCube
- Produce f_S from ϕ

- Analysis performed using SkyLLH
- Public data covering 10 years of IceCube
- Produce f_S from ϕ



• Use power-law point source as null hypothesis:

$$TS(ec{d}_{ ext{src}}, m_{\phi}, g) = -2\log\left(rac{\sup_{n_s, \gamma} \mathcal{L}(n_s, \gamma, ec{d}_{ ext{src}} | \{ec{x}_i\})}{\sup_{n_s, \gamma} \mathcal{L}(n_s, \gamma, ec{d}_{ ext{src}} | \{ec{x}_i\}, m_{\phi}, g)}
ight)$$

• Use power-law point source as null hypothesis:

$$TS(ec{d}_{ ext{src}}, m_{\phi}, g) = -2\log\left(rac{\sup_{n_s, \gamma} \mathcal{L}(n_s, \gamma, ec{d}_{ ext{src}} | \{ec{x}_i\})}{\sup_{n_s, \gamma} \mathcal{L}(n_s, \gamma, ec{d}_{ ext{src}} | \{ec{x}_i\}, m_{\phi}, g)}
ight)$$

• Maximise TS to fit the new flux to data

RESULTS



CONCLUSIONS

CONCLUSIONS

• Astrophysical neutrino fluxes have unique properties (e.g. high energies, large distances)

CONCLUSIONS

- Astrophysical neutrino fluxes have unique properties (e.g. high energies, large distances)
- This allows us to probe fundamental questions about the nature of neutrinos
CONCLUSIONS

- Astrophysical neutrino fluxes have unique properties (e.g. high energies, large distances)
- This allows us to probe fundamental questions about the nature of neutrinos
- As statistics grow over time, the power of these analyses will only improve

CONCLUSIONS

- Astrophysical neutrino fluxes have unique properties (e.g. high energies, large distances)
- This allows us to probe fundamental questions about the nature of neutrinos
- As statistics grow over time, the power of these analyses will only improve
- Future experiments (e.g. JUNO, IceCube Gen2) will improve our knowledge by orders of magnitude!

BACKUP SLIDES

EVENTS