THE RELATIVISTIC GREEN'S FUNCTION MODEL FOR CCQE and NCE SCATTERING

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nuclear response to the electroweak probe



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QE-peak dominated by one-nucleon knockout

QE e-nucleus scattering

 $e + A \Longrightarrow e' + N + (A - 1)$

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- both e' and N detected (A-1) discrete eigenstate n exclusive (e,e'p)
- only e' detected, all final nuclear states included inclusive (e,e')

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FINAL-STATE INTERACTION between the emitted nucleon and the residual nucleus

EXCLUSIVE SCATTERING: FSI

RDWIA

FSI described by a complex OP with an imaginary absorptive part. The imaginary part gives a reduction of the calculated c.s. which is essential to reproduce data

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rROP	only the real part of the OP: conserves the flux but it is conceptually wrong
RMF	RELATIVISTIC MEAN FIELD: same real energy-independent potential of bound states
RGF	GREEN'S FUNCTION complex OP conserves the flux consistent description of FSI in exclusive and inclusive QE electron scattering

with suitable approximations (basically related to the IA) the components of the inclusive response can be written in terms of the s.p. optical model Green's function

the explicit calculation of the s.p. Green's function can be avoided by its spectral representation which is based on a biorthogonal expansion in terms of the eigenfunctions of the non Herm optical potential V and V⁺

matrix elements similar to RDWIA

scattering states eigenfunctions of V and V⁺ (absorption and gain of flux): the imaginary part redistributes the flux and the total flux is conserved

$$W^{\mu\mu}(\omega,q) = \sum_{n} \left[\mathbf{Re} T_{n}^{\mu\mu}(E_{\mathbf{f}} - \varepsilon_{n}, E_{\mathbf{f}} - \varepsilon_{n}) - \frac{1}{\pi} \mathcal{P} \int_{M}^{\infty} \mathbf{d}\mathcal{E} \frac{1}{E_{\mathbf{f}} - \varepsilon_{n} - \mathcal{E}} \mathbf{Im} T_{n}^{\mu\mu}(\mathcal{E}, E_{\mathbf{f}} - \varepsilon_{n}) \right]$$

$$W^{\mu\mu}(\omega,q) = \sum_{n} \left[\mathbf{Re} \left(\mathbf{f}_{n}^{\mu\mu}(E_{\mathbf{f}} - \varepsilon_{n}, E_{\mathbf{f}} - \varepsilon_{n}) - \frac{1}{\pi} \mathcal{P} \int_{M}^{\infty} \mathbf{d}\mathcal{E} \frac{1}{E_{\mathbf{f}} - \varepsilon_{n} - \mathcal{E}} \mathbf{Im} \left(\mathbf{f}_{n}^{\mu\mu}(\mathcal{E}, E_{\mathbf{f}} - \varepsilon_{n}) \right) \right] \mathbf{f}_{n}^{\mu\mu}(\mathcal{E}, E) = \lambda_{n} \langle \varphi_{n} \mid j^{\mu\dagger}(\mathbf{q}) \sqrt{1 - \mathcal{V}'(E)} \mid \tilde{\chi}_{\mathcal{E}}^{(-)}(E) \rangle \langle \chi_{\mathcal{E}}^{(-)}(E) \mid \sqrt{1 - \mathcal{V}'(E)} j^{\mu}(\mathbf{q}) \mid \varphi_{n} \rangle$$



$$W^{\mu\mu}(\omega,q) = \sum_{n} \left[\mathbf{Re} \prod_{n}^{\mu\mu} (E_{\mathbf{f}} - \varepsilon_{n}, E_{\mathbf{f}} - \varepsilon_{n}) - \frac{1}{\pi} \mathcal{P} \int_{M}^{\infty} \mathbf{d}\mathcal{E} \frac{1}{E_{\mathbf{f}} - \varepsilon_{n} - \mathcal{E}} \mathbf{Im} \prod_{n}^{\mu\mu} (\mathcal{E}, E_{\mathbf{f}} - \varepsilon_{n}) \right]$$

$$(T_{n}^{\mu\mu}(\mathcal{E}, E) = \lambda_{n} \langle \varphi_{n} \mid j^{\mu\dagger}(q) \sqrt{1 - \mathcal{V}'(E)} \mid \tilde{\chi}_{\mathcal{E}}^{(-)}(E) \langle \chi_{\mathcal{E}}^{(-)}(E) \mid \sqrt{1 - \mathcal{V}'(E)} j^{\mu}(q) \mid \varphi_{n} \rangle$$

$$\langle \chi^{(-)}(E) \mid j^{\mu}(q) \mid \varphi_{n} \rangle$$

DWIA exclusive (e,e'p)

 $\langle \chi^{(-)}(E) \mid j^{\mu}(q) \mid \varphi_n \rangle$

- j^µ one-body nuclear current
- $\mathbf{\Phi}_{n}$ s.p. bound state overlap function
- $\chi^{(-)}$ s.p. scattering w.f. eigenfunction of an OP

$$W^{\mu\mu}(\omega,q) = \sum_{n} \left[\mathbf{Re} \left(\mathbf{f}_{n}^{\mu\mu}(E_{\mathbf{f}} - \varepsilon_{n}, E_{\mathbf{f}} - \varepsilon_{n}) - \frac{1}{\pi} \mathcal{P} \int_{M}^{\infty} \mathbf{d}\mathcal{E} \frac{1}{E_{\mathbf{f}} - \varepsilon_{n} - \mathcal{E}} \mathbf{Im} \left(\mathbf{f}_{n}^{\mu\mu}(\mathcal{E}, E_{\mathbf{f}} - \varepsilon_{n}) \right) \right] \mathbf{f}_{n}^{\mu\mu}(\mathcal{E}, E) = \lambda_{n} \langle \varphi_{n} \mid j^{\mu\dagger}(q) \sqrt{1 - \mathcal{V}'(E)} \mid \tilde{\chi}_{\mathcal{E}}^{(-)}(E) \rangle \langle \chi_{\mathcal{E}}^{(-)}(E) \mid \sqrt{1 - \mathcal{V}'(E)} j^{\mu}(q) \mid \varphi_{n} \rangle$$



eigenfunctions of V and V⁺







redistribution of the flux among the different channels

Relativistic Green's Function Model

- consistent treatment of FSI in the exclusive and in the inclusive scattering
- the imaginary part of the ROP includes inelastic channels
- with a complex ROP the model can include contributions not included in other models based on the IA
- the use of a phen. ROP does not allow us to disentangle specific contributions
- different phen ROP's available, theoretical uncertainties in the numerical predictions of the model

RGF: comparison with QE (e,e') data





$$E_0 = 1080 \text{ MeV} \ \vartheta = 32^{\circ}$$

$$E_0 = 841 \text{ MeV} \quad \vartheta = 45.5^{\circ}$$

$$E_0 = 2020 \text{ MeV} \quad \vartheta = 20^\circ$$

RGF: comparison CCQE data

$\nu_l(\bar{\nu}_l) + A \longrightarrow l^-(l^+) + N + (A-1)$

RGF: comparison CCQE data

$$\nu_l(\bar{\nu}_l) + A \longrightarrow l^-(l^+) + N + (A-1)$$

only final lepton detected inclusive CC

Differences between Electron and Neutrino Scattering

electron scattering :

beam energy known, cross section as a function of $\ \omega$

neutrino scattering:

beam energy and $\omega\,$ not known

calculations over the energy range relevant for the neutrino flux

the flux-average procedure can include contributions from different kinematic regions where the neutrino flux has significant strength, contributions other than direct 1-nucleon emission

Comparison with MiniBooNe CCQE data

 $^{12}C(\nu_{\mu},\mu^{-})$

 $0.4 < \cos\theta_{\mu} < 0.5$



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Comparison with MiniBooNe CCQE data



Comparison MiniBooNE CCQE neutrinoantineutrino scattering



Comparison MiniBooNE CCQE neutrinoantineutrino scattering

$$12C(\nu_{\mu},\mu^{-})$$

 $[10^{-39} \text{cm}^2]$

ь

 ${}^{12}C(\bar{\nu}_{\mu},\mu^+)$



Comparison MINERvA CCQE neutrinoantineutrino scattering

- higher energy (energy range 1.5-10 GeV)
- models based on the IA underpredict the MiniBooNE data but in general provide a good description of the MINERvA data
- RGF...?

Comparison MINERvA CCQE neutrinoantineutrino scattering



NC v-nucleus scattering

- only the outgoing nucleon is detected: semi-inclusive scattering
- FSI?
- RDWIA: sum of all integrated exclusive 1NKO channels with absorptive imaginary part of the ROP. The imaginary part accounts for the flux lost in each channel towards other inelastic channels. Some of these reaction channels are not included in the experimental cross section when one nucleon is detected. For these channels RDWIA is correct, but there are channels excluded by the RDWIA and included in the experimental c.s.
- RGF recovers the flux lost to these channels but can include also contributions of channels not included in the semi-inclusive cross section
- we can expect RDWIA smaller and RGF larger than the experimental cross sections
- relevance of contributions neglected in RDWIA and added in RGF depends on kinematics

Comparison with MiniBooNE NCE data



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conclusions

RGF

- describes FSI in the inclusive lepton-nucleus scattering
- developed for inclusive QE electron-nucleus scattering, successfully tested in comparison with (e,e') data and then applied to QE neutrino-nucleus scattering
- describes CCQE and NCE MiniBooNE data, MINERvA CCQE data
- the imaginary part of the ROP includes the overall effect of inelastic channels (rescattering, non-nucleonic, multi-nucleon,....)
- the role of different inelastic processes cannot be disentangled
- the use of different phenomenological ROP's may introduce theoretical uncertainties on the RGF results
- better determination of the ROP desirable
- 2p-2h MEC....?