

# The Low Energy Neutrino Factory: Physics Performance

Tracey Li  
IPPP, Durham University



5th IDS-NF Meeting  
Fermilab, Illinois  
8th April 2010

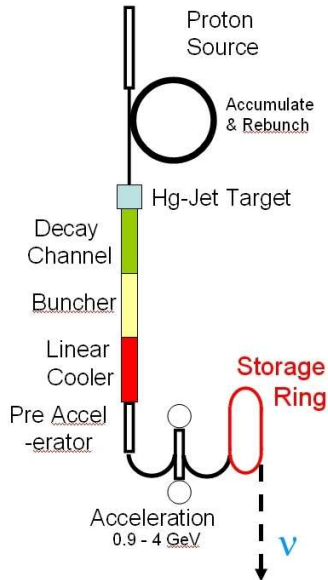
In collaboration with: Alan Bross, Malcolm Ellis, Enrique  
Fernández-Martínez, Steve Geer, Olga Mena and Silvia Pascoli

# Talk outline

- Review of LENF setup and motivation
- Standard oscillations
- Non-standard interactions
- Conclusions.

# LENF review: beam and baseline

- $E_\mu \sim 4.5$  GeV
- $L = 1300$  km
- $1.4 \times 10^{21}$  muon decays per year per polarity
- Running time: 5 + 5 years
- $1.4 \times 10^{22}$  muon decays in total.



# LENF review: detector

Detector options: **20 kton** totally active scintillating detector (**TASD**) or **100 kton** Liquid Argon (**LAr**).

	<b>TASD</b>	<b>LAr</b>
<b>Fiducial mass</b>	20 kton	100 kton
<b>Threshold</b>	0.5 GeV	0.5 GeV
<b>Efficiency (<math>\mu^\pm</math>)</b>	94%	80%
<b>Efficiency (<math>e^\pm</math>)</b>	47%	80%
<b>Systematics</b>	2%	2 – 5%
<b>Energy resolution</b>	10%	5%
<b>Background (<math>\mu^\pm</math>)</b>	$1 \times 10^{-3}$	$1 - 5 \times 10^{-3}$
<b>Background (<math>e^\pm</math>)</b>	$10^{-2}$	$10^{-2} - 0.8$

**Channels:**

$\nu_\mu \rightarrow \nu_\mu$ ,  $\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu$  - **disappearance**

$\nu_e \rightarrow \nu_\mu$ ,  $\bar{\nu}_e \rightarrow \bar{\nu}_\mu$  - **golden**

$\nu_\mu \rightarrow \nu_e$ ,  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  - **platinum**.

The **low energy neutrino factory** was proposed as a next-generation long-baseline experiment in the scenario that  $\theta_{13}$  is large.

S. Geer, O. Mena, S. Pascoli, Phys. Rev. D **75**, 093001 (2007).

**Motivation:** The oscillation spectrum for energies  $\lesssim 5$  GeV is very rich at  $\sim 1300$  km, providing sensitivity to  $\theta_{13}$ ,  $\delta$ , the **mass hierarchy** and  $\theta_{23}$ .

**But why is the LENS suitable for large  $\theta_{13}$ ?**

# $\theta_{13}$ dependence of oscillation probability

This is the oscillation probability for the golden channel,  $\nu_e \rightarrow \nu_\mu$ :

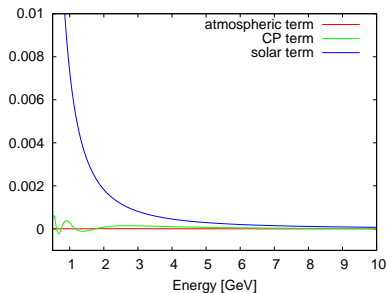
$$\begin{aligned} P_{\nu_e \nu_\mu} = & \textcolor{red}{s_{213}^2 s_{23}^2} \sin^2 \left( \frac{\Delta m_{31}^2 L}{4E} - \frac{AL}{2} \right) \\ & + \textcolor{green}{s_{213} \alpha s_{212} s_{223}} \frac{\Delta m_{31}^2 L}{2EA} \sin \left( \frac{AL}{2} \right) \sin \left( \frac{\Delta m_{31}^2 L}{4E} - \frac{AL}{2} \right) \\ & \times \cos \left( \delta - \frac{\Delta m_{31}^2 L}{4E} \right) \\ & + \textcolor{blue}{\alpha^2 c_{23}^2 s_{212}^2} \left( \frac{\Delta m_{31}^2 L}{2EA} \right)^2 \sin^2 \left( \frac{AL}{2} \right). \end{aligned}$$

- The **atmospheric term** contains information on  $\theta_{13}$  and the mass hierarchy.
- The **CP term** contains information on  $\theta_{13}$ ,  $\delta$  and the mass ordering.
- The **solar term** doesn't tell us anything interesting.

# $\theta_{13}$ dependence of oscillation spectrum

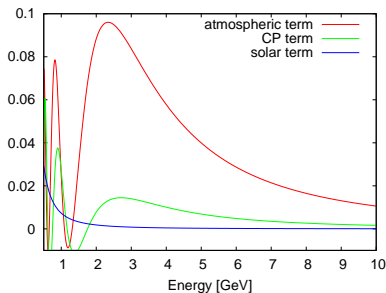
This is how each of the terms vary with the value of  $\theta_{13}$ :

$$\theta_{13} = 0.1^\circ$$



GLOBES 3.0

$$\theta_{13} = 10^\circ$$



GLOBES 3.0

The solar term is dominant when  $\sin^2(2\theta_{13}) \lesssim 10^{-3}$ .

The atmospheric term is dominant when  $\sin^2(2\theta_{13}) \gtrsim 10^{-3}$ .

# Long-baseline matter effects vs CP violation

When the **solar term** is dominant, a long baseline and high energy is the only way to determine the mass hierarchy  $\Rightarrow$  HENF.

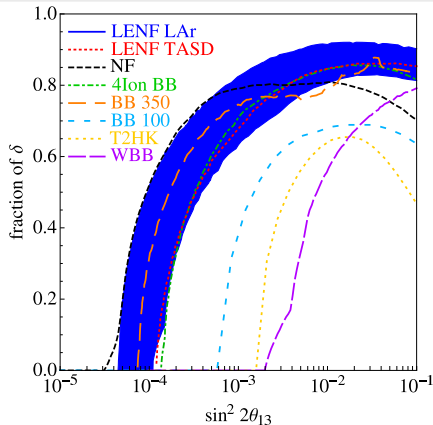
But if the **atmospheric term** is dominant:

- Measurements can be extracted more easily as the oscillations are easier to see.
- **Matter effects and CPV at a HENF become difficult to distinguish.**

$\Rightarrow$  **A shorter baseline with less matter effect is preferable for large  $\theta_{13}$ .**



# LENF sensitivity: CP discovery

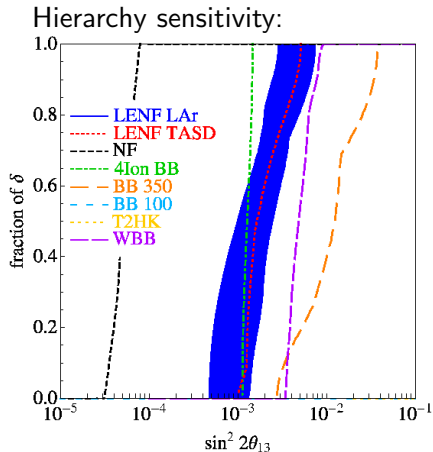
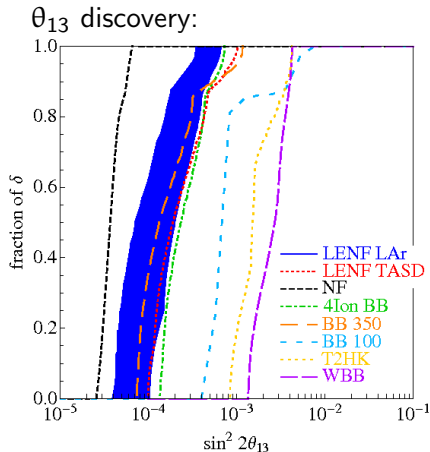


arXiv: 0911.3776

GLoBES 3.0

- The **high energy neutrino factory (NF)** was designed for the scenario that  $\theta_{13}$  is **very small**.
- But the **low energy neutrino factory (LENF)** performs better if  $\sin^2(2\theta_{13}) \gtrsim 10^{-3}$ .

# LENF sensitivity: $\theta_{13}$ and hierarchy



arXiv: 0911.3776

GLOBES

For  $\sin^2(2\theta_{13}) > 1 \times 10^{-3}$  - 100% coverage for  $\theta_{13}$  discovery.

For  $\sin^2(2\theta_{13}) > 4 \times 10^{-3}$  - 100% coverage for hierarchy sensitivity.

# The platinum channels

A full standard oscillation analysis can be found in [arXiv:0911.3776](#).

This includes a study of the effect of adding the **platinum channels** ( $\nu_\mu \rightarrow \nu_e, \bar{\nu}_\mu \rightarrow \bar{\nu}_e$ ).

These channels are experimentally very challenging to observe.

**So precisely how important are the platinum channels?**

# Platinum channels vs statistics

Compare the effect of the **platinum channels** with the effect of **increasing statistics**. Consider 3 scenarios:

- $2.8 \times 10^{22}$  total muon decays, no platinum channel  
Most optimistic statistics - re-optimized accelerator front-end  
 $\Rightarrow$  40% increase, and double the running time.
- $1.0 \times 10^{22}$  total muon decays, no platinum channel  
Same statistics as HENF, with all muons to a single baseline
- $1.0 \times 10^{22}$  total muon decays, with platinum channel.

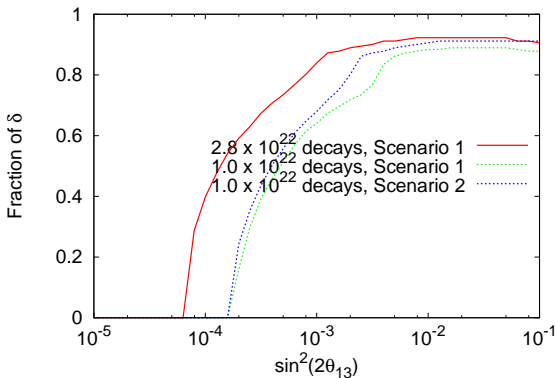
See also J. Tang and W. Winter, Phys. Rev. D **81**, 033005 (2010).

Platinum channel assumptions:

- Efficiency = 37% for  $E < 1$  GeV, 47% for  $E > 1$  GeV
- Background (charge mis-ID rate) =  $10^{-2}$ .

# Platinum channel: CP discovery

**Scenario 1:** no platinum channels, **Scenario 2:** with platinum channels

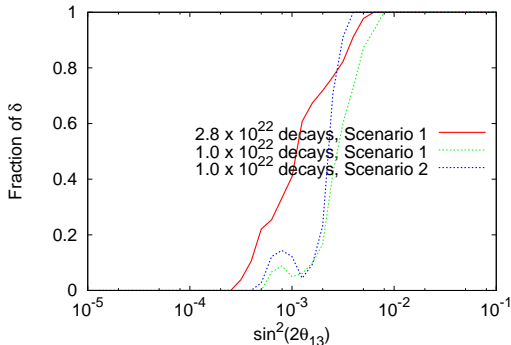


GLOBES

- For  $1.0 \times 10^{22}$  decays and  $\sin^2(2\theta_{13}) \sim 10^{-3}$ , the platinum channel increases CPV sensitivity by  $\sim 10\%$ .
- But higher statistics are always better.

# Platinum channel: hierarchy sensitivity

**Scenario 1:** no platinum channels, **Scenario 2:** with platinum channels



GLoBES

- $\sin^2(2\theta_{13}) \lesssim 4 \times 10^{-3}$ : higher statistics are better.
- $4 \times 10^{-3} \lesssim \sin^2(2\theta_{13}) \lesssim 10^{-2}$ : the platinum channel helps.
- $\sin^2(2\theta_{13}) > 10^{-2}$ : the platinum channel is unnecessary.

# Standard oscillation summary

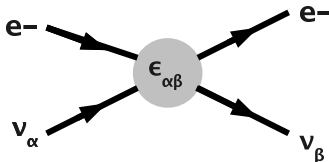
- The LENF has better sensitivity to CP violation than the HENF for  $\sin^2(2\theta_{13}) \gtrsim 10^{-3}$ .
- The LENF has 100% coverage to  $\theta_{13}$  and the mass hierarchy for  $\sin^2(2\theta_{13}) \gtrsim 4 \times 10^{-3}$ .
- The performance of the LENF is very sensitive to statistics.  
J. Tang and W. Winter, Phys. Rev. D **81**, 033005 (2010).
- If the LENF is only to be considered for  $\sin^2(2\theta_{13}) > 10^{-2}$ , the platinum channel adds only a little to the CP sensitivity.

# Non-standard interactions: introduction

We would also like to search for new physics with a  $\nu$  factory, for example **non-standard interactions** (NSI's).

- NSI's are effective 4-point flavour-changing interactions.
- NSI's can be parameterized as  $\epsilon_{\alpha\beta}$  (model-independent) which describe the rate of the transition  $\nu_{\alpha} \rightarrow \nu_{\beta}$ .

T. Ota, J. Sato and N. Yamashita, Phys. Rev. D **65**, 093015 (2002).





# NSI's at LENF

The LENF has leading order sensitivity to the NSI parameters  $\varepsilon_{e\mu} e^{i\phi_{e\mu}}$  and  $\varepsilon_{e\tau} e^{i\phi_{e\tau}}$ :

$$\begin{aligned}
 P_{\nu_e \rightarrow \nu_\mu} = & \quad s_{213}^2 s_{23}^2 \sin^2 \left( \frac{\Delta m_{31}^2 L}{4E} - \frac{AL}{2} \right) \\
 & + s_{213} \alpha s_{212} s_{223} \frac{\Delta m_{31}^2 L}{2EA} \sin \left( \frac{AL}{2} \right) \sin \left( \frac{\Delta m_{31}^2 L}{4E} - \frac{AL}{2} \right) \cos \left( \delta - \frac{\Delta m_{31}^2 L}{4E} \right) \\
 & + \alpha^2 c_{23}^2 s_{212}^2 \left( \frac{\Delta m_{31}^2 L}{2EA} \right)^2 \sin^2 \left( \frac{AL}{2} \right) \\
 & - 4\varepsilon_{e\tau} s_{213} c_{23} s_{23}^2 \sin \left( \frac{AL}{2} \right) \sin \left( \frac{\Delta m_{31}^2 L}{4E} - \frac{AL}{2} \right) \cos \left( \delta + \phi_{e\tau} - \frac{\Delta m_{31}^2 L}{4E} \right) \\
 & + 4\varepsilon_{e\tau} \alpha s_{212} c_{23}^2 s_{23} \sin \left( \frac{AL}{2} \right) \sin \left( \frac{\Delta m_{31}^2 L}{4E} - \frac{AL}{2} \right) \cos \left( \phi_{e\tau} + \frac{\Delta m_{31}^2 L}{4E} \right) \\
 & + 4\varepsilon_{e\tau}^2 c_{23}^2 s_{23}^2 \sin^2 \left( \frac{AL}{2} \right) \\
 & - 4\varepsilon_{e\mu} s_{213} c_{23}^2 s_{23} \sin \left( \frac{AL}{2} \right) \sin \left( \frac{\Delta m_{31}^2 L}{4E} - \frac{AL}{2} \right) \cos \left( \delta + \phi_{e\mu} - \frac{\Delta m_{31}^2 L}{4E} \right) \\
 & - 4\varepsilon_{e\mu} \alpha s_{212} c_{23}^2 s_{23} \sin \left( \frac{AL}{2} \right) \sin \left( \frac{\Delta m_{31}^2 L}{4E} - \frac{AL}{2} \right) \cos \left( \phi_{e\mu} + \frac{\Delta m_{31}^2 L}{4E} \right) \\
 & + 4\varepsilon_{e\mu}^2 c_{23}^2 s_{23}^2 \sin^2 \left( \frac{AL}{2} \right).
 \end{aligned}$$

# Degeneracies and NSI's

When we include NSI's, the parameter space is vastly increased and so the **degeneracy problem** is magnified.

With SO parameters alone we already have a degeneracy problem:

- Data can be fitted to different combinations of  $(\theta_{13}, \delta, \text{sign}(\Delta m_{31}^2))$ .
- From a single measurement, we cannot tell which is the true solution (see next slide).

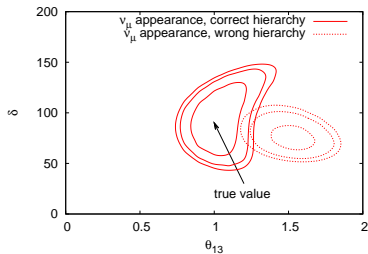
⇒ This severely weakens the precision of measurements.

## Possible solutions:

- Combine information from **complementary channels** (**LENF**).
- Use a **magic baseline** (**HENF**).

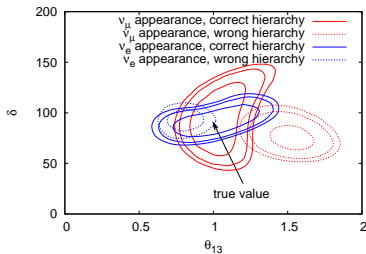
# Using complementary channels to resolve degeneracies

Only  $\nu_\mu$  appearance channel:



GLoBES

$\nu_\mu$  and  $\nu_e$  appearance channels:



The degenerate solutions appear in different regions of parameter space for each channel.

Thus we can eliminate the fake solutions by combining appropriate channels.

The same technique applies to resolving SO-NSI degeneracies.

# Degeneracies and NSI's

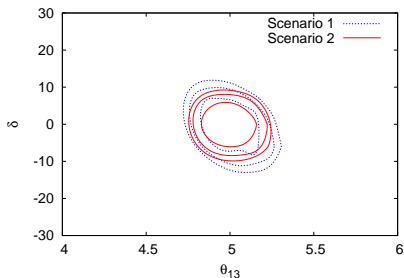
- We must ensure that NSI's do not degrade our sensitivity to the oscillation parameters.
- The high-energy  $\nu$  factory is already immune to this problem because of the magic baseline (sets  $\sin(\frac{AL}{2})$  to zero  $\Rightarrow$  all NSI terms vanish).
- But for the LBNF, we need a solution. We find that the platinum channel ( $\nu_\mu \rightarrow \nu_e$ ) enhances the sensitivity to all parameters.

# Effect of NSI's on standard oscillation measurements

Even if all the NSI's are zero, marginalising over them weakens the precision of the oscillation measurements.

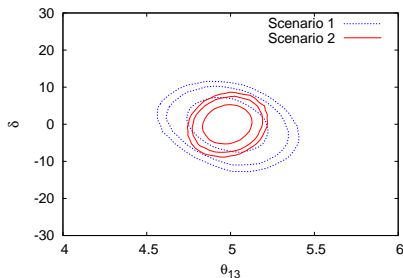
Scenario 1: no platinum channel, Scenario 2: platinum channel included (eff = 47%, bkgd =  $10^{-2}$ ).

Marginalisation over SO parameters only



MonteCUBES

Marginalisation over SO and NSI parameters



MonteCUBES

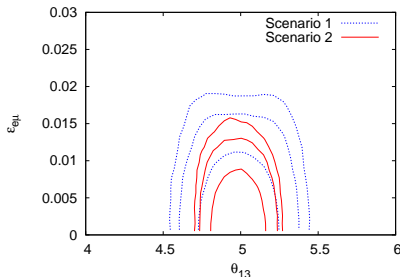
**The platinum channel helps to maintain the sensitivity to the SO parameters.**

# Bounds on $\varepsilon_{e\mu}$ and $\varepsilon_{e\tau}$ : $\theta_{13} = 5^\circ$

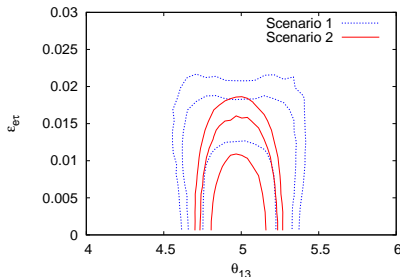
The current upper bounds on  $\varepsilon_{e\mu}$  and  $\varepsilon_{e\tau}$  are  $\sim 10^{-1}$ .

C. Biggio, M. Blennow, E. Fernández-Martínez, JHEP 0908, 090 (2009).

Simulate  $\varepsilon_{e\mu} = \varepsilon_{e\tau} = 0$  and  $\theta_{13} = 5^\circ$ :



MonteCUBES

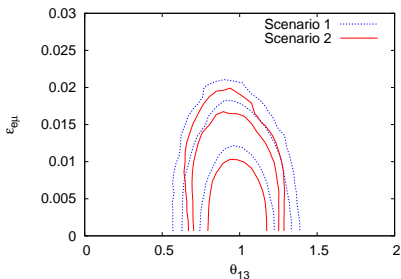


MonteCUBES

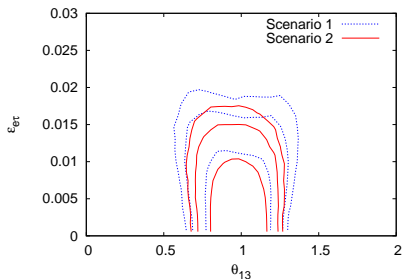
- For  $\theta_{13} = 5^\circ$ , the platinum channel enhances sensitivities.
- We can obtain a 95% upper bound of  $\sim 10^{-2}$  on  $\varepsilon_{e\mu}$  and  $\varepsilon_{e\tau}$  (HENF can reach  $\sim 10^{-3}$ ).

# Bounds on $\varepsilon_{e\mu}$ and $\varepsilon_{e\tau}$ : $\theta_{13} = 1^\circ$

Do the same for  $\theta_{13} = 1^\circ$ :



MonteCUBES



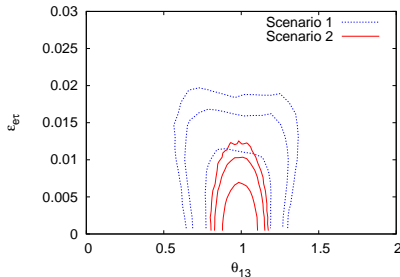
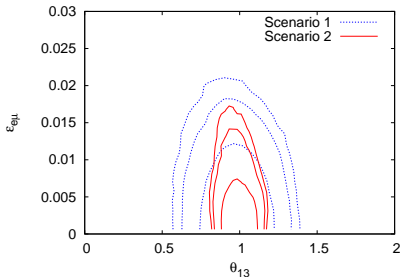
MonteCUBES

The platinum channels are not as effective when  $\theta_{13}$  is small.

We find that the **results are not improved by increasing statistics** (unlike for oscillation measurements).

# Effect of a perfect platinum channel

But if the platinum channels have (hypothetically!) perfect efficiency and no background:



MonteCUBES

MonteCUBES

A perfect platinum channel would enhance sensitivities for all values of  $\theta_{13}$ .

⇒ **There is a critical number of platinum events needed for these channels to be useful.**



- The LENF sensitivity to NSI's is limited by the oscillation-NSI degeneracy.
- To resolve the degeneracy, complementary information from an additional channel or an additional baseline is required.
- Statistics are not important for NSI measurements as they do not help to resolve the degeneracies.
- For the assumed efficiency (47%) and background ( $10^{-2}$ ), the platinum channel is only helpful if  $\theta_{13} \sim 5^\circ$ .
- If the performance can be sufficiently improved (difficult...) the platinum channel will also be useful for small  $\theta_{13}$ .

See also D. Meloni, T. Ohlsson, W. Winter and H. Zhang, arXiv:0912.2735.

# Conclusions

- For  $\sin^2(2\theta_{13}) > 10^{-2}$  and  $1.4 \times 10^{22}$  decays, the LENS has
  - 100%  $\theta_{13}$  discovery potential and hierarchy sensitivity
  - $\sim 80\%$  CP discovery potential.
- The platinum channel will increase the CP sensitivity by a few % but adds nothing else to the oscillation sensitivity.
- The LENS has sensitivity to matter NSI's down to  $\sim 10^{-2}$ .
- The sensitivity to NSI's is limited by degeneracies.
- To resolve the degeneracies we need to include the platinum channels, or a second detector.
- The current estimates of the platinum channel performance indicate that it will only be useful for very large  $\theta_{13}$ .